XANES characterization of UO₂/Mo(Pd) thin films as models for ε-particles in spent nuclear fuel

M A Denecke¹, T Petersmann², R Marsac¹, K Dardenne¹, T Vitova¹, T Prüßmann¹, M Borchert³, U Bösenberg³, G Falkenberg³ and G Wellenreuther³

¹Karlsruhe Institute of Technology, Institut für Nukleare Entsorgung, Postfach 3640, D-76021 Karlsruhe, Germany
²European Joint Research Center, Institute for Transuranium Elements (ITU), D-76344 Eggenstein-Leopoldshafen, Germany
³Deutsches Elektronen-Synchrotron (DESY), Notkestr. 85, D-22607 Hamburg, Germany

E-mail: melissa.denecke@kit.edu

Abstract. X-ray absorption near edge structure (XANES) is recorded for nano- and micro-particles formed in Mo doped UO₂ thin films (TFs) prepared by in situ sputter co-deposition and 1000°C tempering. These UO₂/Mo TFs are intended to serve as synthetic models for ε-particles in spent nuclear fuel. We find that when Si is used as substrate, nano-sized zero-valent Mo particles form as desired. However, these are embedded in USi₃, which forms at the high temper temperatures. Micron-sized Mo-particles are formed when SiO₂ is used as substrate. Using focussed X-ray beams of varying size (500µm, 25µm and 5µm), these particles are characterized to be predominantly hexavalent Mo oxides, potentially with tetrahedral coordinated [Mo(VI)O₄]²⁻ at the surface and a MoO₃-like phase in the bulk. These TFs are poor synthetic models for spent fuel ε-particles but do offer the opportunity to study changes in surface structures in response to stress/charge as a function of particle size.

1. Introduction
Spent nuclear fuel (SNF) consists of 95% UO₂ and 5% mixture of radionuclides (RNs) and fission products, heterogeneously distributed throughout the fuel matrix in different phases, e.g., gases (Xe, Kr, I), oxides (including transuranic elements Pu, Np, etc.) and metallic precipitates or so-called ε-particles (Pd, Mo, Rh, Ru, etc.[1,2]. Systematic investigations of SNF matrix corrosion as an important source term for the mobilization of RNs are essential for evaluating SNF disposal safety [3,4]. We prepare UO₂/Mo thin films (TFs) for potential use as models to study the influence of ε-particles on the surface corrosion of SNF. The TFs provide a less complex system and lower radiation fields than actual SNF samples, which allow systematic mechanistic investigations even in non-radioactive laboratories.

We determine how well these TF models mimic the SNF surface by thoroughly characterizing their composition, structure and morphology using numerous methods. Studies of UO₂/Mo TFs having varying Mo content prepared via in situ sputter co-deposition on different substrates (Si, SiO₂ or LaAlO₃) by means of 2-dimensional (2D) and 3-dimensional synchrotron-based imaging methods (scanning nano-/micro-XRF, scanning nano-XRD and XRF, holographic and ptychographic tomography) reveal that nano- and micro-sized Mo-particles form in TFs prepared on Si and SiO₂.
substrates and tempered at 1000°C (600 °C was not sufficient); when LaAlO3 is used as substrate and tempered at 1000°C, Mo does not aggregate but is found concentrated at the substrate surface [5]. In this report we characterize the Mo-particles observed to form in these TFs using X-ray absorption near-edge structure (XANES) at the Mo K-edge and the L3-edges of both Mo and U.

2. Experimental

Mo L3 edge spectra are recorded at the INE-Beamline for actinide research [6] at the ANKA synchrotron in fluorescence mode, using the low energy setup available there and a Si(111) double crystal monochromator (DCM) calibrated against the first inflection point in the Mo L3-XANES of a Mo0-foil (2.52 keV). Both scanning XRF measurements and XANES measurements using a focused X-ray beam are performed at the hard X-ray Micro/Nanoprobe beamline P06 at PETRA III [7] and at HASYLAB beamline-L experimental station (DORIS at DESY). The X-ray beam is focused at P06 using a Kirkpatrick-Baez (KB) system equipped with Rh-coated silica mirrors, which in our experiment delivered an enlarged focal spot size of approximately 2µm (H) × 7µm (V) due to monochromator instabilities. At beamline-L a polycapillary half-lens is used to focus the beam to 35 µm (H) × 22 µm (V) at 23 keV. At both beamlines a Si(111) DCM is used and the beam energy calibrated against the first inflection point in the Mo K-XANES of a Mo0-foil reference, defined as 20.0 keV.

Two UO2 TFs doped with 8 and 20 wt% Mo and around 500 nm thick are prepared via co-deposition of U and Mo onto SiO2 and Si substrates and tempering, designated as sample I and H, respectively (sample preparation is described in [5]). Samples are mounted at a 45° angle to the incident beam on a positioning stage for scanning in the y and z directions (2D) and positioning the x-axis into the beam focus. A Vortex-90EX silicon drift detector (SII Nanotechnology USA) is used at the INE-Beamline and beamline-L and a Vortex-EM at P06 for collecting X-ray fluorescence radiation. Mo K (U L3) edge XANES are recorded from Mo (U)-rich areas identified in distribution maps reconstructed from the scanning µ-XRF data.

3. Results and discussion

In previous µXRF/µ-XRD studies performed at the ESRF ID22Ni using a 100 nm focussed beam, we found that for TFs prepared on Si substrates and tempered at 1000°C (sample H) the UO2 partially reacts with the substrate, thereby forming USi3 with Mo localized as 100-400 nm nanoparticles in USi3 grain boundaries [5]. Uranium L3 XANES recorded at beamline-L using a ~25 µm focussed-beam at different sample areas identified to be composed of UO2, USi3 and a mixture of these two phases are depicted in Figure 1 (left). The XANES from the UO2 area (U-A) is similar to a bulk UO2 reference spectrum but with a broadened white line at ~17180 eV, as well as indication of a lengthening of the average U-O distance, evident as a shift of the first EXAFS resonance at ~17216 eV to lower energy (due to the relation ΔE ∝ 1/R(U-O)2 [8]). The U L3 XANES spectra measured at the sample area comprised the mixed UO2/USi3 phases (U-C) are more similar to
the XANES recorded in an area composed solely of USi$_3$ (U-B) suggesting that U is mostly present as USi$_3$ in the mixed phase.

All Mo K XANES recorded at various areas randomly distributed across sample H lie within 1.2 eV to one another and although their edge energies are near that for Mo(0), their XANES fingerprint does not resemble Mo$^0$-foil at all. The Mo K-XANES recorded at exactly the same sample positions as the U L3 spectra are shown in Figure 1 (right). The edge energies are near that for Mo(0) and well below energies expected for Mo(IV) and Mo(VI) (around 20003.5 and 20007.5 eV [9]). In contrast to the U L3 results above, Mo K XANES recorded in the mixed phase area of sample H (Mo-C) is essentially identical to that from the UO$_2$ area (Mo-A) and less similar to the USi$_3$ sample area XANES (Mo-B). The bulk Mo L3 edge data measured at the INE-Beamline at ANKA for sample H at different positions on the sample is shown in Figure 3 (lower curves, top). The L3 edge spectra line up with that for the Mo(0) L3 XANES and lie below the Mo(IV)O$_2$ and Mo(VI)O$_3$ reference spectra, indicating that the mean valence of Mo nanoparticles observed to form in the tempered UO$_2$/Mo TFs on Si is zero. Similar to the K edge data, the sample H XANES fingerprint does not resemble that for Mo$^0$-foil; the white lines are at least twice as intense and the resonance above the white line visible in the Mo(0) spectrum is absent. Presumably these differences result from the small size of the particles (narrowing of final state 4d band and loss of extended coordination structure).

In Mo distributions reconstructed from scans of small areas of sample I using a nano-sized beam at ESRF ID22-NI, no Mo hot spots were found. In large area XRF scans of sample I (Figure 2) measured at beamline-L with a µ-sized beam, 25-150 µm sized Mo hot spots are revealed, which are distributed randomly and widely spaced in the UO$_2$ TF. A Mo K-XANES recorded from one of the larger hot spots is shown in Figure 1 (right, bottom curve). The spectrum show a single broad main absorption and its edge energy lies approximately at 20007.3 eV, which is close to the value expected for Mo(VI) [9]. Bulk Mo L3 edge data (Figure 3, top) for sample I recorded at two different positions on the sample show well-resolved features, as the life time broadening due to the L3 core hole is considerably reduced. The white line is split into t$_{2g}$ and e$_g$ states for the (distorted) O$_6$ symmetry of the reference oxides; the t$_{2g}$ being partially occupied in Mo(IV) (d$^4$). The sample I spectra qualitatively more resemble the Mo(VI) L3 signal with a strong t$_{2g}$ intensity. The L3 XANES energy position is evaluated from the value of the arctan step function obtained in XANES fits, compared to the reference spectra. The result of this analysis is shown in Figure 3 (bottom). The relationship between reference oxidation state and edge energy was not strictly linear so we modelled the calibration curve with a polynomial function. Using this curve, the Mo oxidation state in sample I determined for the two L3 XANES spectra is 3.8 (I1) and 5.2 (I2). For comparison, fits to the spectra as linear combination of the MoO$_2$ and MoO$_3$ reference spectra gave slightly different mean Mo valence values in I1 and I2, 4.7 and 5.4, respectively. The quality of least square fit results, however, indicates that these two reference compounds are not sufficient to model sample I L3 spectra. This is likely due to the presence of tetrahedral coordinated Mo(VI) in this sample (see below).

In order to explore the reason for measuring different mean Mo oxidation states in one and the same sample I, we use a ~5 µm beam at the P06 microprobe experimental station to measure Mo K XANES of individual Mo particles <5µm in size and compared these to those recorded at different positions of
a particle at least 10 times larger. The results are shown in Figure 4. Particles studied having dimensions of the beam size or smaller exhibit a strong pre-peak at around 20005 eV, which is a signature for tetrahedral coordinated \([\text{Mo(VI)}\text{O}_4]^2\) (cf. the XANES for CaMoO₄). This feature is absent or reduced to a shoulder (similar to MoO₃) on the rising absorption edge in most pixels studied (spectra 1-3), but clearly (and reproducibly) visible in the XANES measured on the outer rim of the particle (#4). The Mo K energy position at half the edge jump of the XANES without a pre-peak all lie about 8 eV above the Mo(0) XANES, indicative of hexavalent molybdenum. We find no evidence for Mo(IV). Summarizing results for this particular particle: we observe a mixture of a MoO₃-like phase and also tetrahedral coordinated \([\text{Mo(VI)}\text{O}_4]^2\), which is particularly evident at the pixel located at the particle rim. The average XANES from spectra 1 through 4 is shown in Figure 4 (#5), where the main absorption peak is broadened and the contribution of tetrahedral coordinated Mo(VI) reduced to a mere shoulder, which is imaginably a trend towards the sample I spectrum in Figure 1 measured with a ~25 µm beam. We interpret the two observations that 1) both the smallest particles and the rim of a large particle exhibit the largest tetrahedral coordinated Mo(VI) signal and that 2) the particles studied have Mo in its hexavalent state as suggesting charge compensation at the Mo-particles surface occurs through reduction of the number of O atoms coordinated to Mo (tetrahedral surface versus octahedral bulk coordination).

4. Conclusions
These results show that none of the substrates or temper programs used produces satisfactory SNF ε-particle models. While nano- and micron-sized Mo particles are formed in the UO₂/Mo TF systems studied, these particles are either not embedded in the desired UO₂ substrate or are present in oxide form, instead of the desired metallic state. The oxide particles, however, do offer the opportunity to study changes in surface structures in response to stress/charge as a function of particle size. It may be possible to prepare Mo(0)-particles embedded in UO₂ on Si substrates through tempering at temperatures below 1000°C, but above 600 °C.

Acknowledgments
We thank PETRAIII, HASYLAB and ANKA for the granted beamtime.

References
[1] Kleykamp, H 1985 J Nucl Materials 131(2-3) 221
[2] Bruno J, Ewing R 2006 Elements 2 34.
[3] Johnson LH, et al. Spent Fuel. In Radioactive waste forms for the future, Lutze, W (Ed) 1988, pp 635-698
[4] Shoesmith DW 2000 J Nucl Materials 282 1
[5] Denecke MA, et al. 2012 MRS Spring Meeting Proceedings 1444, mrrs12-1444-y01-05doi:10.1557/opl.2012.1159
[6] Rothe J, et al. 2012 Rev Sci Instrum 83 043105
[7] Schroer CG, et al. 2010 Nucl Instrum Methods A616 93
[8] Denecke MA Proceedings of the OECD-NEA Workshop on Speciation, Techniques and Facilities for Radioactive Materials at Synchrotron Light Sources, Grenoble, France, 4-6 October 1999, pp. 135-41.
[9] Ressler T, Jentoft RE, Weinhold J, Günter MM, Timpe O 2000 J Phys Chem B104, 6360