Kinetic Models for the in Situ Reaction between Cu-Ti Melt and Graphite

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Abstract: The in situ reaction method for preparing metal matrix composites has the advantages of a simple process, good combination of the reinforcing phase and matrix, etc. Based on the mechanism of forming TiCₙ particles via the dissolution reaction of solid carbon (C) particles in Cu-Ti melt, the kinetic models for C particle dissolution reaction were established. The kinetic models of the dissolution reaction of spherical, cylindrical, and flat C source particles in Cu-Ti melt were deduced, and the expressions of the time for the complete reaction of C source particles of different sizes were obtained. The mathematical relationship between the degree of reaction of C source and the reaction time was deduced by introducing the shape factor. By immersing a cylindrical C rod in a Cu-Ti melt and placing it in a super-gravity field for the dissolution reaction, it was found that the super-gravity field could cause the precipitated TiCₓ particles to aggregate toward the upper part of the sample under the action of buoyancy. Therefore, the consuming rate of the C rod was significantly accelerated. Based on the flat C source reaction kinetic model, the relationship between the floating speed of TiCₓ particles in the Cu-Ti melt and the centrifugal velocity (or the coefficient of super-gravity G) was derived. It was proven that, when the centrifugal velocity exceeded a critical value, the super-gravity field could completely avoid the accumulation behavior of TiCₓ particles on the surface of the C source, thereby speeding up the formation reaction of TiCₓ. The goal of this study is to better understand and evaluate the generating process of TiCₓ particles, thus finding possible methods to increase the reaction efficiency.

Keywords: in situ reaction; kinetics; metal matrix composites; TiC; Cu

1. Introduction

TiCₓ powder has a high melting point, corrosion resistance, and hardness (micro hardness up to 3000 kg/mm²), which is just lower than diamond. It is a high-quality basic material for cutting tools, wear-resistant parts, high-temperature alloys, etc., and it has important application value in the field of high-mechanical-property material preparation [1]. However, its synthesis conditions are very harsh [2,3], as the carbothermal reduction method with TiO₂ powder, which is commonly used in the industry, requires a high synthesis temperature (1700–2100 °C) and long reaction time of 10–24 h. Due to the limitation on mass transfer conditions of solid-solid reactions, the products often contain unreacted C and Ti. Therefore, the synthesized TiCₓ powder has a wide range of particle size distribution, which demands an additional process of ball milling. The direct carbonization method [4] uses Ti powder and C powder to generate TiCₓ. The reaction takes 5–20 h, and the reaction process is difficult to control. The reactants are severely prone to agglomeration, and fine grinding is required to prepare fine-sized TiCₓ powder. In order to obtain a purer product, the fine powder after ball milling needs to be chemically purified. In addition, due to the high price of Ti metal powder, the cost of synthesizing TiCₓ is also high with this method. The self-propagating high-temperature
synthesis (SHS) [5–7] method uses the reaction heat of the TiC$_x$ formation reaction to keep the reaction going. This synthesis method requires high-purity and fine Ti powder as raw materials [8,9], and the yield is limited. The reaction of the ball milling method [10–12] uses Ti powder and C powder to generate TiC$_x$ via their combination reaction under the action of mechanical energy during the ball milling process. Because the reaction process is very slow, this method is usually used to produce nanoscale TiC$_x$ powder on a small scale. Other methods for preparing TiC$_x$ powder, such as the chemical vapor deposition method and the microwave synthesis method, all have some defects, which cannot meet the needs of industrial production.

In view of the above problems, the authors herein propose a method for preparing TiC$_x$ powder via the solid-liquid reaction between Cu-Ti melt and dissolvable solid C coupled with super-gravity separation. The solid-liquid in situ reaction between solid C and Cu-Ti melt was firstly used to generate TiC$_x$ particles [13–15]. After super-gravity filtration [16] and acid leaching purification treatments, the qualified TiC$_x$ powder products could be obtained. The feasibility of this method was proven by experiments. Compared with the most commonly used TiO$_2$ carbothermal reduction method for producing TiC$_x$ particles, this method has several advantages. Firstly, the use of the Cu-Ti binary alloy melt reduces the reaction temperature from above 1700 °C to about 1000 °C. The decrease in reaction temperature helps to reduce production costs and equipment burden. Secondly, TiC$_x$ is synthesized via solid-liquid reaction. Compared with the solid-solid reaction, it can effectively improve mass transfer conditions and, thus, increase reaction speed. The reaction time can be shortened from more than 10 h to 2 h or even shorter. Thirdly, the particle size of TiC$_x$ particles precipitated in the Cu-Ti melt is not affected by the particle size of raw Ti material; thus, this method can use low-purity bulk sponge titanium as the raw material to get rid of the limitation of using high-purity fine-grained metal Ti powder or TiO$_2$ powder, which can help greatly reduce the cost of raw materials. Lastly, because the TiC$_x$ particles precipitated in the melt have good dispersibility, and the operating temperature is relatively low, there is no particle bonding problem.

In most cases, TiC$_x$ powder is compounded into a metal matrix as a particulate reinforcing material to prepare MMCs (metal matrix composites). The preparation of this composite material can be divided into two ways. The first approach is to prepare TiC$_x$ powder, and then prepare MMCs by powder metallurgy [13,17–20], or to mix TiC$_x$ powder into a metal solution and cast it [21]. The second method is to combine the preparation of TiC$_x$ powder and the composite process. The one-step preparation of composite materials can be achieved via in situ reaction [22,23], which can help obtain a high interface bonding strength between the TiC$_x$ particles and the matrix metal [24]. Compared with ex situ processing techniques, synthesis approaches based on the in situ formation of particles in metal matrices offer more possibilities to control the particle size and particle distribution pattern in the composites [25,26].

We prepared a composite with the Cu-11.1wt.% Ti alloy as a matrix material, and the TiC$_x$ particles were prepared as the reinforcement phase via a previously described second method [27]. For the two methods of preparing TiC$_x$ particle-enhanced MMCs, the analysis of their reaction mechanism and the description of their kinetics are of great significance for better regulating the reaction. Therefore, we combined the phenomena in the experimental study to perform a kinetic modeling analysis of the reaction process.

A lot of research work on liquid-solid separation at high temperatures with super-gravity technology was carried out in our previous work [28–34], and we found that solid particles would rapidly float or settle in high-temperature melt due to their density difference. Previous work mainly focused on the super-gravity separation process, and little research was done on the effects of super-gravity fields on chemical reaction kinetics. Therefore, in this experiment, an attempt was made to introduce a super-gravity field into the in situ reaction of generating TiC$_x$. It was found that the super-gravity field could accelerate the floating speed of TiC$_x$ particles in the Cu-Ti melt, thereby obtaining the effect of accelerating the C source dissolution reaction speed. In this paper, the theoretical analysis and a mathematical model for TiC$_x$ particle floating behavior in a super-gravity field and its influence on TiC$_x$ formation were established.
The mathematical model presented in this article helps to better understand and evaluate the process of the solid-liquid in situ reaction of generating TiC\(_x\) particles, so as to provide ideas for seeking measures to accelerate the dissolution reaction of the C source and promote the complete reaction of the C source. At the same time, it can provide a theoretical basis for future numerical simulation studies of this process.

2. Model Descriptions

2.1. Basic Assumptions

The modeling work in this study is based on the experimental phenomenon of the solid-liquid in situ reaction between graphite carbon and Cu-Ti alloy. The reaction of solid C source particles in the Cu-Ti melt is [Ti] + xC(s) = TiC\(_x\)(s). This reaction can be regarded as the “dissolution” process of the solid phase in the liquid phase. It was found in the experiment that, during the reaction, a TiC\(_x\) transition layer forms on the surface of the solid C source particles, which was also found in the study of Dudina et al. [26] when investigating the solid-state interactions in the Ti\(_{25}\)Cu\(_{75}\) + C system with different sources of carbon. For a solid C particle wrapped by solid TiC\(_x\), if the reaction interface is between the liquid alloy and the TiC\(_x\) layer and the reaction is determined by the diffusion of C, then the volume of the solid C should decrease gradually, and there will be a gap between the solid particle and the layered TiC\(_x\). The above phenomenon was not detected in our previous study [27]; instead, the solid C particle combined closely with the solid TiC\(_x\) layer. Thus, we deduced that there is a high possibility that the interface where the chemical reaction occurs should be at the interface between the C source and the layered TiC\(_x\). Ti atoms migrate to the reaction interface within the layered TiC\(_x\) via solid-state diffusion [35]. This point was also followed by some other researchers [36], and we think it needs to be further studied to give clear evidence in later work. As the TiC\(_x\) generation reaction proceeds, the diameter of the C source particles decreases steadily, the reaction interface advances toward the core of the particles, and the layered TiC\(_x\) advances toward the C source. Simultaneously, the layered TiC\(_x\) releases fine TiC\(_x\) particles to the Cu-Ti melt, and these TiC\(_x\) particles gradually disperse into the Cu-Ti melt.

The reaction kinetics of the solid C source in the Cu-Ti melt can be divided into the following three steps: (1) Ti atoms diffuse in the Cu-Ti melt to the surface of solid particles participating in the reaction under the effect of a concentration gradient; (2) Ti atoms diffuse through the solid TiC\(_x\) product layer; (3) an in situ chemical reaction occurs at the interface between the solid C source and the layered TiC\(_x\). In order to facilitate the mathematical modeling and analysis of previous processes, the following assumptions were made in this study: (1) the diffusion of Ti atoms in the Cu-Ti melt and in the layered TiC\(_x\) conform to Fick’s first law; (2) the TiC\(_x\) particles dispersed around the solid particles will not have a significant effect on the diffusion of Ti atoms in the Cu-Ti melt; (3) the speed of the reaction interface advances to the interior of the C source particles, and the depletion rates of the layered TiC\(_x\) during the reaction are the same, i.e., the thickness of the layered TiC\(_x\) remains constant. Based on the above theories and assumptions, the reaction kinetics of spherical, cylindrical, and flat C sources were modeled in this paper. The detailed content is described below.

2.2. Spherical C Source

In order to simplify the reasoning process, in a previous study [27], the concentration of Ti atoms at infinity relative to the carbon source particles was taken as the original concentration of Ti in the Cu-Ti melt (C\(_{Ti,b}\)). In fact, there should be a concentration transition layer of Ti atoms near the C source. Its thickness is likely to be limited relative to the diameter of the C source particles. Therefore, the thickness of this concentration transition zone is defined as “h”, as shown in Figure 1. The in situ reaction kinetic of spherical C source particles in the Cu-Ti melt was re-derived in this paper. The specific process is described below.
Figure 1. The schematic diagram of the reaction mechanism of spherical solid C source particles in Cu-Ti melt.

For step (1) defined in Section 2.1 and shown in Figure 1,

\[
\dot{n}_{Ti,1} = 4\pi r^2 D_{Ti,1} \frac{dC_{Ti}}{dr},
\]

where \( \dot{n}_{Ti,1} \) is the diffusion flux of Ti atoms from the distant Cu-Ti melt to the layered TiCx surface (mol/s), \( r \) is the radius of the C source particles (m), \( D_{Ti,1} \) is the diffusion coefficient of Ti atoms in the Cu-Ti melt (m²/s), and \( C_{Ti} \) is the concentration of Ti atoms (mol/m³). If this process is considered as steady-state diffusion, then the variables are separated and integrated in the above formula as follows:

\[
\dot{n}_{Ti,1} \int_{r+h}^{r+d+h} \frac{dr}{r^2} = 4\pi D_{Ti,1} \int_{C_{Ti,b}}^{C_{Ti,s}} dC_{Ti}
\]

That is,

\[
\dot{n}_{Ti,1} = 4\pi \frac{(r+d) \cdot (r+d+h)}{h} D_{Ti,1} (C_{Ti,b} - C_{Ti,s}),
\]

where \( C_{Ti,b} \) and \( C_{Ti,s} \) are the concentration of Ti atoms in the Cu-Ti melt and on the layered TiCx surface, respectively (mol/m³), \( d \) is the thickness of the layered TiCx (m), and \( h \) is the concentration transition layer thickness of Ti atoms in the Cu-Ti melt (m). Equation (3) is an expression for the diffusion of Ti atoms from the Cu-Ti melt to the surface of the particles enclosed by the TiCx product layer. For step (2), the solid-state diffusion of Ti atoms through the layered TiCx is

\[
\dot{n}_{Ti,2} = 4\pi r^2 D_{Ti,2} \frac{dC_{Ti}}{dr},
\]

where \( \dot{n}_{Ti,2} \) is the diffusion flux of Ti atoms through the layered TiCx product layer (mol/s), and \( D_{Ti,2} \) is the diffusion coefficient of Ti atoms in the layered TiCx product layer (m²/s). If the process is considered steady-state diffusion, then we have

\[
\dot{n}_{Ti,2} \int_{r+d}^{r+d+h} \frac{dr}{r^2} = 4\pi D_{Ti,2} \int_{C_{Ti,b}}^{C_{Ti,s}} dC_{Ti}
\]

That is,
\[ \dot{n}_{Ti,2} = \frac{4\pi r(r+d)D_{Ti,2}(C_{Ti,s} - C_{Ti,r})}{d}, \]  

(6)

where \( C_{Ti,r} \) is the Ti atom concentration on the surface of the solid C source particle (mol/m\(^3\)). For step (3), in the formation reaction of TiC\(_x\) at the interface between the solid C source and the layered TiC\(_x\), if this reaction is regarded as a first-order irreversible reaction, then we have

\[ \dot{n}_{Ti,r} = 4\pi r^2 k_i C_{Ti,r}, \]  

(7)

where \( \dot{n}_{Ti,r} \) is the reaction rate of TiC\(_x\) generated at the surface of C source (mol/s), and \( k_i \) is the chemical reaction rate constant. Assuming that the previous steps are in a steady state, the speed of each step is equal, i.e.,

\[ \dot{n}_{Ti,i} = \dot{n}_{Ti,2} = \dot{n}_{Ti,r} = r_{Ti}, \]  

(8)

where \( r_{Ti} \) is the Ti atom consumption rate of the overall process (mol/s). Substituting Equation (3) into Equation (6) to remove \( C_{Ti,s} \) and then bringing the obtained expression of \( C_{Ti,r} \) into Equation (7), we can get

\[ r_{Ti} = \frac{4\pi r^2 k_i C_{Ti,b}}{1 + \frac{r^2 k_i h}{(r+d)(r+d+h)D_{Ti,1}} + \frac{r k_i d}{(r+d)D_{Ti,2}}}, \]  

(9)

assuming that \( r_c \) is the consumption rate of solid C source during the reaction (mol/s). Then, we get

\[ r_{Ti} = \frac{1}{x} r_c = -\frac{\rho_c}{x} \frac{d}{dr} \left( \frac{4\pi r^3}{3} \right) = -\frac{\rho_c}{x} 4\pi r^2 \frac{dr}{dt}, \]  

(10)

where \( \rho_c \) is the molar density of the solid C source (mol/m\(^3\)), and \( x \) is the stoichiometry of the product TiC\(_x\). From Equations (9) and (10), we get

\[ \frac{dr}{dt} = \frac{x k_i C_{Ti,b}}{\rho_c} \left[ 1 + \frac{r^2 k_i h}{(r+d)(r+d+h)D_{Ti,1}} + \frac{r k_i d}{(r+d)D_{Ti,2}} \right] \]  

(11)

Because both \( D_{Ti,1} \) and \( D_{Ti,2} \) are independent of the radius \( r \) of the solid C source particle, the variables of above equation are separated and integrated as

\[ -\int_{r_0}^{r} \rho_c \left[ 1 + \frac{r^2 k_i h}{(r+d)(r+d+h)D_{Ti,1}} + \frac{r k_i d}{(r+d)D_{Ti,2}} \right] dr = \int_{0}^{t} x k_i C_{Ti,b} dt' \]  

(12)

where \( r_0 \) is the initial radius of the solid C source particles (m), and \( t \) is the reaction time (s). The mathematical analytical solution of the reaction time \( t \) can be obtained from the above equation as follows:

\[
\frac{-\rho_c}{x k_i C_{Ti,b}} \left[ r - \frac{k_i h}{D_{Ti,1}} \left[ \ln(r+d+h) + \frac{2d - d^2}{h} \ln \left( \frac{r+d+h}{r+d} \right) \right] - \frac{k_i d}{D_{Ti,2}} \left[ r - d \ln(r+d) \right] \right]_{r_0}^{r}
\]  

(13)

When \( r = 0 \), the time required for the solid C source to be completely consumed is
Through the above model, the kinetic process of the reaction of spherical solid C source particles in Cu-Ti melt under certain conditions can be evaluated, and the time for the spherical C source particles to be consumed thoroughly can be calculated.

2.3. Cylindrical C Source

The reaction process of cylindrical solid C source particles in Cu-Ti melt is shown in Figure 2.

![Figure 2](image)

**Figure 2.** The schematic diagram of reaction mechanism of cylindrical solid C source particles in Cu-Ti melt.

For step (1),

\[
\dot{n}_{T_{i,1}} = 2\pi r D_{T_{i,1}} \frac{dC_{Ti}}{dr} \tag{15}
\]

When step (1) is in a steady state, we can get the equation below after separating and integrating the variables in Equation (15).

\[
\dot{n}_{T_{i,1}} \int_{r+d+h}^{r+d} \frac{dr}{r} = 2\pi r D_{T_{i,1}} \int_{C_{Ti,b}}^{C_{Ti,a}} dC_{Ti} \tag{16}
\]

That is,

\[
\dot{n}_{T_{i,1}} \left[ \ln(r+d) - \ln(r+d+h) \right] = 2\pi r D_{T_{i,1}} (C_{Ti,a} - C_{Ti,b}) \tag{17}
\]

\[
\dot{n}_{T_{i,1}} = \frac{2\pi r D_{T_{i,1}} (C_{Ti,a} - C_{Ti,b})}{\ln \left( \frac{r+d}{r+d+h} \right)} \tag{18}
\]

For step (2),

\[
\dot{n}_{T_{i,2}} = 2\pi r D_{T_{i,2}} \frac{dC_{Ti}}{dr} \tag{19}
\]

After separating and integrating of variables in Equation (19), we get
\[
\dot{n}_{Ti,2} \int_{r+d}^{r} \frac{dr}{r} = 2\pi Id_{Ti,2} \int_{C_{Ti}}^{C_{Ti,2}} r \, dC_{Ti},
\]

(20)

\[
\dot{n}_{Ti,2} [\ln r - \ln(r + d)] = 2\pi Id_{Ti,2} (C_{Ti,f} - C_{Ti,b})
\]

(21)

Then, we have

\[
\dot{n}_{Ti,2} = \frac{2\pi Id_{Ti,2} (C_{Ti,f} - C_{Ti,b})}{\ln \left(\frac{r}{r + d}\right)}
\]

(22)

Taking step (3) as a first-order irreversible reaction, ignoring the area of the upper and lower sides of the cylinder (that is, “l” is much larger than “r”), then we can get Equation (23).

\[
\dot{n}_{Ti,2} = 2\pi r k_{i} C_{Ti,f}
\]

(23)

When the in situ reaction is in a steady state, i.e., the rates of steps (1–3) are equal, then the following equation can be obtained from Equations (8), (18), (22), and (23):

\[
r_{Ti} = \frac{2\pi r k_{i} C_{Ti,b}}{1 - \frac{r k_{i}}{D_{Ti,1}} \ln \left(\frac{r + d}{r + d + h}\right) - \frac{r k_{i}}{D_{Ti,2}} \ln \left(\frac{r}{r + d}\right)}
\]

(24)

This is possible, because

\[
r_{Ti} = \frac{1}{x} r_{c} = -\frac{D_{c}}{x} \times \frac{d}{dt} \left(\frac{\pi r^{2} l}{l}ight) = -\frac{2\pi r_{c} r l}{x} \times \frac{dr}{dt}
\]

(25)

From Equations (24) and (25), we get

\[
\frac{dx}{dt} = \frac{xk_{i} C_{Ti,b}}{\rho_{c} \left[1 - \frac{r k_{i}}{D_{Ti,1}} \ln \left(\frac{r + d}{r + d + h}\right) - \frac{r k_{i}}{D_{Ti,2}} \ln \left(\frac{r}{r + d}\right)\right]}
\]

(26)

Upon separating and integrating variables in Equation (26), we can get

\[
-\int_{x}^{x(t)} \rho_{c} \left[1 - \frac{r k_{i}}{D_{Ti,1}} \ln \left(\frac{r + d}{r + d + h}\right) - \frac{r k_{i}}{D_{Ti,2}} \ln \left(\frac{r}{r + d}\right)\right] \, dx = \int_{0}^{t} \frac{xk_{i} C_{Ti,b}}{x} \, dt,
\]

(27)

from which we have

\[
-\frac{\rho_{c}}{xk_{i} C_{Ti,b}} \left[r + \frac{k_{i}}{D_{Ti,1}} \left[\frac{1}{2} r^{2} - (d + h)^{2}\right] \ln (r + d + h) - \frac{1}{4} r^{2} + \frac{1}{2} r (d + h)\right] =
\]

(28)

\[
-\frac{k_{i}}{D_{Ti,2}} \left(\frac{1}{2} r^{2} \ln r - \frac{1}{2} r^{2}\right) + \left(\frac{k_{i}}{D_{Ti,1}} - \frac{k_{i}}{D_{Ti,2}}\right) \left[\frac{1}{2} (r^{2} - d^{2}) \ln (r + d) - \frac{1}{4} r^{2} + \frac{1}{2} r d\right]
\]

\[
t = \frac{\rho_{c}}{xk_{i} C_{Ti,b}}
\]

When \( r = 0 \), the time required for complete reaction is
The model was used to describe the kinetics of the reaction of cylindrical solid C source particles in Cu-Ti melt under certain conditions, and to calculate the time for the cylindrical C source particles to achieve a complete reaction.

2.4. Flat C Source

The reaction process mechanism of flat solid C source particles in Cu-Ti melt is shown in Figure 3.

For step (1), we have

\[ \dot{n}_{T_1,1} = abD_{T_1,1} \frac{dC_{Ti}}{dr}, \]  

where \( a \) and \( b \) are the length and width of the flat solid C source. When step (1) achieves steady-state diffusion, the variables in Equation (30) are separated and integrated analogously to Equation (2) or (16); afterward, we can get

\[ \dot{n}_{T_1,1} = abD_{T_1,1} \frac{(C_{Ti,b} - C_{Ti,s})}{h} \]  

Similarly, for step (2), we have

\[ \dot{n}_{T_1,2} = abD_{T_1,2} \frac{(C_{Ti,s} - C_{Ti,r})}{d} \]  

If the TiC\(_x\) formation reaction is a first-order irreversible reaction, then

\[ \dot{n}_{Ti,r} = abk\_r C_{Ti,r} \]
When the in situ reaction to generate TiC\textsubscript{x} is in a steady state, that is, the rates of steps (1–3) are equal, the following can be obtained from Equations (8), (31), (32), (33):

\[ r_{Ti} = \frac{abk_r C_{Ti,b}}{1 + \frac{dk_r}{D_{Ti,2}} + \frac{hk_r}{D_{Ti,1}}} \]  

(34)

According to the stoichiometric relationship of Ti and C elements in the generated TiC\textsubscript{x},

\[ r_{Ti} = \frac{1}{x} r^c = -\frac{\rho_c}{x} \frac{d}{dt}(abr) = -\frac{\rho_c ab}{x} \frac{dr}{dt} \]  

(35)

From Equations (34) and (35), we can get

\[ -\frac{dr}{dt} = \frac{x r_{Ti}}{ab \rho_c} = \frac{x k_r C_{Ti,b}}{\rho_c \left(1 + \frac{dk_r}{D_{Ti,2}} + \frac{hk_r}{D_{Ti,1}}\right)} \]  

(36)

Upon separating and integrating the variables in the above equation, we have

\[ \int_{r_0}^{r} \rho_c \left(1 + \frac{dk_r}{D_{Ti,2}} + \frac{hk_r}{D_{Ti,1}}\right) dr = \int_0^t dt \]  

(37)

Then,

\[ (r_0 - r) \rho_c \left(1 + \frac{dk_r}{D_{Ti,2}} + \frac{hk_r}{D_{Ti,1}}\right) = \frac{r_0 \rho_c \left(1 + \frac{dk_r}{D_{Ti,2}} + \frac{hk_r}{D_{Ti,1}}\right)}{x k_r C_{Ti,b}} \]  

(38)

When \( r = 0 \), that is, for the total reaction time,

\[ t_{total} = \frac{r_0 \rho_c \left(1 + \frac{dk_r}{D_{Ti,2}} + \frac{hk_r}{D_{Ti,1}}\right)}{x k_r C_{Ti,b}} \]  

(39)

If \( r_{Til} \) is the rate of reaction defined by the consuming molar rate of Ti per unit time per unit surface area of the C source, then the area term in the expression of step (1) can be omitted. The relationship between \( r_{Til} \) and \( r_{Ti} \) becomes

\[ r_{Til} = \frac{r_{Ti}}{S} \]  

(40)

where \( S \) for a spherical C source is \( 4\pi r^2 \), \( S \) for a cylindrical C source is \( 2\pi rl \), and \( S \) for a flat C source is \( ab \). When the chemical reaction is the limiting link (i.e., \( C_{Ti,r} = C_{Ti,b} \)) in the condition of, e.g., the mechanical agitation, the diffusion resistance of Ti atoms can be ignored. At this time, for the three types of C sources,

\[ r_{Til} = k_r C_{Ti,r} = -\frac{\rho_c}{x} \frac{dr}{dt} \]  

(41)

After separating variables of Equation (41),
For all the three types of C sources, we have
\[
\frac{1}{F_p V_p} \frac{A_p}{r_0} = F_p \rho_p \frac{C_{Ti} \, d r}{t \, d t} = -\frac{\rho_c}{F_p V_p} C_{Ti} \, d t
\]

where \( A_p \) is the original surface area of the C source, \( F_p \) is the shape factor of the C source (sphere \( F_p = 3 \), cylinder \( F_p = 2 \), flat plate \( F_p = 1 \)), and \( V_p \) is the original volume of the C source \[37\].

Upon dividing by \( r_0 \) at both sides of Equation (42), we get
\[
\frac{1}{F_p V_p} \frac{A_p}{r_0} \frac{C_{Ti} \, d r}{t \, d t} = -\frac{\rho_c}{F_p V_p} \frac{C_{Ti} \, d r}{r_0}
\]

For large plates, \( \frac{F_p V_p}{A_p} \) is half of their original thickness, and, for spheres or cylinders, \( \frac{F_p V_p}{A_p} \) is their original radius. Integrating Equation (44) gives
\[
\int_0^r \frac{1}{F_p V_p} \frac{A_p}{r_0} \frac{C_{Ti} \, d r}{t \, d t} = -\int_0^r \frac{\rho_c}{F_p V_p} \frac{d t}{r_0}
\]

That is,
\[
\frac{1}{F_p V_p} \frac{A_p}{r_0} \frac{C_{Ti} \, t}{C_{Ti} \, t} = 1 - \frac{r}{r_0}
\]

Suppose that \( t^* = \frac{1}{F_p V_p} \frac{A_p}{r_0} C_{Ti} \), \( \xi = \frac{A_p}{r_0} r = \frac{r}{r_0} \), then Equation (46) can be transformed into the dimensionless form as follows:
\[
\frac{d \xi}{d t^*} = -1
\]

\( \xi=1 \) and \( t^*=0 \) are the initial conditions of Equation (47), from which we can get the relationships below after integrating Equation (47).
\[
\begin{align*}
\xi & = 1 - t^* \quad \text{when} \quad t^* \leq 1 \\
\xi & = 0 \quad \text{when} \quad t^* \geq 1
\end{align*}
\]

The above reasoning process assumes that the C source particles all maintain their original particle shape during the reaction. The degree of reaction of the C source \( X \) can be expressed as
\[
X = 1 - \frac{\xi}{F_p}
\]

From Equations (48) and (49), the relationship between the degree of reaction of the C source and time can be obtained as follows:
\[
t^* = \frac{1}{F_p} \frac{A_p}{r_0} \frac{C_{Ti} \, t}{C_{Ti} \, t} = 1 - (1 - X)^\frac{1}{F_p}
\]

Thus, the time required for the C source to be completely consumed can be obtained, that is, when \( X = 1 \), we have
We then define

$$t_{x_{i=1}} = \frac{\rho_c F_p V}{xk_i C_{Ti} t A_{pp}}$$  \hspace{1cm} (51)$$

For large C source particles, the shape factor can be obtained by visual observation. For small C source particles that cannot be easily observed directly, or when their shape factor cannot be simply determined, we can set different $F_p$ values and bring them into Equation (52). The numerical relationship between $g_{F_p}(X)$ and the reaction time $t$ can be obtained by experiment; when they show a good linear relationship with certain $F_p$ values, then the value of $F_p$ can be used as the shape factor of the C source. Thereafter, when the effect of Ti atom diffusion resistance is excluded in the experiment, the reaction order, activation energy, and pre-exponential factor of the reaction can be inferred through experiments under different Ti atom concentrations and different temperatures.

The above models are adjustable for different kinds of C in terms of the reaction kinetics. When it comes to some procedures like the peeling of C, the dispersion of TiC particles, etc., further experimental investigations are needed to provide clues for the refining of this model.

2.5. The Reaction Behavior in Super-Gravity Field

In this paper, a super-gravity field was introduced in the reaction of solid C particles in the Cu-Ti melt. The effect of the super-gravity field on the kinetics of the in situ reaction of solid C in the Cu-Ti alloy melt was experimentally verified. The sketch map of the super-gravity centrifuge used in the experiment is shown in Figure 4. In the experiment, 32 g of copper powder (25–75 μm) was mixed with 8 g of titanium powder (165–665 μm) in a graphite crucible as shown in Figure 5a, and then the graphite crucible was sealed with the graphite lid equipped with a graphite rod, before placing the set-up in the resistance furnace of the centrifugal device (Figure 4). The temperature was raised to 1250 °C at a heating rate of 10 °C/min, and the centrifuge was turned on after holding for 30 min. The centrifuge was maintained at a speed of $N = 1892$ rpm (gravity coefficient $G = 1000 \times g$) for 30 min. After the test, the centrifuge and heating program were turned off, and the samples were taken out after cooling to room temperature. Comparative samples were prepared under normal gravity conditions (the remaining experimental conditions except for the gravity coefficient were the same as above). The longitudinal profile of the two samples from the center position is shown in Figure 5c,d.

![Figure 4. Sketch map of the centrifugal apparatus: 1 counterweight, 2 centrifugal axis, 3 conductive slipping, 4 thermocouple, 5 insulating layer, 6 temperature controller, 7 TiC particles, 8 Cu-Ti melt, 9 graphite crucible.](image-url)
Figure 5. Comparison of the size change in solid C rods dissolved by Cu-Ti melt under normal gravity and super-gravity conditions (a) image of the designed graphite crucible; (b) the image of the sample after heating treatment; (c) the cross-section morphology of the sample after super-gravity treatment with $G = 1000$; (d) the cross-section morphology of the sample after heating in normal gravity.

The initial diameter of the C rod used in this test was 6 mm. As can be seen from Figure 5d, under normal gravity conditions, the diameter of the C rod did not change significantly, i.e., the diameters of the top and bottom of the C rod were 6 mm. Under the condition of a super-gravity field, the difference between the diameters of the top and the bottom was about 0.9 mm (Figure 5c). The main reason for this difference was that the reaction between the Cu-Ti melt and solid C was strengthened under the condition of super-gravity ($G = 1000$). The kinetic model for the flat C source was used as an example to analyze the reaction process under the condition of super-gravity.

If the TiCx particles and the molten Cu-Ti alloy are considered as hard spherical particles and a viscous liquid, respectively, then the moving behavior of the TiCx particles in the super-gravity field can be simulated with Stokes law [34]. There are two major forces acting on the particles: the centrifugal force and the viscous drag force, which are in opposite directions. The force balance of the particles can be expressed as follows [38,39]:

$$m_p \frac{d^2 R}{dt^2} = \rho_{\text{TiCx}} - \rho_{\text{Cu-Ti}} \frac{4}{3} \pi \left( \frac{D_p}{2} \right)^3 G g - 3 \pi \eta D_p \frac{d^2 R}{dt^2},$$

(53)

where $R$ is the distance between the particles and the rotating axis of the centrifugal machine (m), $dR/dt$ is the moving velocity (m/s), $d^2 R/dt^2$ is the acceleration speed (m/s²), $\rho_{\text{TiCx}}$ is the density of TiCx particles (4930 kg/m³), $\rho_{\text{Cu-Ti}}$ is the density of molten Cu-Ti alloy (kg/m³), $D_p$ is the diameter of TiCx particles (m), $G$ is the super-gravity coefficient (dimensionless), $g$ is the gravitational acceleration (9.8 m/s²), and $\eta$ is the viscosity of liquid (Pa·s). The left term in Equation (53) represents the buoyant force [39]. The first term on the right side is the centrifugal force and the second term on the right side is the viscous drag force [39]. The moving direction of the particles is determined by the relative values of densities. In this study, $\rho_{\text{TiCx}} < \rho_{\text{Cu-Ti}}$; thus, the particles move toward the opposite direction of the centrifugal force. The terms $\rho_{\text{TiCx}}, \rho_{\text{Cu-Ti}}, \eta,$ and $D_p$ in Equation (53) were assumed to be time-independent; then, Equation (53) could be solved under the initial condition of $d^2 R/dt^2 = 0$ at $t = 0$. The solution is
where $\omega$ is the rotational angular velocity (rad/s). As the moving direction of the TiC$_x$ particles is opposite to the direction of super-gravity, a negative sign is added to Equation (54). It is known that $\rho_{\text{TiC}_x} = 4930 \text{ kg/m}^3$, $\rho_{\text{Cu}} = 8960 \text{ kg/m}^3$, and $\rho_{\text{Ti}} = 4500 \text{ kg/m}^3$. In this experiment, $D_p = 3 \times 10^{-6} \text{ m}$ and $R = 0.25 \text{ m}$. The rotation speed at $G = 1000$ was 1892 rpm ($\omega = 198 \text{ rad/s}$). Suppose that the viscosity of molten Ti is $\eta_{\text{Ti}} \approx 5.2 \times 10^{-3} \text{ Pa.s}$ [40,41], and the viscosity of molten Cu is $\eta_{\text{Cu}} \approx 4 \times 10^{-3} \text{ Pa.s}$ [42]. To simplify the calculation, suppose that $\rho_{\text{Cu-Ti}} = f_{\text{Cu}} \times \rho_{\text{Cu}} + f_{\text{Ti}} \times \rho_{\text{Ti}}$ and $\eta_{\text{Cu-Ti}} = f_{\text{Cu}} \times \eta_{\text{Cu}} + f_{\text{Ti}} \times \eta_{\text{Ti}}$, where $f_{\text{Cu}}$ and $f_{\text{Ti}}$ are the mass fraction of Cu and Ti in the Cu-Ti alloy. Then, the values of $\Delta \rho (|\rho_{\text{TiC}_x} - \rho_{\text{Cu-Ti}}|)$ and $\eta_{\text{Cu-Ti}}$ with different Cu-Ti proportions can be obtained as listed in Table 1.

Table 1. The values of $\rho_{\text{Cu-Ti}}$, $\Delta \rho$, and $\eta_{\text{Cu-Ti}}$ with different Cu-Ti proportions (e.g., “Cu-20 wt.% Ti” means the weight percentage of Ti in the Cu-Ti alloy is 20 wt.%).

| Cu-20 wt.% Ti | Cu-40 wt.% Ti | Cu-60 wt.% Ti | Cu-80 wt.% Ti |
|----------------|----------------|----------------|----------------|
| $\rho_{\text{Cu-Ti}}$ (kg/m$^3$) | 8068 | 7176 | 6284 | 5392 |
| $\Delta \rho$ (kg/m$^3$) | 3138 | 2246 | 1354 | 462 |
| $\eta_{\text{Cu-Ti}}$ ($\times 10^{-3}$ Pa·s) | 4.24 | 4.48 | 4.72 | 4.96 |

Then, according to Equation (54) and the data in Table 1, the moving velocity of TiC$_x$ particles in the Cu-Ti melt under the super-gravity field can be calculated as shown in Figure 6. It can be seen that the moving velocity of TiC$_x$ particles increases with the decrease in Ti content in the Cu-Ti melt, the increase in rotation velocity of the centrifugal apparatus, and the increase in particle diameter.

Figure 6. The correlation between the TiC$_x$ particle floating speed in the Cu-Ti melt and the centrifugal velocity (super gravity coefficient $G$).

The viscosity of the melt increases with the increased number of TiC$_x$ particles in the melt, which can be expressed as follows [43]:
\[ \eta = \frac{\eta_0}{\left(1 - \frac{\epsilon}{\epsilon_{\text{max}}}\right)^{2.5}} \]  
(55)

where \(\eta_0\) is the viscosity of the molten metal without particles (Pa·s), \(\epsilon\) is the particle volume fraction of TiC\(_x\), and \(\epsilon_{\text{max}}\) is the maximum packing fraction. The volume of TiC\(_x\) produced per unit time is

\[ V_{\text{TiC}_x} = -\frac{ab \rho_c}{\rho_{\text{TiC}_x}} \frac{dr}{dt} \frac{48 + 12x}{12x} \]  
(56)

Assuming that the bulk density of TiC\(_x\) particles precipitated around the C source at the primary stage of the in situ reaction under the constant gravity field is “\(\epsilon\)”, then the volume of TiC\(_x\) particles precipitated in the melt per unit time is

\[ V_{\text{occ}} = \frac{V_{\text{TiC}_x}}{\epsilon} \]  
(57)

To ensure that the TiC\(_x\) particles generated on the surface of the C source do not accumulate, the TiC\(_x\) generated in the previous stage needs to make room for the newly generated TiC\(_x\), i.e., the TiC\(_x\) particles on the surface of the C source need to float the distance of \(\frac{V_{\text{occ}}}{ab}\) in a unit time; then,

\[ \frac{dR}{dt} = \frac{V_{\text{occ}}}{ab} \]  
(58)

From Equations (36), (54), and (56–58), we can get

\[ \frac{|\rho_{\text{TiC}_x} - \rho_{\text{Cu-Ti}}| D_p^2 \omega^2 R}{18 \eta} = \frac{k_{\text{Ti,C}_l} (4 + x)}{\epsilon \rho_{\text{TiC}_x} \left(1 + \frac{D_{\text{Ti},2}}{D_{\text{Ti},1}} + \frac{h_{k_{\text{Ti}}}}{D_{\text{Ti},1}}\right)} \]  
(59)

Then,

\[ \omega = \sqrt{\frac{18 \eta k_{\text{Ti,C}_l} (4 + x)}{\epsilon \rho_{\text{TiC}_x} \rho_{\text{Cu-Ti}} D_p^2 R \left(1 + \frac{D_{\text{Ti},2}}{D_{\text{Ti},1}} + \frac{h_{k_{\text{Ti}}}}{D_{\text{Ti},1}}\right)}} \]  
(60)

That is, when \(\omega \geq \) the right-hand value in Equation (60), the TiC\(_x\) accumulation effect can be completely eliminated by the super-gravity field, thereby speeding up the TiC\(_x\) generation reaction.

3. Conclusions

(1) For larger-size solid C sources, the kinetic models for spherical, cylindrical, and flat C-sources derived in this paper can be used to describe the reaction process. For C source particles with a small size that cannot easily distinguish the shape of the particles, a kinetic model incorporating a shape factor can be used for analysis.

(2) The introduction of a super-gravity field can promote the release of TiC\(_x\) particles from the surface of C source particles, while accelerating the mass transfer. It apparently shows the speed acceleration of the in situ reaction of solid C source in Cu-Ti melt.

(3) According to the theoretical derivation, when the centrifugal velocity exceeds a certain threshold, the effect of the super-gravity field can completely avoid the accumulation of TiC\(_x\) particles on the surface of the C source.
**Author Contributions:** L.G. established the mathematical model and wrote the manuscript. X.W. did the super-gravity experiments. Z.G. gave constructive suggestions for this work and helped design the experiments. All authors have read and agreed to the published version of the manuscript.

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**Abbreviations**

\( \bar{n}_{Ti,1} \) Diffusion flux of Ti atoms from the distant Cu-Ti melt to the layered TiCx surface (mol/s);
\( r \) Radius of the C source particles (m);
\( D_{Ti,1} \) Diffusion coefficient of Ti atoms in the Cu-Ti melt (m²/s);
\( C_{Ti} \) Concentration of Ti atoms (mol/m³);
\( C_{Ti,b} \) Concentration of Ti atoms in the Cu-Ti melt (mol/m³);
\( C_{Ti,s} \) Concentration of Ti atoms on the layered TiCx surface (mol/m³);
\( d \) Thickness of the layered TiCx (m);
\( h \) Concentration transition layer thickness of Ti atom in the Cu-Ti melt (m);
\( \bar{n}_{Ti,2} \) Diffusion flux of Ti atoms through the layered TiCx product layer (mol/s);
\( D_{Ti,2} \) Diffusion coefficient of Ti atoms in the layered TiCx product layer (m²/s);
\( C_{Ti,r} \) Ti atom concentration on the surface of the solid C source particle (mol/m³);
\( k_r \) Chemical reaction rate constant;
\( r_{Ti,b} \) Ti atom consumption rate of the overall process (mol/s);
\( r_C \) Consumption rate of solid C source during the reaction (mol/s);
\( \rho_C \) Molar density of the solid C source (mol/m³);
\( x \) Stoichiometry of the product TiCx;
\( r_0 \) Initial radius of the solid C source particles (m);
\( t \) Reaction time (s);
\( a, b \) Length and width of the flat solid C source (m);
\( r_{Ti,p} \) Rate of reaction defined by the consuming molar rate of Ti per unit time per unit surface area of C source (mol/(s·m²));
\( S \) Surface area of the C source particle (m²);
\( A_p \) Original surface area of the C source;
\( F_p \) Shape factor of the C source;
\( V_p \) Original volume of the C source;
\( R \) Distance between the particles and the rotating axis of the centrifugal machine (m);
\( \rho_{TiCx} \) Density of TiCx particles (kg/m³);
\( \rho_{Cu-Ti} \) Density of molten Cu-Ti alloy (kg/m³);
\( \Delta \rho \) \( |\rho_{TiCx} - \rho_{Cu-Ti}| \)
\( D_p \) Diameter of TiCx particles (m);
\( G \) Super-gravity coefficient;
\( g \) Gravitational acceleration (m/s²);
\( \eta \) Viscosity of liquid (Pa·s);
\( \eta_{Cu} \) Viscosity of molten Cu (Pa·s);
\( \eta_{Ti} \) Viscosity of molten Ti (Pa·s);
\( \eta_{Cu-Ti} \) Viscosity of Cu-Ti melt (Pa·s);
\( \omega \) Rotational angular velocity (rad/s);
\( f_{Cu} \) Mass fraction of Cu in the Cu-Ti alloy;
\( f_{Ti} \) Mass fraction of Ti in the Cu-Ti alloy;
\( \eta_0 \) Viscosity of the molten metal without particles (Pa·s);
\( \varepsilon \) Particle volume fraction of TiCx;
Maximum packing fraction of TiCx;  
Volume of TiCx produced per unit time;  
Volume of TiCx particles precipitated in the melt per unit time.

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