Instrumentation for comparing night sky quality and atmospheric conditions of CTA site candidates

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ABSTRACT: Many atmospheric and climatic criteria have to be taken into account for the selection of a suitable site for the next generation of imaging air-shower Cherenkov telescopes, the “Cherenkov Telescope Array” CTA. Such data are not available with sufficient precision, thus a comparison of the proposed sites and final decision based on a comprehensive characterization is impossible. Identical cross-calibrated instruments have been developed which allow for precise comparison between sites, the cross-validation of existing data, and the ground-validation of satellite data. The site characterization work package of the CTA consortium opted to construct and deploy 9 copies of an autonomous multi-purpose weather sensor, incorporating an infrared cloud sensor, a newly developed sensor for measuring the light of the night sky, and an All-Sky-Camera, the whole referred to as Autonomous Tool for Measuring Observatory Site Conditions Precisely (ATMOSCOPE). We present here the hardware that was combined into the ATMOSCOPE and characterize its performance.

KEYWORDS: Large detector systems for particle and astroparticle physics; Gamma telescopes; Instrument optimisation; Gamma detectors

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1 Introduction

Ground-based observations in the regime of very high energy (VHE) gamma-rays using Cherenkov light produced in extensive air showers constitute a young branch of astronomy. The aim of the future Cherenkov Telescope Array (CTA) [1], currently finishing its preparatory phase, is to leave behind the pioneering phase of current Imaging Air-shower Cherenkov Telescopes (IACTs) [2] and provide the astrophysical community with one mature and reliable gamma-ray observatory in both hemispheres, for the observation of gamma-rays with energies from a few tens of GeV to more than 100 TeV [3]. This entails a strong requirement to maximize the duty cycle of the observatories and hence minimize the influence of the various atmospheric disturbances, as well as possible anthropogenic contributions to the Light of the Night Sky (LoNS), at the main wavelengths covered by the Cherenkov light, namely the UV and blue. A very careful selection of the two sites is therefore crucial.

The requirements for an IACT site are different in many ways from those for optical or infrared (IR) telescopes. Since IACTs observe the development of air showers throughout the atmosphere at different heights, the altitude of atmospheric phenomena plays an important role when it comes to possible degradation of energy threshold, collection area and angular resolution. Typical observations by IACTs require long integration times, hence the achievable effective duty cycle of the observatory (now energy and angular resolution dependent) becomes a key parameter for site selection. Very high clouds may lie entirely above the observed air showers and do not hinder observations above a certain threshold energy. Air turbulence or seeing, in turn, are not so critical
since the requirements on the quality of the optics of IACTs are not comparable to those for modern optical telescopes. Finally, the prohibitive cost of providing domes for these large telescopes means that extreme weather phenomena such as low temperatures, high winds or sandstorms become crucial parameters for their design.

To help finding the best site for the CTA, all information about observation and climate conditions that is available from existing instrumentation including satellite data can be evaluated. However, satellite instruments do not always provide comparable measurements and require validation from the ground. Moreover, satellite data archives offer many possibilities, but comparing the candidate sites based on such datasets is not possible in some cases, due to limited temporal and/or spatial resolution and the impact of site topologies, which can be difficult to interpret.

It therefore makes sense to cross-calibrate each candidate site using identical hardware, and for this purpose the CTA community decided to design and set up standardized sensor stations for collecting LoNS, cloud and weather data. With the remote location of some candidate sites, implying a lack of infrastructure and availability of experts for maintenance, the requirements on such an instrumentation are challenging: relative compactness, independence from power grid, easy maintenance and remote internet control. The devices should be combined into a compact and rugged station, supplied with solar panels, providing off-grid power supply, computing power and possibly also a communication link to the Internet or temporary storage and regular pick up of the data. Such requirements imply reliability resulting from mechanical simplicity, low power consumption and limited bandwidth. The eventual choice of instruments was: a multi-purpose weather sensor, an infrared cloud sensor and specially developed sensors for measuring the LoNS. The basis and main instruments of this sensor station were developed and designed at the Max-Planck-Institute for physics (MPI) in Munich, Germany, and later improved by the Centre de Physique des Particules de Marseille, France. The device was given the name Autonomous Tool for Measuring Observatory Site Conditions PrEcisely (ATMOSCOPE). At the same time, the Joint Laboratory of Optics of the Palacky University Olomouc, Czech Republic, and the Institute of Physics of the Czech Academy of Science developed All-Sky-Cameras which were later incorporated in all but one ATMOSCOPE station.

2 Mechanical design and power supply

The primary requirement of the design of the site monitoring device was to construct a compact and robust station that can operate independently of the grid and under harsh outdoor conditions. The power is therefore provided by solar panels and stored in lead-acid batteries. The different instruments, shown in figure 1, are mounted on a mast 2.5 m above the ground. The weather station and the cloud sensor were relocated to a separate 10 m mast for the final installation. Figure 2 shows a picture of the installation, before relocation of the weather station and cloud sensor.

All the structural elements such as the mast and quadropod are made from 5 cm × 5 cm profile aluminium tubes. The power for the ATMOSCOPEs is generated by solar panels providing a peak power of 100W, which is used to charge two lead-acid batteries with a capacity of 100Ah each. All electric and electronic parts are housed in a cabinet suitable for outdoor electric installations, placed below the solar panels for additional protection from rain and direct sunlight. The power system is designed to avoid shutdown even in the case of several days of bad weather: out of all the
Figure 1. ATMOSCOPE instrumentation on top of the mast.

Figure 2. Picture of an Autonomous Tool for Measuring Site COnditions PrEcisely (ATMOSCOPE), a long term site evaluation measurement station for CTA.
Table 1. Summary of different instruments and the most important technical details of the weather station.

| measured parameter | technical properties                     |
|--------------------|------------------------------------------|
| temperature        | -40 to 60 °C (±0.3 °C)                   |
| pressure           | 600 to 1000 hPa (±0.8 hPa)               |
| humidity           | 1 to 100% (±2%)                         |
| wind speed         | 0 to 150 km/h (±2 km/h), starting at 0.5 m/s |
| wind direction     | 0 to 360° (±5°, hysteresis < 8°, starting at 0.5 m/s) |

ATMOSCOPEs operated over several years, the batteries were exhausted on one site only, leading to a cumulative loss of 26 hours. The cause was persistent bad weather in winter.

The mast is kept in place by four steel ropes under tension, going to the tips of the quadropod. The whole station is to be set up in a leveled position with the mast side pointing towards the equator. The orientation of each ATMOSCOPE with respect to the local coordinates is hence fixed, and can later be used for the analysis of star light captured by the LoNS sensors. The precision of the alignment of the optical devices is of the order of one to a few degrees, depending on the site. This characteristic can be measured by comparing the evolution of the light of bright stars passing in the field of view with the expected one.

3 Weather station with cloud sensor

For collecting the basic weather parameters relevant for site selection, a compact, commercially available weather station was selected. The decision was made to employ a Reinhardt MWS4 [4], which can be read out via a RS232 interface and measures temperature, barometric air pressure, relative humidity, wind speed and direction (see table 1). In addition, it was equipped with a cloud sensor offered by the same company [5], which uses a thermopile to measure the difference between the ambient temperature and the radiation temperature of the sky. It then infers the altitude of a cloud assuming a constant temperature lapse rate of −6.5 K/km.

The interval at which weather data were taken, initially set at one minute as for the LoNS measurement, was decreased to 2 seconds after a few months of operation in order to better estimate the time profile of weather parameters, especially the wind.

IR radiometer measurements, such as that from the thermopile, are affected by changing temperature profiles, dry or wet adiabatic lapse rates and cloud or haze layers, which may be too thin and not 100% opaque. The cloud sensor has a full field of view of the order of 10°, thus for an inhomogeneous cloud coverage, only a mean value is measured. Saturation is reached for an estimated cloud altitude of ~7 km; however, experience with the installed instruments showed that only clouds below about 4000 m above ground could be detected reliably, depending on the site.

The cloud sensors have been cross-calibrated against each other by placing them next to each other and exposing them to the same sky. The first such cross-check between 3 ATMOSCOPEs at MPI in Munich served as the basis for the overall cross-calibration. Figure 3 shows the data of about 3 days of measurements during mostly covered skies from the cross-check dataset before and after correction for the different response of the individual sensors. By plotting the measured cloud
Figure 3. Thermopile cloud altitude measurements vs time (top) and mean of the 3 sensor measurements (bottom), before (left) and after (right) cross-calibration. The individual devices show considerable measurement differences, but these follow nicely a linear correlation allowing corrections to be made.

altitude of each instrument against the mean of the 3 sensor measurements a linear relationship became apparent. The data from each cloud sensor were therefore treated with a linear function, such that all reproduce the same mean value. This procedure yields only a relative calibration; on the other hand we were not too interested in precision with respect to the cloud height, but rather were aiming to measure the average cloud coverage.

Nevertheless, once installed at the sites, we discovered that residual dependencies of the cloud height on at least ground temperature and relative humidity exist, which need to be corrected for as well. Figure 4 (top) shows an example of such a dependency. The retrieved cloud height maxima (corresponding to cloudless nights) were fitted with a linear function with respect to both parameters and corrected (see figure 4 bottom). The temperature gradients obtained range from -25 m/K to +65 m/K for the different ATMOSCOPEs, while the relative humidity gradients are always negative, ranging from -25 m/% to -4 m/%. After correction, each ATMOSCOPE was found to be sensitive to quite different maximum cloud heights, ranging from 4500 to 7300 m a.s.l.

Tests in the lab had revealed that the measured cloud altitude depends both on water and dust deposits on the lens. The thermopiles were never cleaned on-site, and typically show sensitivity degradations of roughly 1 m per day with time (see figure 5 left). During long periods of rainfall, wet dust deposits may solidify on the lens and then disappear only slowly and cause temporary degradation of sensitivity. Figure 5 (right) shows such an example.

While this method offers data for relatively low clouds at low cost and low instrumental complexity for comparative studies of the average cloud coverage during the nights, the obtained wide range of sensitivities caused some concern about the feasibility of such studies. Consequently, the device was finally only used for rough estimates.
Figure 4. Top: cloud heights, obtained from one deployed Reinhardt thermopile during a half-year measurement campaign, as a function of ground temperature (left) and relative humidity on the ground (right). Bottom: the corrected cloud heights, plotted against temperature and relative humidity. The coloured entries correspond overwhelmingly to cloudless nights, hence the black line shows the assumed sensitivity limit of this device.

Figure 5. Cloud heights, obtained from two deployed Reinhardt thermopiles as a function of time, expressed as Modified Julian Days (MJD). The coloured entries show the measured cloud height and the black line the derived sensitivity limits of the devices. While on the left side a gradual decrease of sensitivity with time is observed, mainly attributed to dust, the right side shows a two months period of frequent rain, starting around MJD = 56460. At the end of that period, the thermopile has lost sensitivity which it slowly recovers during the next 250 days. That the drop of sensitivity was not mimicked by clouds has been cross-checked with images from the All-Sky-Camera.
4 All-Sky-Camera

A more precise measurement of the cloud coverage during nights is achieved with the use of All-Sky-Cameras (ASCs), installed at each candidate site on the ATMOSCOPE structure, opposite the LoNS sensor\(^1\) (see figure 6). The ACS system consists of a fish eye lens of type FUJINON YV2.2x1.4A-SA2 [6], an astronomical CCD camera of type G1-2000 [7], a control computer and associated electronics. The ACS is fed by the power supply of the ATMOSCOPE and uses the same internet connection, but the measurements are taken independently and do not interrupt measurements by the other devices of the ATMOSCOPE. The CCD is equipped with an ICX274AL chip which has a quantum efficiency greater than 50% in the range from 450-550 nm. Its resolution is 1600 × 1200 pixels, and the pixel depth is 16-bit monochrome. The fish-eye varifocal lens has a 185° field of view, and its iris can be controlled electronically. It is covered by a “zero” meniscus lens, correcting for the divergence of the incident beam.

The camera and lens are housed in a solid aluminium waterproof body and connected to a miniPC via USB and a separate serial I/O controller, which controls the iris and the power switch of the camera. The optical switching electronics turns the system ON after sunset and keeps it going during astronomical night-time. The system is OFF during the day, in order to save energy.

Images are taken automatically every five minutes and processed using the control computer of the device. The resulting analysis output is the cloud fraction (CF), defined as the percentage of the sky covered by clouds. The data are analysed on-site, and only the results are sent to the data server using the ATMOSCOPE internet connection.

A calibration of the system’s fish-eye and zero meniscus lenses was performed in the lab to correct for optical aberrations and vignetting effects. For this aim, a point-like optical source (a

\(^1\)This is true except for one site where an ASC of a different type was already installed.
spot as small as possible) was placed at a distance of at least 10 m from the lens. Then, the body of the camera was rotated around the CCD detector’s sagittal axis from 0 to 180 degrees. This procedure was repeated for the transverse axis. Intensity and pixel position of the light source as a function of incidence angle were then calculated (see figure 7).

The cloud fraction was defined as the ratio of the number of recognized image stars and catalogue stars in a field of view from 0 to 60° zenith angle. To estimate the second number, we used the visual magnitudes [8]. The catalogue stars have been filtered with respect to their magnitude and calibration. A magnitude limit for stars close to 60° zenith angle of slightly larger than 4\textsuperscript{m} was chosen relaxing gradually towards zenith, where > 6\textsuperscript{m} was applied. For each star in the filtered catalogue, we search for a detected star within an angular distance of < 1° (see figure 8). The ratio of paired divided by unpaired stars gives then the CF. Typically about 900 catalogue stars are tested over the whole observed night sky. The precision of the algorithm was estimated using artificial cloud simulations which yielded an uncertainty of the calculated CF of always less than 5% in absolute CF, depending slightly on the simulated CF (see figure 9).

5 LoNS monitor

The main instrument of the ATMOSCOPEs for measuring the night sky brightness is the LoNS sensor. To avoid too much complexity and keep the data volume low, an integrating rather than imaging type of device was used. The final design uses imaging optics to map a certain part of the sky onto a large area PIN diode and measure the integral light flux via the photo current. A starlight model predicting the amount of “natural light” from this field of view can then be used to quantify

![Figure 7. CCD positon and star magnitude limit vs. zenith angle, as obtained from the ASC calibration.](image-url)
Figure 8. An example of the cloud analysis of an ASC image of the sky partly covered by clouds. Left: raw image without analysis, right: with the result of the analysis superimposed. The yellow crosses show all detected stars, out of which the green squares show those detected stars whose position matches a catalogue star; the small red circles depict catalogue stars without correspondence in the ASC image. These are typically obscured by clouds. In the lower right area of the images, the 10 m mast with the weather station is visible. The large red circle shows the analysis limit of $60^\circ$ zenith angle.

Figure 9. Dependence of the cloud cover retrieved vs. simulated one. The blue points show all trials while the red points with error bars show the mean and RMS for each simulated CF bin. The red points have been linearly interpolated; the curve is used as a correction function to avoid small residual biases.
Figure 10. LoNS measurement device, open box with the filter wheel removed. From left: PCB with PIN diode, motor for turning the filter wheel, Hall sensors for checking the position of the wheel.

Figure 11. Quantum efficiency of the Hamamatsu S3584-08 PIN diode, measured with the MPI Quantum Efficiency Measurement Device (left). Transmission curves of the two filters, which have been combined to achieve the “ATMOSCOPE-B” filter for the LoNS sensor (right).

The fraction of light pollution from artificial light sources as well as to detect clouds inside the field of view, which also modify the light flux.

The light sensor is a $28 \times 28$ mm Silicon PIN diode of type S3584-08, manufactured by Hamamatsu [9] (see figure 10). The quantum efficiency curve of the device, measured in the lab at MPI, is shown in figure 11. The PIN diode is operated in photo-voltaic mode. The photo current is amplified by a trans-impedance amplifier with $10^8$ VA$^{-1}$ in the first stage and by two optional secondary stages with an amplification of 10 and 100 maximizing the dynamic range. The precise
measurement of very weak currents in our first amplification stage requires an amplifier with a very low input-bias current and a low temperature drift. A device that offers these properties is the OpAmp AD549 [10] from Analog Devices. The manufacturer claims an input bias of 250 fA and a temperature offset drift of $5 \mu VK^{-1}$. A capacity of 1 nF in the loop-back of the amplifier guarantees an integration time $RC = 0.1 \text{s}$, to avoid noise pickup and oscillations.

To be able to measure the LoNS in different wavelength bands and for taking dark frames before each measurement, the instrument is equipped with a 5-position filter wheel (see figure 12).
Figure 13. Transmission curves of the Hamamatsu S3584-08 PIN diode (full red), the lens (dotted brown) and the ATMOSCOPE filters (violet and yellow). The blue and orange lines show the standard Johnson/Bessell blue and visible filters for comparison. The dotted lines show the normalized transmission of the filters only, the full lines after convolution with the quantum efficiency of the PIN diode. For comparison, the normalized spectral acceptance of a Sky Quality Meter (SQM) [11] is shown in green.

The ‘closed’ position enables dark frames to be taken. Another position is equipped with an astronomical V-band filter from Lot-Oriel [12]. The next position houses a set of two colour glass filters from Schott [13] that in combination resemble an astronomical B-band filter but have an extension towards the UV. The remaining two positions are empty.

Figure 13 shows the transmission curves for the filters used by the ATMOSCOPE. These match more-or-less the standard Johnson/Bessell visible and blue filters [14–16], but with some shift of the mean wavelength. The “ATMOSCOPE-B” filter moreover shows broader spectral acceptance in the UV than the standard blue filter (see figure 13) and hence covers the major part of the spectrum of Cherenkov light from air showers observed on ground. The effect of these shifts needs to be corrected for if star catalogues from standard filter photometry are used for calibration, or in order to compare results with other instruments. Peak transmission ranges from 0.85 to 0.88 for the visible band and from 0.61 to 0.71 in the blue/UV band.

Special care has been taken to control filter leakage at other wavelengths, especially in the near infrared, where the night sky is much brighter (see figure 14). The visible filter shows transmission values always below the $10^{-4}$ level outside the nominal window, and the cumulative contribution of leaking light from longer wavelengths can be conservatively estimated to be below 0.1% of the light registered inside the nominal window. The blue/UV filter shows stronger leakage, estimated to account for $\sim 3\%$ for typical night-sky spectra [17]. A small leak around 560 nm shows transmission up to $2 \cdot 10^{-2}$, and coincides with the OI 557.7 nm emission line from airglow. Another broader
Figure 14. V-filter and “ATMOSCOPE-B” filter leakages measured for an ATMOSCOPE produced at MPI. Shown are the transmissions, with and without folding with the quantum efficiency of the PIN diode, and the result folded with a typical light-of-night-sky spectrum [from 17]. The integral light-of-night-sky leakage into the V-filter window amounts to 0.04%, the B-filter leakage into the V-filter window amounts to 2.2% and into the infrared leakage to about 0.4%. Measurements courtesy of Jose Luis Rasilla, IAC.

Leak allows $10^{-3}$ of the light to pass through at wavelengths around 690 nm. We conservatively estimate that the first leak may contribute up to 3% of the nominal filter wavelengths, but this is well controlled by the measurements performed with the visible filter, which has overlapping acceptance. The red and near infrared leakage is smaller and estimated to contribute always less than 1%.

To demonstrate that the ATMOSCOPE incorporates a good choice of filters for the site selection of the CTA, figure 15 shows the natural light of night sky spectrum, due to air-glow and the minimum zodiacal light [17], folded with the quantum efficiency of photomultipliers foreseen for the 23 m Large-Size Telescope (LST) [18]. That telescope type has been chosen because its sensitivity will be limited most severely by the LoNS, compared to the Medium-Size (∼12 m) and Small-Size (∼5 m) telescopes. One can see that the “ATMOSCOPE-B” filter covers almost half of the background light responsible for photomultiplier noise, while the V-filter covers about 20%.
Figure 15. Example of a night-sky spectrum from an astronomical site, here the Observatorio del Roque de los Muchachos at La Palma [17], in units of photons per nano-seconds, per nano-meter per pixel of the camera of the large-size telescopes of the CTA (black line). To a good approximation, anthropogenic backgrounds can be considered to be limited only to the two lines Hg\textsubscript{435.8} and NaD\textsubscript{589.0/6} for the ranges of interest here. The brown line shows the same spectrum, folded with the quantum efficiency of the Hamamatsu S11920 PMT, used for the large-size telescopes of the CTA. Also shown are the responses of the two filters used in the ATMOSCOPE, together with those of the two standard Johnson/Bessell visible and blue filters for comparison. The numbers below the filter curves show the integral in units of photo-electrons per nano-second, and their relative contribution to the observed photo-electron background.

Figure 16. Example of one night of night-sky measurements with both filters (yellow: V-filter, blue: "ATMOSCOPE-B"-filter”). Left: northern hemisphere with Jupiter entering the field of view, right: southern hemisphere with Sirius entering the field of view.

There is only a small overlap between both filters, and the contribution of light from 600 nm onwards is largely suppressed by the spectral acceptance of the photomultiplier. Note that several cameras developed for the smaller telescopes of the CTA are planned to be equipped with Geiger-mode avalanche photo diodes (often referred to as “Silicon PMs”), which show significant spectral acceptance up to 800 nm. However efforts are ongoing to cut out the (unwanted) acceptance by using filters in this range.
The lens of the LoNS device has a physical diameter of 5 cm and a focal length of 6 cm. Its orientation has been chosen in order to produce a sharp cut-off at the edges of the acceptance function, such that bright point sources, like planets or bright stars, can be distinguished by a sudden change in the diode current when they enter the field of view. This feature was later used for fine-calibration of the effective orientation of the LoNS sensor (see figure 16). With the PIN diode in its focal plane, the field of view is $\approx 0.38$ sr. Ray-tracing simulations (figure 17, right) show that not all the light collected on the surface of the lens is imaged onto the PIN diode. Assuming normal incidence, the maximum radius at which light does not undergo total reflection on the inside of the lens can be calculated using $R_{\text{max}} = 1/n_c \cdot R_c$, where $R_c = 30.9$ mm is the radius of curvature of the lens, and the $n_c$ the refractive index of glass, which can be approximated by $n_c \approx 1.52$ for the V-filter wavelengths, and $n_c \approx 1.53$ for the “ATMOSCOPE-B” filter. This results in an entrance pupil of $12.9 \text{–} 13.0$ cm$^2$.

Due to the square shape of the PIN diode, the acceptance of the light sensor is not symmetric in azimuth ($\Phi$). The angular acceptance of one ATMOSCOPE has been measured in the laboratory of the MPI Munich. This was done by scanning the radial incidence angle ($\theta$) in $1^\circ$ steps for 3 different $\Phi$ positions, $0^\circ$, $30^\circ$ and $45^\circ$ (see figure 17). Due to mechanical constraints in the setup, the measurements were not extended to zenith angles larger than $60^\circ$. Values between $60^\circ$ and $70^\circ$ were obtained through extrapolation. The resulting acceptance is symmetric with respect to the sign of the zenith angle (see figure 18). One can see that the acceptance drops down to less than $1\%$ at $30^\circ$ incidence angle, as expected, but it rises again to reach a secondary maximum of $3\%$ around $45^\circ$ incidence angle. Although $3\%$ seems a small number, the enhanced solid angle between $35^\circ$ and $60^\circ$ makes this contribution significantly larger than $10\%$. The secondary acceptance can be explained with single scattering of light on the surface of the ring which holds the lens. Although the ring is painted with black anti-reflective paint, a $3\%$ reflectivity is quite possible. Any tertiary acceptance due to multiple scattering should be suppressed by at least a factor $0.03^2 \sim 10^{-3}$, and is hence expected to contribute only insignificantly to the overall acceptance.

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2Alternatively a circular diaphragm could have been used to make the acceptance rotationally symmetric. This way one would however have sacrificed some of the angular acceptance.
Figure 18. Interpolated data from measurements at MPI of the angular acceptance of the ATMOSCOPE, as a function of the radial incidence angle $\theta$. The three curves show different azimuthal angles, where $\Phi = 0$ corresponds to incidence on the side of the PIN diode, $\Phi = 45$ along its diagonal. Left: linear scale, right: logarithmic scale.

Figure 19. Angular acceptance of the ATMOSCOPE, shown in polar coordinates up to an incidence angle of 30 degrees (left, linear scale) and up to 70 degrees (right, logarithmic scale).

A perfect 8-fold symmetry of the system was then assumed, and the statistical measurement uncertainties doubled at every 10° distance from the closest azimuth measurement. Even so, the statistical uncertainties of the overall measurement of the cumulative angular acceptance are less than 1%. The result of this analysis is shown in figure 19. These measurements were made with a blue LED, and hence no spectral dependency of the acceptance could be measured. However, given the small variation of the refractive index of the lens glass with wavelength, this effect is likely limited to less than 3%.
The resulting cumulative angular acceptance is $0.375 \pm 0.005$ sr (see figure 20). The statistical and systematic uncertainty stems almost entirely from the extrapolation from 60° to 70°. Degradations of the acceptance due to dirt and aging are not included here, and have been calibrated on-site about once per two months.

The three analogue PIN diode current measurement channels are read by a MAXIM 128 ADC, with a range of 0–4.095 V and a resolution of 1 mV. The communication to the control computer is via I²C, a simple serial protocol, using one data line and one clock line. There is a second optional temperature channel that can be used to monitor the temperature inside the electric control cabinet. The humidity inside the LoNS sensor box is also monitored by a sensor connected to the MAXIM 128.

The custom-made filter wheel is turned by a DC gear motor which is switched on and off with a PhotoMOS relay from the computer. The positioning feedback comes from two Hall sensors mounted on the filter wheel, that detect the iron plates passing through. One sensor gives the computer the stop signal at each of the five positions, and the other, which is mounted at the position of the dark measurement (taken before each run) provides absolute positioning. The diode and the electronics are installed inside a water proof box, with the lens elevated slightly above the cover, positioned at the end of a tube.

The average LoNS on a dark site between 300 and 600 nm is about $1.7 \cdot 10^{12}$ ph m$^{-2}$ sr$^{-1}$ s$^{-1}$ [20]. Assuming an average QE of about 50% for the “ATMOSCOPE-B” filter, an average filter transmission of about 0.65, the collection area of the lens of $1.3 \cdot 10^{-3}$ m$^2$ and a field of view of 0.38 sr, roughly $2.7 \cdot 10^8$ ph s$^{-1}$ would be expected for the ATMOSCOPE to be measured under dark sky conditions. The corresponding PIN diode produces about 0.27 photo-
Figure 21. Our measurement of the linear temperature drift of the empty AD549 amplifier (diode not plugged in).

electrons per nano-second which yields a diode current of 43 pA. In the three different readout channels with amplification of \((10^8 \, \text{VA}^{-1}, 10^9 \, \text{VA}^{-1} \text{ and } 10^{10} \, \text{VA}^{-1})\), this gives signals of 4.3 mV, 43 mV and 430 mV, respectively.

The PIN diode itself produces a temperature-dependent dark current, even in the absence of light. Since it is operated in photo-voltaic mode, the PIN diode will reach a saturation reverse current which can be described by [see e.g. 21]:

\[
I_s = A \cdot \frac{D_p \cdot c_n^2 \cdot T^{(3+\gamma)/2}}{L_p \cdot N_D} \cdot \exp\left(-\frac{E_g}{kT}\right) \\
\approx \frac{10^{11} \, \text{ns}^{-1} \, \text{K}^{-3}}{(L_p/100 \, \mu\text{m}) \cdot (N_D/10^{21} \, \text{m}^{-3})} \cdot T^{(3+\gamma)/2} \cdot \exp\left(-\frac{E_g}{kT}\right) \tag{5.1}
\]

\(A = 7.84 \cdot 10^{-4} \, \text{m}^2\) is the PIN diode area. \(D_p = 1.2 \cdot 10^{-3} \, \text{m}^2/\text{s}\) is the diffusion coefficient of holes in silicon. \(c_n = 3.28 \cdot 10^{21} \, \text{m}^{-3} \, \text{K}^{-3/2}\) is the effective density of states in the un-doped layer. \(L_p\) is the diffusion length of holes, and \(N_D\) is the donor density. The additional exponent \(\gamma\) parameterizes the temperature dependency of the diffusion length. Finally, \(E_g = 1.12 \, \text{eV}\) is the band gap energy of silicon, and \(k\) the Boltzmann constant.

The saturation current \(I_s\) is expressed here in units of electrons per nano-second. Additionally, the amplifier shows a linear temperature drift (see figure 21). A small voltage offset was chosen by design, in order to ensure always a positive dark current measurement.

Figure 22 (left) shows the measured dark currents as a function of temperature of one ATMOSCOPE, fitted to appropriate temperature dependencies for the PIN diode, amplifier bias current and offset. The measurement points were taken continuously from dark measurements throughout about one year. The fit describes the data well, as can be seen in the fit residuals plot on the right side. Less than 4% of the fit residuals lie outside the fitted Gaussian; this slight deviation may be explained by circumstances in which the PIN diode has a slightly different temperature than the sensor on the PC board. Fitting all 9 deployed ATMOSCOPEs with the combination of these
dependencies, we obtain \((2.3 \pm 0.2)\) for the term \(1/((L_p/100 \mu m) \cdot (N_D/10^{15} \text{ cm}^{-3}))\), which is compatible with the assumption that \(L_p\) is about the same size as the depletion region, i.e. \(~30 \mu m\), and an \(N_D\) of the order of \(10^{15} \text{ cm}^{-3}\). The fits yield results for the additional exponent \(\gamma\) always smaller than 0.014, but larger than 0.0002. Finally, the temperature drift of the amplifier contributes with \(dI_{\text{amp}}/dT = (7 \pm 4) \cdot 10^{-5} \text{ ns}^{-1} \text{K}^{-1}\).

The temperature dependency of the dark current is hence well understood, and always meets the specifications of the manufacturer. On the other hand, fluctuations around the fitted dependencies are still too large to predict the dark current with sufficient precision using temperature measurements alone. For this reason, care was taken to measure the dark current with sufficient precision a few seconds before each filter measurement. This is achieved by putting the filter wheel into the completely closed position for recording the dark current before each LoNS measurement is made with the filters. At each position the ADC is read 1000 times and the average value as well as the RMS are stored. The measured average dark current is immediately subtracted from the filter measurements, and stored in the data for cross-checks.

We define the signal-to-noise ratio (SNR) as the background subtracted filter current, divided by the statistical uncertainty of the measurement:

\[
\text{SNR}_{B,V} = \frac{I_{B,V} - I_{\text{dark}}}{\sqrt{\text{RMS}(I_{\text{dark}})^2 + \text{RMS}(I_{B,V})^2}}.
\] (5.2)

The obtained SNR is always above 50 for the V-filter measurements, and above 30 for the “ATMOSCOPE-B” filter, even in the absence of anthropogenic light. This is enough to reduce statistical fluctuations of all measurements to well below the systematic uncertainties.

Even in one case of a faulty high-gain channel, the SNR obtained from the medium gain always yielded values greater than 8.

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**Figure 22.** Left: the measured mean dark current (expressed in photo-electrons per nano-second) vs. the temperature measured at the PC-board and fitted to diffusion dark current contribution, an amplifier bias offset drift, and a constant offset (red line). The measurement points correspond to a one-year measurement on a site with relatively small temperature gradients.
Calibration of the LoNS sensors

Before the deployment of the devices a cross-calibration was performed in the laboratory. A reference light source made of a continuous, wide-spectrum LED and two polarizers was used for this purpose. Tuning the angle between the two polarizers allows for variation of the level of light according to Malus law for measuring the response of a LoNS sensor, and for checking the linearity across at least two gain ranges. Examples of calibration curves are shown figures 23 and 24, the latter presenting the layout of the device. The knowledge of the absolute level of light generated by the device is not necessary, as this apparatus aims at providing a reproducible set of light intensities over a wide range to compare the responses of different LoNS sensors. The ratio of the measured responses gives the cross-calibration factor, which covers the whole system: lens, filters, photo-diode, and readout chain. The accuracy of a cross-calibration between two sensors is of the order of ±2% at same LED current and temperature.

One of the LoNS sensors has been kept in the laboratory and has served as reference for calibrations: each LoNS sensor has been calibrated relative to this reference. Figure 23 shows the response of a sensor which was later sent to a candidate site versus the response of the reference sensor: a difference of 18% is measured for the response with the “ATMOSCOPE-B” filter, while only a 4% difference is measured for the V filter. This is consistent with the observation that the spread between batches of Schott B-filters was larger than for the Lot Oriel V-filters. Of the 9 LoNS sensors, the cross-calibration factors w.r.t. the reference sensor range from 0.96 to 1.07 when the V filters are in place, and from 1 to 1.2 with the “ATMOSCOPE-B” filters, respectively.

The sensitivity of the sensors is monitored on-site from time to time with a device as described above. In this case the main purpose is to measure the loss of transparency of the lens due to
Figure 24. Example of measurement of the “dirtiness”: the response of the LoNS sensor to a tunable light source is measured before and after cleaning the lens. The ratio between the two response curves yields the “dirtiness”, which is 5% in this example. The layout of the calibration device is also shown. Angle = 0 corresponds to crossed polarization directions.

Figure 25. Applied “dirtiness” corrections (red lines) to the light sensor data of one typical ATOMSCOPE, obtained from measurements (blue squares) as described figure 24.
Table 2. Compilation of systematic uncertainties of the ATMOSCOPE LoNS sensor.

|                           | V-filter | ATMOSCOPE “B”-filter | Comments                                      |
|---------------------------|----------|-----------------------|-----------------------------------------------|
| **Absolute Calibration LoNS Sensor** |          |                       |                                               |
| Filter transmission      | 1%       | 2%                    | spectral acceptance                           |
| Filter leakage           | 1%       | 3%                    | for typical dark sites                         |
| Angular acceptance       | 3%       | 3%                    | limitations of lab. measurements (spectral dependency, etc.) |
| Acceptance above 70°     | <2%      | <2%                   |                                               |
| Shadows ATMOSCOPE        | <1%      | <1%                   |                                               |
| Dirtiness/deposits       | 2%       | 2%                    |                                               |
| Light leaks              | <1%      | <1%                   |                                               |
| Lab. calibration         | 4%       | 4%                    |                                               |
| Entrance pupil lens      | 2%       | 2%                    |                                               |
| Dark current subtraction | <2%      | <3%                   |                                               |
| Temperature effects      | <0.5%    | <0.5%                 |                                               |
| **Total**                | <7%      | <8%                   |                                               |

|                           | V-filter | ATMOSCOPE “B”-filter | Comments                                      |
|---------------------------|----------|-----------------------|-----------------------------------------------|
| **Relative inter-calibration LoNS Sensor between sites, after deployment** |          |                       |                                               |
| Filter transmission      | 0.5%     | 1%                    | spectral acceptance                           |
| Filter leakage           | 1%       | 2%                    | for typical dark sites                         |
| Angular acceptance       | 1%       | 1%                    | for typical dark sites                         |
| Shadows Surroundings     | 1%       | 1%                    |                                               |
| Interpolation Dirtiness  | 2–8%     | 2–8%                  | on average <2%, larger only for indiv. days   |
| Inter-calibration lab.   | 4%       | 4%                    |                                               |
| Entrance pupil lens      | <1%      | <1%                   | gluing, etc.                                   |
| Dark current subtraction | <2%      | <3%                   | only for exceptionally strong temperature gradients |
| Temperature effects      | <1%      | <1%                   |                                               |
| **Total**                | 5–9.5%   | 6–10%                 | typically 5–6%                                |
dust deposition: a measure of the LoNS sensor response is performed before and after cleaning the lens, and the difference is attributed to dust deposits on the lens, the so-called “dirtiness” (see figure 24). Then, an approximation of the evolution of the loss of transparency along the time is performed by performing a linear interpolation between the day of the last cleaning and the current one. Typical correction factors range from zero to eight percent over periods of several months (see figure 25). An absolute calibration of the LoNS sensor can be attempted either by comparison with other devices, or with the help of known stars passing across the field of view of the devices. The description of these methods is however outside the scope of this paper, and will be treated in a future publication.

Table 2 lists the residual systematic uncertainties after calibration, divided in absolute calibration uncertainties, valid for comparison of the results with those obtained from a different instrument — such as external measurements by optical telescopes — and relative uncertainties valid for the inter-comparison of sites. The totals have been calculated by adding the individual items in quadrature.

7 The Sky Quality Meter

A commercial LoNS measurement device, a Sky Quality Meter SQM-LE, from Unihedron, is also part of the ATMOSCOPE (see figure 26). While the SQMs are not as accurate as the LoNS, we have decided to use them for the following reasons: the SQMs are extremely robust, and can work for extended periods of time with no maintenance. Moreover they are the standard tools to measure the night sky background providing a check of the results as well as easy comparison to many sites worldwide.

The precision is $\pm 0.1 \text{ mag/arcsec}^2$ according to the provider, and the operating temperature range is $-40$ to $85^\circ$C. The SQM is housed in a protective tube with a glass window, and operates continuously, with one measurement being performed each minute. The angular acceptance of the SQM is $20^\circ$ FWHM. SQMs are known to have considerably broader spectral acceptance than through a B and V-filter alone (recall figure 13). Cinzano [11] has shown that the broader acceptance may result in typical LoNS measurements taken with the SQM differing from standard V-band measurements by up to $0.25 \text{ mag/arcsec}^2$.

The actual accuracy of the SQM was found to be worse than the specifications. To cross-calibrate the SQMs, they were placed side by side in a controlled low light environment, as the light intensity was varied. We present an example of the results in figure 27. The figure shows the measured differences between three SQMs and one that was chosen as reference. The results show substantial offsets of up to $\pm 0.4 \text{ mag/arcsec}^2$ between various units. However, for each unit individually, the spread between the measurements is smaller than the $\pm 0.1 \text{ mag/arcsec}^2$ listed in the specifications. The offsets measured in the laboratory were later used in the data analysis.

Table 3 lists the most important technical details of the two devices employed to measure the LoNS.
Figure 26. Details of the housing of the SQM: the SQM is powered using the Ethernet cable (Power over Ethernet, 5V) and the readout is performed each minute across the Ethernet by a simple script. The glass housing protects the SQM against the rain.

Table 3. Summary of different instruments to measure the LoNS and the most important technical details.

| instruments | technical properties |
|-------------|----------------------|
| LoNS sensor | acceptance: 0.375 ± 0.005 sr, intercalib. precision: 0.06 mag/arcsec² |
| SQM         | acceptance: 0.17 ± 0.02 sr, intercalib. precision: 0.17 mag/arcsec², temperature range: −40 to 85°C |

8 Control computer and software

As control computer we use a TS-7260 single board computer (see figure 28). It uses an ARM9 CPU, running at 200 MHz, has an SD-card slot that can be equipped with up to 2 GB memory and a battery-supported real time clock. The overall power consumption is less than 1 W. In addition, it has three serial ports, one of which is used to read out the data from the weather station. There are also 30 independently usable I/O lines that are used for the I²C communication and the motor control, and a watch dog timer, to restart the board in case it hangs. The internal clocks of the ATMOSCOPEs show considerable drifts, advancing by on average 20 µs/s, with a spread of 5 µs/s between different clocks (see figure 29). Monthly synchronizations of the clocks have been made since autumn 2012. This has reduced the clock drift corrections to always less than one minute, the data between synchronizations being corrected assuming a linear drift with time.

3Official manufacturer claim: 20° FWHM, corresponding to 0.17 ± 0.02 sr according to the shape of the efficiency curve provided by Unihedron.
Figure 27. An example of the cross-calibration measurement of the SQMs. Shown are the differences between the reference SQM and three other units as a function of the reference SQM value.

The computer runs a Debian Linux operating system that is installed on and booted directly from SD-card. This makes it possible to compile the software on board. Setting up the system for a second or third station means just changing the boot procedure to SD-boot and copying the SD card. The data are stored on a USB key connected to one of the 2 available USB ports on the board. The data files are uploaded each hour on a data server located in Warsaw, using the `rsync` Linux utility. For this purpose, all ATMOSCOPEs are connected to the network via 3G connections or wireless links when a nearby observatory is reachable. This allows also for an hourly monitoring of ATMOSCOPE measurements, as well as remote checks of the system and the monthly tuning of clocks mentioned above. The network link is achieved using a router (3G or Wireless) and a switch serving a local network loop to which the TS-7260 board is connected, as well as the SQM, the All Sky Camera PC, and, on some sites, auxiliary devices.

All programs are written in C and called by entries in the `crontab`. After the measurement process is launched, the filter wheel is still in the dark frame position, as set after the last measurement. First the temperature sensors and the humidity monitor, attached to the ADC, are read, and the physical values are calculated from the voltage values. Then the three ADC channels, assigned to the LoNS sensor amplification stages, are read out 1000 times each and averaged to increase the precision. The statistical uncertainties are also calculated to get an estimate of the fluctuations.
Figure 28. TS-7260 single board computer, used for the operation and readout of the ATMOSCOPE.

Figure 29. Internal clock delays, as estimated from one typical ATMOSCOPE. The jumps correspond to those moments when the clock was synchronized.
All voltages are converted back to the corresponding currents using the appropriate amplification factors. During daytime, too much light reaches the PIN diode for the measurements to be useful. Therefore, if the voltage value in the first amplification stage exceeds 100mV, the wheel does not turn at all so that the PIN diode is protected. If it is dark enough outside, the four other filters are selected in sequence and the voltage measurement procedure described above is repeated each time. The weather station data taken simultaneously are added to the measurement results of the LoNS device. Finally, all values and errors, together with the system date and the Unix time stamp, gathered in one single string, are appended to a file.

The SQM data are recorded each minute by sending a simple reading request over the local network loop; they are stored in a separate file. Each record consists of the system date, the magnitude per square arc second, the temperature at the level of the light sensor, and few auxiliary parameters.

9 Conclusion

Ten ATMOSCOPEs (Autonomous Tool for Measuring Observatory Site COnditions PrEcisely) have been built and installed at 9 CTA candidate sites in the Northern and Southern hemisphere, while one reference unit stayed in the laboratory. Each apparatus allows the measurement of weather parameters by means of a commercial weather station measuring wind, temperature, humidity, atmospheric pressure and cloud altitude, and allows the measurement of the Light of the Night Sky (LoNS) by means of two sensors: a commercial one (Sky Quality Meter) and one specific sensor designed and built at the MPI Munich. The latter is based on a large-area calibrated PIN-photodiode and two filters, in the V band and a customized “ATMOSCOPE-B” band, which matches the spectral acceptance of super-bialkali photomultipliers. The ATMOSCOPEs are complemented by All-Sky-Cameras to derive a cloud fraction parameter during the night.

On all sites the weather station, initially held by the ATMOSCOPE structure at 2.5 m above the ground, has been relocated to the top of a mast in 10 m height for being less sensitive to aerodynamic ground effects caused by surface roughness and vegetation. The optical devices, assisted by the weather measurements, have performed a precise LoNS estimate during the clear nights, measuring the light of the sky each minute over periods of up to 2.5 years, depending on the site. By means of wireless or 3G communication, depending on the presence of neighboring observatories, data have been retrieved hourly to allow for online monitoring.

The data provided by the devices equipping the ATMOSCOPEs have been analyzed and used to give an evaluation of each site. The analysis methods applied and results obtained will be presented in another publication. Figure 30 shows an example of raw data obtained by the two optical sensors taking data simultaneously over 8 days, with the cloud altitude given by the weather station and the cloud fraction measured with the All-Sky-Camera. The correlation between the measurements can be seen directly: a darkening of the sky seen by the optical devices (larger mag/arcsec$^2$) correlates with the presence of clouds, indicated by the decreasing cloud altitude measured by the cloud sensor, the increasing cloud fraction (here within 20° of the zenith direction) given by the All-Sky-Camera and large variability of the measured light. The V-shape of the LoNS measurement is due to the Galactic plane passing across the field of view of the optical devices. The last night displayed on the top panel shows the influence of the Moon which transits above the horizon at the beginning of the night.
Figure 30. Example of data provided by the LoNS sensor, the SQM (top) and the cloud sensor and All-Sky-Camera (bottom panels) running together. The Galactic latitude observed by the sensors in the zenith direction is shown on the second panel.

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