Nanomaterials for Energy Conversion and Storage
Nanomaterials for Energy Conversion and Storage

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Progress [1] in the area of nanoscience and nanotechnology has pervaded almost all areas of science and technology. Over last couple of decades the ability to manipulate and control materials at an atomic and molecular level (nanometer range) and subsequent understanding of the fundamental processes at nanoscale have led to new avenues. The knowledge thus acquired can be translated into innovative processes, leading to design or fabrication of better products. More importantly, new scientific phenomena and processes have emerged that could provide either revolutionary or novel solutions to the energy, environmental, and sustainable mobility challenges that will face humanity in the 21st century.

With demand for clean and sustainable energy sources increasing at an exponential rate [1], new material technologies are being explored that could provide cost-effective and environmentally clean solutions to the world’s energy problems. Developments in the areas of alternative fuels or energy storage technologies like advanced batteries, fuel cells, ultra capacitors, and biofuels are emerging as strong contenders to petroleum-based sources. Energy derived from clean and renewable sources like solar and wind power have tremendous potential, but the practical use of these sources of energy requires efficient electrical energy storage (EES) technologies that can provide uninterrupted power on demand. In all of these new technologies, nanomaterials are increasingly playing an active role by either increasing the efficiency of the energy storage and conversion processes or by improving device design and performance.

Lithium-ion batteries [2] are one of the great successes of modern materials electrochemistry. Their science and technology have been extensively reported. A lithium-ion battery consists of a lithium-ion intercalation negative electrode (generally graphite) and a lithium-ion intercalation positive electrode (generally the lithium metal oxide, LiCoO$_2$), these being separated by a lithium-ion conducting electrolyte, for example, a solution of LiPF$_6$ in ethylene carbonate-diethylcarbonate. Although such batteries are commercially successful, we are reaching the limits in performance using the current electrode and electrolyte materials. For new generations of rechargeable lithium batteries, not only for applications in consumer electronics but especially for clean energy storage and use in hybrid electric vehicles, further breakthroughs in materials are essential, such as the use of nanomaterials devices.

Supercapacitors [2] are of key importance in supporting the voltage of a system during increased loads in everything from portable equipment to electric vehicles. There are two general categories of electrochemical supercapacitors: electric double layer capacitors (EDLC) and redox supercapacitors. In contrast to batteries, where the cycle life is limited because of the repeated contraction and expansion of the electrode on cycling, EDLC lifetime is in principle infinite, as it operates solely on electrostatic surface charge accumulation. For redox supercapacitors, some fast faradic charge transfer takes place and results in large pseudocapacitance. Progress in supercapacitor technology can benefit by moving from conventional to nanostructured electrodes. In the case of supercapacitors, the electrode requirements are less demanding than in batteries, at least in terms of electrode compaction, because power prevails over energy density. Thus, the benefits of nanopowders with their high-surface area (primary nanoparticles) are potentially more
important, hence the staggering interest in nanopowders and their rapid uptake for supercapacitor-based storage sources.

Fuel cell technologies [2] are now approaching commercialization, especially in the fields of portable power sources—distributed and remote generation of electrical energy. Already, nanostructured materials are having an impact on processing methods in the development of low-temperature fuel cells \((T < 200\, ^\circ C)\), the dispersion of precious metal catalysts, the development and dispersion of nonprecious catalysts, fuel reformation and hydrogen storage, and the fabrication of membrane-electrode assemblies (MEA). Polymer electrolyte membrane fuel cells (PEMFCs) have recently gained momentum for application in transportation and as small portable power sources; whereas phosphoric acid fuel cells (PAFCs), solid oxide fuel cells (SOFCs) and molten carbonates fuel cells (MCFCs) still offer advantages for stationary applications, and especially for cogeneration. Platinum-based catalysts are the most active materials for low-temperature fuel cells fed with hydrogen, reformate, or methanol. To reduce the costs, the platinum loading must be decreased (while maintaining or improving MEA performance), and continuous processes for fabricating MEAs in high volume must be developed. A few routes are being actively investigated to improve the electrocatalytic activity of Pt-based catalysts. They consist mainly of alloying Pt with transition metals or tailoring the Pt particle size.

Due to the potential applications in the field of environmental protection, the photochemistry of TiO\(_2\) is a fast growing area both in terms of research and commercial activity [3]. Over the past decade, the level of research activity concerning TiO\(_2\) can be appreciated through the exponential increase of relevant research literature produced (11,500 publications between 1993 and 2003) and the number of patents regarding, for instance, photocatalysis (3000 between 1996 and 2001, and more recently 2500 in Japan and 500 in USA). These features account for the intense fundamental research activity involving the scientific community. This particular interest is related to the remarkable photoactive properties of TiO\(_2\) and therefore to its numerous applications, which are related to two main environmental priorities: environmental protection through heterogeneous catalysis (water purification, air cleaning, and self cleaning materials) and renewable energy production in photovoltaic solar cells, dye sensitised solar cells (DSSC). In addition to this the use of TiO\(_2\) nanostructures has enhanced the photo electrochemical properties of these materials.

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Research Article

Lithium Ion Storage Characteristics of Mechanically Fractured Titanate Nanotubes

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The effect of mechanical milling on the formation of short titanate nanotube and structural change induced is investigated. Mechanical milling produces the short nanotubes with the length of 30–160 nm. The lithium ion intercalation characteristics of the obtained short titanate nanotube were studied to verify the effect of the newly formed cross-sections of nanotubes. It was found that the protonated titanate nanotubes maintained long shapes until 30 min of mechanical milling and were transformed into agglomerated nanosheets and finally anatase granules depending on the treatment duration. Through galvanostatic investigation, the nanotubes with milling of 15 min exhibited the highest discharge capacity of 336 mAh·g⁻¹ in first cycle, 12.4% larger than pristine.

1. Introduction

Nanostructured electrodes fabricated from nanotubes, nanowires, or nanofibers have shown unique rate capability for lithium ion batteries since the distance Li⁺ ions penetrate is as small as several nanometers in the radius direction which is significantly smaller than that in the usual power. As a result, there is tremendous research interest in the development of nanostructured Li-ion battery electrode with high capacity and excellent cycling stability [1–6].

Many researchers have recently prepared titanate nanotubes and developed their electrochemical properties by precise control over their morphology for various applications. It has been found that the properties of these nanomaterials depend strongly on their size and shape. Therefore, it is important to control their size and the morphology to make use of them in suitable applications. Especially much attention has been paid to nanotubes for electrochemical lithium storage due to their open, mesoporous structure, large specific surface, efficient transport of lithium ions, and impressive ion-exchange properties, which result in a high value of charge/discharge capacity [7–12].

Li et al. and Zhou et al. reported on the effect of the titanate morphology produced by changing process parameters on the battery performance [13–16]. Bavykin et al. reported the rate of ion exchange along length of titanate nanotubes [17]. For the shortest nanotubes (having an average length ca. 50 nm), the characteristic time of the ion exchange was shorter than that for long nanotubes, suggesting that the diffusion of ions along the length of the tube could be the rate-limiting step. Although this finding was pioneering, they did not investigate the effect of titanate nanotube length on the electrochemical characteristics as an anode material for lithium battery. It is our main objective to carry out the basic estimation on the battery performance of titanate nanotubes with various lengths and verify the possibility of improving the poor cyclability of titanate nanotubes, by changing the nanotube length.

For an anode for lithium ion battery, it was confirmed that, at the first discharge stage, lithium intercalation occurred at inner and outer surface of nanotube. Some of these intercalated lithium ions could not be extracted from deep interlayer sites on inner surface of the nanotube and, hence, electrochemical reactions occur actively only at the both end of nanotubes [15, 16, 18]. In the interests of large surface area for electrochemical reaction sites of LiₓTiO₂ and short nanochannel for easy and fast transport
the produce of mechanical milling.

of Li ion for lithium intercalation, we synthesized the short length of nanotubes by low-energy mechanical milling from the hydrated titanate nanotubes synthesized via alkaline hydrothermal treatment of TiO₂ and subsequent acid treatment [19, 20]. To verify the effect of mechanical milling treatment on titanate nanotube for battery performance, the length of the titanate nanotube was controlled by various milling condition, and basic charging-discharging behavior of nanotubes as anode was estimated.

2. Experimental Methods

Titanate nanotubes were prepared by alkali hydrothermal method. Commercial titanium oxide powder (P25, Degussa AG) was used as a precursor to synthesize titanate nanotubes (TNTs). 4 g of the TiO₂ precursor and 180 mL of 10 M NaOH (98%, SAMCHUN) aqueous solution were mixed under stirring for 3 h. The suspension was poured into a Teflon vessel and placed in autoclave for hydrothermal treatment at 130°C for 24 h and cooled to room temperature. The synthesized solution was washed for 24 h with HCl (35%, DAEJUNG) acid solutions in pH 1 and then washed continuously with distilled waters to pH 7 to remove to extract sodium. The samples were filtered and dried at 60°C for 24 h. Titanate nanotubes were treated in a mechanical ball milling using identical milling parameters. The milling performed in air using a Fritsch planetary ball mill, with balls made of zirconium dioxide. The powder-to-ball weight ratio was chosen to be 1 (powder): 15 (ball), while the milling time was varied from 15 to 60 min under the rotational speed of 100 rpm.

The morphology and structural changes of titanate nanotubes with various milling conditions were characterized using BET (Surface Characterizing Analyzer, ASAP2010 & TriStar3000, USA), (X-ray diffraction XRD, Rigaku d-max 3000, Japan), (High Resolution Transmission Electron Microscopy TEM, JEM-2100F, Japan), and (Field Emission Scanning Electron Microscopy FE-SEM, Hitachi S-4800, Japan).

3. Results and Discussion

Figure 1 shows the XRD patterns of the as-prepared titanate nanotubes and the mechanically milled in condition of 100 rpm for 15, 30, 45, and 60 min. All the powders exhibited characteristic peaks at around 20 = 10°, 24°, and 28°, which are assigned to the diffraction of H₂Ti₂O₅·H₂O (JCPDS, Card 47-0124) with peak broadening by nanometer size [20–23]. It was known that the nanotubes showed weak diffraction lines because of their tubular structure and poor crystallinity [24]. The gradual changes in the major diffraction peaks suggest that the mechanical milling affected both the length and the structure of nanotubes. The intensities of main peaks at 20 = 10° and 28° of titanate nanotubes were slightly decreased and broadened with increasing milling time, which suggests that nanotube structure was gradually distorted along with shortening by mechanical milling process. Unlike main peaks of nanotube, the peaks at 20 = 24° and 48.5° were maintained under whole milling conditions. Plodinec et al. confirmed that during the mechanical milling treatment the phase transitions from titanate nanotubes to anatase phase and finally rutile phase were observed since the high-energy ball milling generated defects, and it influenced the composition and the stoichiometry [25]. Similar phase change was observed in this work; nanotube structure was distorted and converted into anatase phase gradually. After milling for 60 min, most of nanotubes seemed to collapse into anatase crystalline granules.

The variation of surface area of the mechanically milled product is summarized in Table 1. The surface area increases with milling and reached the highest value of 550 m²·g⁻¹ at 15 min and then decreases with an increase in the milling time. This is believed to be due to the agglomeration of fractured nanotubes and the anatase granules. At 60 min of milling the surface area was decreased to 273 m²·g⁻¹, suggesting that the highly agglomerated anatase granules were formed as supported from XRD data in Figure 1.

Figures 2 and 3 show the HRTEM and FESEM images of the titanate nanotubes. As-prepared titanate nanotubes were (Thermogravimetric-differential thermal analysis TGD TA, TGA7, PERKIN ELMER) was performed with a temperature range of room temperature to 500°C.

For fabrication of electrodes, the active materials prepared with the mixture of carbon black, and isopropyl alcohol, and the obtained titanate nanotubes were coated on copper foil and, then, dried at 120°C for 12 hours to remove the solvents. The tests were performed on a coin-type cell assembled with as-prepared electrode, lithium metal as the counter, separator, and an electrolyte. The electrolyte solution was LiPF₆ dissolved in a mixture of EC (ethylene carbonate) and (dimethyl carbonate DMC) with the volume ratio of 1:1.

Electrochemical properties measuring were using TOCAT' cell test instrument (Toyo system, Japan). Charge and discharge were conducted between 3 and 1 V versus Li⁺ reference electrode at 0.1 C rate. Cyclic voltammogram was measured from 3 and 1 V range at 0.2 mV·s⁻¹, using electrochemical interface (SI 1287, UK).
Figure 2: TEM images of (a) the as-prepared protonated titanate nanotubes, and the titanate nanotubes with mechanical milling in condition of (b) 15, (c) 30, (d) 45, and (e) 60 minutes at 100 rpm.
Figure 3: SEM images of (a) the as-prepared protonated titanate nanotubes, and the titanate nanotubes with mechanical milling in condition of (b) 15, (c) 30, (d) 45, and (e) 60 minutes at 100 rpm.

Table 1: Surface area of the starting material, the as-prepared, and mechanically milled.

| Condition of milling (min) | P25 | 0  | 15 | 30  | 60 |
|----------------------------|-----|----|----|-----|----|
| Surface area, $S_{BET}$ (m²·g⁻¹) | 55–60 | 500 | 550 | 513 | 273 |

long and slim tubular shape, whose diameter was around 10 nm and length was several hundred nanometers. HRTEM image in as-prepared protonated titanate nanotubes shows that the synthesized nanotubes have 4-5 layers with open ends. The nanotubes were slightly shortened by milling for 15 min to approximately ~100 nm in length and milling for 30 min to 30–100 nm. However, the number of wall layers and the size of interspacing remained similar to those of the as-prepared long titanate nanotubes. After 45 min, the particle type granules were observed that the titanate nanotubes were fractured and transformed to anatase particles with excessive milling treatment. XRD data in Figure 1 support that short titanate nanotubes maintained the $\text{H}_2\text{Ti}_2\text{O}_5$·$\text{H}_2\text{O}$ structure up to 30 min of milling. Figure 3 shows that the partial change of powder morphology took place with milling treatment. It is clear that the degree of agglomeration is slightly intensified with milling time. After 60 min, tubular type morphology is not shown as TEM image, and anatase granules were agglomerated by increased treatment duration.

Figure 4 shows the initial charge/discharge behavior of lithium intercalation in the titanate nanotubes before
and after milling treatments at current rate of 10 mA·g⁻¹ between 1.0 V and 3.0 V. The initial discharge capacities were varied from 335 mAh·g⁻¹ to 210 mAh·g⁻¹, and initial charge capacities were varied from 280 mAh·g⁻¹ to 114 mAh·g⁻¹ depending on milling duration. As-prepared titanate nanotubes exhibited the initial charge and discharge capacity of 219 mAh·g⁻¹ and 299 mAh·g⁻¹, respectively, with a reversible efficiency of 73.2%. However, charge and discharge capacities were increased to 280 mAh·g⁻¹ and 336 mAh·g⁻¹ for milling of 15 min with a reversible efficiency of 83.3%. After milling of 15 min, the capacities rapidly decreased with increasing milling time. One can see that morphological and structural change, controlling the length of titanate nanotubes, seems to have a noticeable effect on the cell performance. At the initial stage of milling treatment, long nanotubes were fractured into short nanotube that had the highest value of surface area for the absorption of Li ion and short nanochannel for easy and fast transfer of Li⁺ for lithium intercalation. Further milling treatment, however, starts to produce nanosheets with lower intercalation capability. Due to the increase of less active phases, even though nanotubes are shorten further, the charge/discharge capacities decrease. After 30 min of milling, the nanotubes transformed to anatase granular structure, and consequently a distinguishing characteristic of nanotubes was disappeared.

Figure 5 shows the cyclic voltammograms of the samples at a scan rate of 0.2 mV·s⁻¹. In the first cycle, there is only a pair of broad redox peaks for nanotubes as prepared and milling 13 and 30 min. The pseudocapacitive behavior of this peak is explained in terms of the interaction taking place on the nanosheet surface, hence curved one. The TiO₂ nanotubes hold the curved nanosheet-like morphology with thin-wall structure and high specific surface area.

In the case of the nanotubes with milling 45 min, a pair of shoulder redox peaks near the larger broad redox peaks appears at 1.75 and 1.95 V (versus Li⁺/Li, depending on the scan rate), which is characteristic for lithium insertion into anatase lattice, revealing the possible phase transition from titanate to anatase. Obviously, the pair of shoulder redox peaks related to lithium insertion/extraction in anatase becomes larger with increasing milling time, accompanied by a decrease of the broad redox peaks and a structural transformation from titanate to anatase [26].

Figure 6 shows the galvanostatic charge/discharge recyclability for the titanate nanotube and the materials produced by milling treatments. The cutoff potentials for the charge/discharge were set to be 1.0 V to 3.0 V versus Li⁺/Li. All samples exhibited an obvious drop of discharge capacity during first several cycles that may be due to the formation of Li₂O₂ or LiOH by left remained water in nanotubes on the surface of electrode [27]. The electrochemical discharge capacities of titanate nanotubes with milling decreased gradually and showed poor cycle stability for lithium insertion and extraction compared to nanotube. In spite of highest value of initial charge and discharge capacities, the short titanate nanotubes produced by milling contain defects in composition and structure, and these defects cause irreversible capacities accumulating every cycles.

4. Conclusion

Titanate nanotubes were synthesized by the hydrothermal method and treated in a high-energy ball milling in various duration conditions. The structure of nanotubes was H₂Ti₂O₅·H₂O with inner and outer diameter of 5 nm and 10 nm. Increasing milling duration, the titanate nanotubes were fractured into short length of nanotubes, and finally transformed to anatase structure. The 15 min sample shows that high initial charge and discharge capacities of 280 mAh·g⁻¹ and 336 mAh·g⁻¹ observed in the titanate nanotube, which shows an effect of shortened titanate nanotubes. The length of nanotubes had an effect on cell performance due to high specific surface area for absorption of Li ion and short nanochannel for easy and fast transfer of Li⁺ for lithium intercalation; however, the cyclic stability was not greatly improved because of structural and compositional defects caused by high-energy ball milling. Further study is required the synthesis of short nanotube without structural
Figure 5: Cyclic voltammogram plots of TiO₂ nanotube electrodes (a) and with mechanical milling in condition of (b) 15, (c) 30, (d) 45, and (e) 60 minutes at 0.2 mV·s⁻¹.
change and defect, which can solve the problems of low initial capacity and poor cyclability of titanate nanotubes.

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Research Article

Nafion Titania Nanotubes Nanocomposite Electrolytes for High-Temperature Direct Methanol Fuel Cells

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Nafion-based nanocomposite membranes containing various amounts of titania nanotubes (TNTs) as an inorganic filler have been prepared using melt-mixing method and have been investigated for proton exchange membrane applications. The one-dimensional TNTs have been prepared from potassium hydroxide using hydrothermal route and conventional heating. Nafion R1100 in a protonated form was used, and TNT contents were in a range of 0.5–2.0 wt%. The acid-treated composite membranes, at lowest inorganic additive content, exhibited improved properties in terms of thermal stability and methanol (MeOH) permeability. The best performing nanocomposite was the membrane containing only 0.5 wt% TNTs showing ionic conductivity value of $7.2 \times 10^{-2} \text{ S} \cdot \text{cm}^{-1}$ at 26°C and 100% of relative humidity.

1. Introduction

Currently, titanium dioxide (TiO$_2$) or titania-derived nanotubes (TNTs) are examined in respect with their possible prominent applications such as catalyst supports [1], photocatalyst [2], gas sensing elements [3], and anode materials for batteries [4]. Titanium-dioxide-supported noble metal catalysts, such as gold [5], silver, or platinum [6, 7], are effective in many important processes, such as selective oxidation and hydrogenation and automobile engine exhaust systems [8].

Principal advantages of the nanotubular form seem to be related to the property that they have high specific surface area, which allows diffusion of active species. The crystal structure of TNTs seems to be a derivative for many photocatalytic applications. Another property is that TNTs are often open at the end and accessible for the interaction with gaseous substances. The TNTs are also thermally stable up to 600°C [6–9]. These are some of the essential criteria for fuel cell electrolytes membranes. This promotes to study the Nafion/TNTs nanocomposite membranes behaviour with the aim to improve Nafion properties such as fuel permeability and thermal and mechanical stability.

Nafion, whose primary structure consists of acid tipped side chains dangling from a perfluorinated backbone, is expected to be compatible with TNTs, due to its reactive nature. The metal ions present in TNTs such as K$^+$ is expected to increase thermal stability of Nafion membrane. The resistance of TNTs in presence of H$_2$ is likely to improve swelling and mechanical properties of Nafion membranes at higher temperatures greater than 100°C. Therefore, to understand the effect of the interactions between Nafion and TNTs on the proton exchange membrane properties, further investigations were carried out. In this study, the effects of TiO$_2$ nanotubes on Nafion properties such as water uptake, thermal stability, methanol (MeOH) permeability, and ion conductivity were investigated. The nanocomposite membranes were prepared using melt-extrusion method by varying the nanofiller loadings. The mechanical, thermal properties, and ionic conductivity of membranes were...
investigated using dynamic mechanical analyser (DMA), thermogravimetric analyser (TGA), and electrochemical impedance spectroscopy (EIS). The vibration studies were carried out by FTIR spectroscopy, providing information on the nanofiller/Nafion, nanofiller/H$_2$O, Nafion/H$_2$O interactions, and the conformation of fluorocarbon chains of Nafion hydrophobic domains. Methanol permeability was also investigated using the gas chromatography.

2. Experimental Details

2.1. Materials. Nafion precursor (R1100) was obtained from Ion Power. Potassium hydroxide (KOH) and titanium dioxide powder was purchased from Merck chemicals and used as received.

2.2. Synthesis of TiO$_2$ Nanotubes. TiO$_2$-derived nanotubes (TNTs) were synthesized using a procedure described elsewhere [9–11]. In a typical procedure about 23 g of TiO$_2$ powder, P25 Degussa, was mixed with 200 mL of 18 M of aqueous solution of KOH. The mixture was heated at 150°C in an autoclave with constant stirring for 24 hours. The product was then washed with deionized water and then dried in an oven at 120°C for 14 hours.

2.3. Membrane Preparation. The TNTs were used as filler during the preparation of polymer nanocomposite membranes. The nanocomposite membranes were prepared by melt-extrusion method using a microextruder Reomix OS (HAAKE) instrument. Nafion (R1100) and TNTs were melt-mixed at 250°C at a rotor speed of 60 rpm for 10 min. The TNTs were added after 2 min of melting Nafion inside the extruder. For each nanocomposite, the amount of TNTs loaded was varied from 0.5 to 2.0 wt% to study the effect of TNTs loading on the properties of Nafion. The nanocomposites samples were then converted into sheets or films with a thickness of about 0.2 cm using a Carver laboratory press at 2 MPa and 250°C. The nanocomposites were then hydrolysed, to introduce cation exchange properties into the membranes. The nanocomposites were immersed in a mixture of 15 wt% potassium hydroxide, 50 wt% of deionized water, and 35 wt% of dimethyl-sulfoxide at 80°C for 2 hours. This was followed by the repeated immersion (three times) in a fresh 3 M nitric acid (HNO$_3$) for 1 h to complete protonation process.

2.4. Characterisation Techniques

2.4.1. Scanning Electron Microscopy (SEM). Dry membranes were manually fractured after cooling in liquid nitrogen, to expose their cross section. The morphology of the fractured sample surfaces was analysed using SEM (LEO 1525), operating at an accelerating voltage of 3–10 kV.

2.4.2. The Fourier Transform Infrared (FTIR). FTIR spectroscopy 800 from Perkin Elmer that was used to study the elementary composition of the nanocomposites studies were carried out from 500 to 4000 cm$^{-1}$.

2.4.3. X-Ray Diffraction. XRD patterns were measured on a Phillips X’Pert diffractometer using secondary graphite monochromatic CuK$_\alpha$ radiation ($\lambda = 1.54060\,	ext{ Å}$) at 40 kV/50 mA. Measurements were taken using angle of incidence detector at an angle of 2°, for 2θ values over 10–70° in steps of 0.05°.

2.4.4. The Thermogravimetric Analysis. TGA was conducted on a TGA Q500 (TA Instruments) at a heating rate of 10°C/min in air, from room temperature to 800°C.

Mechanical properties of pure Nafion and its TNTs containing composite membranes were studied by a Perkin Elmer DMA 8000 in the dual cantilever bending geometry. The temperature dependence of loss tangent (tan δ, loss modulus/storage modulus) of all samples were measured at a constant frequency of 1 Hz with the strain amplitude of 0.02% from room temperature to 180°C with a heating rate of 2°C/min.

2.4.5. Electrical Conductivity. The electrical resistivity was measured at room temperature using the four point collinear probe method. The equidistant tungsten carbide probes have a separation distance (s) and probe radius of 0.127 cm and t 0.005 cm, respectively. The 1 × 1 cm$^2$ samples were prepared with a thickness of about 0.2 cm. The Keithley 4200-SCS Semiconductor Characterisation System, equipped with four supply and measure units (SMUs) and an amplifier, was used to perform high-precision direct current characterisation by supplying currents ranging from 1 fA to 1 mA. The resistivity changes as a function of probe spacing (0.127 cm) and is given by:

$$\rho = \frac{V}{2\pi s I}, \quad \text{for } t \gg s,$$

where $V$ is the measured voltage, $I$ is the supplied current, and $t$ is the thickness of the films [10, 12]. The electrical conductivity is then given by

$$\sigma = \frac{1}{\rho}.$$

2.4.6. Electrochemical Impedance. Proton conductivity of fuel cell membrane was measured using a galvanostatic four-point probe AC electrochemical impedance spectroscopy (EIS) technique, which is relatively so insensitive to the contact impedance that it could be adequate to accurately test membranes with high conductivity. A standard electrochemical conductivity cell was used as a reference. The experimental data was fitted to the equivalent circuits and the electrolyte resistance ($R$) was estimated from fitting procedure. The proton or ionic conductivity ($\sigma$) was calculated by

$$\sigma = \frac{d}{Rlt},$$

where $\sigma$ is the proton conductivity, $d$ is the distance between the electrodes and $R$ denotes the electrolyte resistance, respectively. $l$ and $t$ are the length and thickness of
the membrane, respectively. The electrochemical cell was connected using a two-point probe technique to an Autolab model 4.90006 potentiostat and Frequency Response Analyser (FRA). The home made cell was used with the electrode spacing of 1 cm, and the $3 \times 3$ cm² membranes were used. The FRA electrochemical impedance software was used for the impedance measurements and analysis from 1 MHz to 1 Hz with the AC voltage amplitude of 5 mV.

2.4.7. Liquid Uptake. The nanocomposites films were completely dried in vacuum at 100°C for 24 hours and weighed ($W_{\text{dry}}$) and then placed in water/methanol (50/50) by volume solution at 25°C for 24 h. The nanocomposites films were then wiped dry quickly with filter paper and weighed ($W_{\text{wet}}$). The water/methanol uptake was then calculated as

$$\text{Water/methanol uptake (\%) } = \frac{W_{\text{wet}} - W_{\text{dry}}}{W_{\text{dry}}} \times 100. \quad (4)$$

A methanol diffusion homemade cell was used based on previous reports [13]. Compartment A was filled with methanol solution and compartment B was filled with the same volume of deionized water. The membrane was clamped between the two compartments, and these were kept under stirring conditions during the experiment to avoid any concentration build-up in the two compartments. Methanol concentration profiles were determined over a period of time at different operating temperatures. Samples were kept for about 30 min between compartments. The concentration of methanol was determined by gas chromatography (GC).

2.4.8. Methanol (MeOH) Permeability. The measurements were carried out using a testing device constructed in our laboratories. The polymer membranes with surface area of $3 \times 3$ cm² were inserted between vessels A and B. A 250 mL solution of MeOH (2 M) was transferred to vessel A, whilst vessel B was filled with 250 mL of deionised water (H₂O). After 24 h, the amount of methanol that crossed through the membrane to vessel B was determined by gas chromatography (GC 17A, Shimadzu model). The gas chromatograph was equipped with a capillary column (14% cyanopropyl phenyl methyl polysiloxane, 30 m × 0.25 mm × 1.0 μm) and a flame ionisation detector (FID).

The permeability of methanol through the membrane can be expressed with diffusion coefficient, $P$ (cm² · s⁻¹) using equation (4) [14]

$$P = \frac{kV}{L \Delta C}, \quad (5)$$

where $k$ is the slope of the concentration profile (mol/cm³ · s), $V$ is the volume of a solution in compartment (cm³), $L$ is thickness of the membrane (cm), $A$ is the active surface area of the membrane (cm²), and $\Delta C$ is the difference of methanol concentration between two compartments (mol/cm³).

3. Results and Discussions

3.1. Structure and Morphology of Titania Nanotubes. The SEM results of the TNTs samples are shown in Figure 1. The samples show the tubular structures with diameter ranging from 10 to 20 nm. Large bundles of TNTs structures agglomeration are observed (see low magnification). This is attributed to the presence of K⁺ ions in the TNT structures, which can lead to an increase in the surface charge and electrostatic attraction between the tubes leading to agglomeration. This thermodynamic synthesis mechanism results in the formation of KTiO₂(OH) due to the presence of K⁺.

Figure 2(a) shows the FTIR spectra for TNTs that have been annealed at 300°C. The peaks corresponding to physically adsorbed water (O–H stretching mode at ~3400 cm⁻¹ and O–H bending mode at ~1640 cm⁻¹) all disappeared due to the higher annealing temperature of 300°C. This is consistent with the literature [15]. Only the Ti–O stretching mode is observed at ~640 cm⁻¹ and 400 cm⁻¹; this is attributed to the amorphous and anatase phase of TNT. XRD patterns for TNTs are shown in Figure 2(b) revealed the presence of both anatase and rutile phase. The additional peaks are ascribed to the interaction of K⁺ ions and TiO₂ and have been previously reported [10, 15].

3.2. Structural Properties of Nanocomposite Membranes. The FTIR spectra showing vibration bands of Nafion and Nafion/TNTs nanocomposites are shown in Figure 3. The bands around 1200–620 cm⁻¹ denote CF₂–CF₂ stretching, and the bands around 1050 cm⁻¹ denote the–SO₃-symmetric
Figure 2: (a) FTIR spectra of TNTs and (b) XRD patterns of TNTs synthesised using a hydrothermal conventional heating method and annealed at 300°C.

stretching. All these bands are characteristics of the Nafion membrane [14, 16–18]. The vibration bands of Nafion and nanocomposite membranes are similar; this might be due to the very small amount of TNTs loading or poor interaction with polymer matrix.

3.3. Morphology. The SEM investigation was performed on the nanocomposite membranes, and micrographs are shown Figure 4. Samples were frozen and fractured to expose the cross-sectional surface of the nanocomposite membranes containing various amounts of TNTs in wt% 4(a) 0.5, 4(b) 1.0, and 4(c) 2.0. The micrographs show the presence of inorganic structures, homogeneously dispersed into polymer matrix. This indicates the presence of bundle TNTs into Nafion matrix. SEM images show similar dispersion for all TNTs loadings. This is in agreement with XRD results, demonstrating that TNTs nanocomposite had the same d-spacing independent of TNTs loadings. This good dispersion suggests that the number of accessible sites on the surface of TNTs is high contributing to the adhesion and electrostatic interaction.

3.4. Water/Methanol Uptake and Electrical Conductivity. The hydration of the Nafion membranes is crucial to maintain higher proton conductivity. However, the methanol or fuel crossover must be reduced to avoid the catalyst poising and improve the fuel cell performance. The water/methanol uptake was analysed to maintain higher fuel cell performance. The water uptake increases from 19% for pure Nafion to a maximum of about 22% for 1.0 wt% TNTs composite as shown in Figure 5. The water uptake of N-0.5 wt% TNTs and N-2 wt% TNTs remains at about 20%. The increase in water uptake with addition of TNTs suggests an increase in hydrophilic nature of Nafion.

The bulk electrical conductivity of Nafion and its TNTs-based nanocomposite membranes is shown in Figure 5.

The electrical conductivity of Nafion increases from 0 to 0.06 S·cm⁻¹, when modified with 0.5 wt% of TNTs; it further increases to about 0.08 S·cm⁻¹, with 1.0 wt% incorporation. As the amount of TNTs into Nafion matrix increases to 2 wt%, the electrical conductivity decreases to zero, which is the same as that of pure Nafion. This decrease is associated with the electrical percolation threshold of Nafion and also the poor interaction between Nafion and TNTs as observed in FTIR spectra for 2 wt%.

3.5. Thermal Stability. Thermogravimetric analyses (TGA) of Nafion and its TNTs-based nanocomposite membranes are shown in Figure 6. Thermal degradation profiles of the nanocomposite samples are similar and show that the membranes are more thermally stable up to about 400°C compared to Nafion. The first degradation stage in a range of 30–120°C is assigned to the decomposition of physically adsorbed water present in the membranes. The weight loss in the range of 250–400°C is associated with the degradation of –SO₃H-groups of Nafion [17]. This increase in thermal stability is attributed to the presence of TiO₂ nanostructures into Nafion. The results show a slightly higher thermal stability for 0.5 wt% of TNTs. This is attributed to the good interaction between Nafion matrix, which is in agreement with the SEM results.

Figure 3: FTIR spectra of Nafion and its TNTs-based nanocomposite membranes with different TNTs loadings in wt% taken from 500 to 4000 cm⁻¹.

3.6. Proton Conductivity and Methanol (MeOH) Permeability. To evaluate the influence of TNTs on the membranes electrochemical behaviour, electrochemical impedance spectroscopy measurements were performed. Figure 7 shows a Nyquist plot of Nafion and TNTs-based nanocomposites at 100% relative humidity. The plot gives electrolyte resistance (R) of the materials. The nanocomposite containing 0.5 wt% TNTs has the lowest R = 210 Ω (see Table 1). The electrolyte
resistance increases as the TNTs loading increases to 1.0 and 2.0 wt%. This is indicative of a decrease in proton conductivity. This suggests that the amount of TNTs strongly influence the proton conductivity of Nafion. This decrease in proton conductivity can be due to the TNTs structures hindering the permeability of hydrogen ions. Whereas the electrolytes resistance decreases when 0.5 wt% TNTs was incorporated, indicating an increase in proton conductivity compared to pure Nafion. This increase in proton conductivity is related to the increase in water uptake observed with 1 wt% TNTs nanocomposite. The MeOH permeability was increased to $3 \times 10^{-8}$ cm$^2$ · s$^{-1}$ with 1.0 wt% TNTs nanocomposite and sharply decreases to $1 \times 10^{-8}$ m$^2$ · s$^{-1}$ as the amount of TNTs increases to 2 wt%. This behaviour is due to the MeOH permeability percolation threshold, which is affected by TNTs dispersed in Nafion.

The ratio of proton conductivity ($\sigma$) and methanol permeability ($P$) was calculated in order to estimate the performance of these nanocomposite membranes in fuel cell applications. The calculated values of proton conductivity and MeOH permeability are shown in Table 1. A decrease in the ratio was observed with an incorporation of TNTs in difference proportions.

4. Conclusion

In this paper, TNTs-based Nafion nanocomposite membranes were prepared in different proportions by melt-extrusion method. Water uptake of nanocomposite
membranes resulted higher than that of pure Nafion, with N-1.0 wt% TNTs reaching maximum of 21.5%. The highest electrical conductivity was observed with 1.0 wt% TNTs nanocomposite but remains at zero with 2.0 wt% TNTs nanocomposite. The thermal stability of Nafion increases with incorporation of TNTs. A reduced MeOH permeability was observed with 2.0 wt% TNTs nanocomposite membrane. The proton conductivity of Nafion slightly increases when 0.5 wt% TNTs and then decreases with an increase in the amount of TNTs. The best performing nanocomposite was the membrane containing only 0.5 wt% TNTs showing ionic conductivity value of $7.2 \times 10^{-2}$ S·cm$^{-1}$ at 26°C and 100% relative humidity.

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Challenges Found When Patterning Semiconducting Polymers with Electric Fields for Organic Solar Cell Applications

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A material-independent, contactless structuring method of semiconducting organic materials for the fabrication of interface-enhanced bilayer solar cells is not available so far. Patterning of thin films using electrohydrodynamic instabilities possesses many desired characteristics and has convincingly been used as a simple method to structure and replicate patterns of nonconducting polymers on submicrometer length scales. However, the applicability of this technique to a wider range of materials has not been demonstrated yet. Here, we report attempts to structure poly(p-phenylene vinylene) in a similar way. We found that thin films of poly(2-methoxy-5-(2′-ethylhexyl-oxy)-1,4-phenylene-vinylene) (MEH-PPV) and poly(2-methoxy-5-(3′,7′-dimethyloctyloxy)-1,4-phenylene-vinylene) (MDMO-PPV) could not be destabilized at all in the limited accessible range of the experimental parameters set by the delicate chemical nature of these materials. We discuss failure origins and present possible loopholes for the patterning of semiconducting polymers using electric fields.

1. Introduction

Applications ranging from the medical science to the semiconductor industry require the control of polymer thin film structures on submicron length scales [1, 2]. A variety of techniques such as lithography or imprinting methods are used for polymer structuring [3]. Some of these methods are still expensive, however, when lower than 100 nm resolution is required, or may not be applicable for patterning of larger areas. Alternatively, polymers can be structured by self-assembly processes, such as pattern formation via instabilities in thin polymer films or polymer demixing of multicomponent blends [4].

The most efficient organic photovoltaic cells today are based on a bulk heterojunction device consisting of a polymer donor/fullerene acceptor mixture, but the blend nanoscale morphology is difficult to control and depends largely on the processing conditions [5]. Therefore, new approaches that allow the controlled structuring of bicomponent mixtures of semiconducting materials are of general interest [6]. For organic solar cells [7, 8], the targeted morphology consists of a bilayer structure (~200 nm thick, to absorb the entire incident light) of a donor and acceptor component that interdigitate markedly on a ~20 nm length scale to enable efficient charge generation within the diffusion length (~5–40 nm [9]) of the photoexcited electron-hole pair. At the same time, the active components should wet their respective electrode completely, to provide conduction paths for the selective extraction of charges [10] (Figure 1). Concepts for the creation of such nanostructured interface topographical features with high aspect ratios (feature heights : widths) using bottom-up (polymer demixing) or top-down (nanoimprint lithography) have been presented [11, 12] (among others), but a general low-cost and large-area compatible pattern method for organic solar cells is not available so far.

With the intention of fabricating such highly folded interpenetrating donor-acceptor bilayers for organic solar cells, we tested here electric fields to destabilize thin films of semiconducting polymers. An electric field applied normal
to an interface between two dielectric materials polarizes the dielectric, which results in an effective displacement charge density and an interfacial electrostatic pressure. This electrostatic pressure couples to the capillary wave spectrum of the liquid interface, selecting a narrow band of wavelengths that is amplified. Electrohydrodynamic patterning (EHP) has been shown to be an effective means of overcoming interfacial interactions and to produce arrays of nanoscale structures [13]. Using mainly classical, non-conducting polymers so far, electric fields have been applied to induce instability and pattern formation at a liquid/air or at a liquid/liquid interface [14, 15]. Electric fields have also been used to replicate patterned electrodes with submicron lateral dimensions [16].

2. Materials and Methods

Polystyrene (PS, \(M_n = 150,000\), Aldrich) and poly(2-methoxy-5(3′,7′-dimethyloctyloxy)-1,4-phenylene-vinylene) (MDMO-PPV, \(M_n = 50,000–500,000\), American Dye Source) were used as received, poly(2-methoxy-5-(2′-ethylhexyloxy)-1,4-phenylene-vinylene) (MEH-PPV, \(M_n = 40,000–70,000\), Aldrich) was purified before use [11]. Films were spin coated from chlorobenzene (CB) on a first planar silicon electrode (2.5 \(\times\) 2.5 cm²) and were dried overnight before use. On a second silicon electrode, LiF was evaporated as an x-shaped spacer pattern in high vacuum; afterwards, the two silicon plates were assembled into the capacitor geometry leaving a polymer/air gap. EHP was carried out under vacuum in an oven, with electrical feedthroughs to a Keithley source meter. Experiments were carried out at \(T = 180\) °C, well above the glass transition temperatures and below the decomposition temperatures for both MEH-PPV and MDMO-PPV [17–19]. Rheology measurements (ARES RHEOMETRIC Scientific, parallel plate geometry, plate diameter 2.5 cm, plate gap 0.25 cm, shear rates 0.1–100 rad/s, under ambient air) were carried out for PS and MDMO-PPV in the temperature range of 160–240 °C. AFM experiments were carried out on a Nanosurf Mobile S instrument. We used the software WsXM to analyse the AFM images [20].

3. Results and Discussion

Figure 2 shows a schematic view of the experimental setup. The polymer film is deposited onto the surface of a bottom electrode and covered with a top electrode, separated at a small distance using an insulating spacer. The whole setup is heated above the glass transition temperature (\(T_g\)) of the polymer, and an electric field is applied between the two electrodes. If the electrostatic pressure exceeds the capillary pressure generated by the curvature of the film due to thermal-induced surface fluctuations, the film is destabilized and pillars evolve with a characteristic separation distance \(\lambda \sim \gamma^{1/2}H^{3/2}V^{-1}\) (for planar electrodes), where \(\gamma\) is the interfacial tension, \(V\) the applied voltage, and \(H\) the electrode spacing [13]. Decreasing the temperature below \(T_g\) finally freezes the morphology.

Since the electric field-induced instability is not material specific, it should be possible to destabilize essentially any liquid with reasonably low viscosity, as long as the polymer does not degrade when heated to high temperatures. The experiment involves low-cost equipment and allows to pattern larger areas, provided that the electrodes can be arranged precisely parallel to each other. These are essential prerequisites for organic solar cell applications, where a proof-of-concept device would measure an active area of several square millimetres at least. As an important potential advantage for semiconducting polymers, patterning involves minimal contact between the material and the electrode, whereas lithographic or printing techniques can compromise optical and electrical properties [3].

We compared the effects of electric field-induced destabilization using both insulating (PS) and semiconducting polymers (MEH-PPV and MDMO-PPV). Figure 3(a) displays a destabilized PS film with characteristics closely resembling those reported when using similar experimental parameters. We used planar silicon electrodes separated by a 280 nm thick LiF spacer. The PS film had a thickness of 100 nm, and a voltage of 25 V was applied. During a total annealing time (\(T = 180\) °C) of 22 hours, a small dielectric leakage current (9 mA) was flowing through the device [21].

This result indicates that the patterning method is robust and can easily be reproduced. Most often, destabilized features show a correlation over a length of several periods, beyond which domains with different orientations are found [13]. This would be of no concern for photovoltaic applications, where only the local heterojunction interface area should be maximised for charge transfer within the exciton diffusion length [9]. Enlarging the viewing angle for the PS sample, however, revealed that also the wavelengths of destabilized features varied considerably over a distance of 150 \(\times\) 150 \(\mu\)m² (see Figure 3(b)). This is probably a signature of small lateral variations of the electric field, introduced by slight variations in the plate spacing [4]. We conclude that the precise electrode adjustment and control of the air gap and capacitor plate distances for larger areas requires a more sophisticated experimental setup than used here, for example, where a gap monitoring scheme is used to fix the electrodes at a constant distance and parallel to each other [22]. In addition, the strong dependence of the developing
Figure 2: Molecular structure of the semiconducting polymers MEH-PPV and MDMO-PPV and the experimental EHP setup for the fabrication of organic solar cells. (a) A thin electron-donor polymer film is sandwiched between two electrodes separated by a distance $H$ using insulating spacers. (b) A voltage applied across the electrodes induces instabilities in the film that evolve with a characteristic wavelength $\lambda$. (c) To complete the solar cell, the patterned polymer film is covered by an electron acceptor material and a top electrode. Arrows represent sunlight exciting the solar cell. More elaborated experimental setups could involve the destabilization of donor-acceptor polymer bilayers directly or the guided destabilization via patterned electrodes or prestructured polymer films [11].

Figure 3: Atomic force microscope (AFM) images of (a) polystyrene (PS) and (c) MEH-PPV films after EHP using the same experimental conditions. Optical microscope image (b) of the patterned PS film. (d) Cross-section height profile of the MEH-PPV film shown in (c).
wavelength on the local distance indicates that the electrodes must be precisely even. This might limit the substrate choice in actual organic solar cell devices, where cost-effective applications in various fields, such as portable electronics or smart textiles, are expected in the near future on flexible devices.

Figure 3(c) shows an AFM image of a treated MEH-PPV film using similar experimental parameters as for PS (100 nm thick film, LiF spacer 280 nm, $T = 180\, ^\circ\text{C}$, $V = 20\, \text{V}$, $t = 6\, \text{h}$, leakage current $I = 9\, \text{mA}$). Apparently, the film did not destabilize and remained essentially flat with roughness below 1 nm (Figure 3(d)). Attempts with different experimental parameters, such as film thickness (20, 30, 80 nm), airgap (130 nm), or reversed polarity and intensity of the applied electric field, were unsuccessful. The same was observed for the structurally closely related polymer MDMO-PPV. Both PPV polymers were much less stable than PS when exposed to higher voltage and longer destabilization times, what could be visually observed by the development of black spots in the films. This limited the range of parameters that could be used.

The inability to destabilize the semiconducting polymer films was unexpected since at high enough fields, the electrostatic driving force overpowers other interfacial pressures and the materials characteristics should no longer be determining [13]. One reason for the inability to induce instabilities in MEH-PPV and MDMO-PPV might be related to the semiconducting nature of these PPV polymers. While PS behaves to a good approximation as a dielectric and its pattern formation behaviour can be explained by the perfect dielectric models [13, 23], charges are more easily injected and rearranged in semiconducting materials. For such materials, the “leaky dielectric model” [24] predicts an additional tangential force at the surface due to the presence of free charges, which will tend to decrease charge concentration. If the normal electrostatic force is capable of overcoming this effect, the presence of free charges is expected to increase the electrostatic stress at the surface, while lowering the electric field in the polymer. This would have little influence on the developing instability wavelength, but the temporal evolution of the instability would be accelerated by several orders of magnitude [13, 21]. Clearly, these predictions contradict with our experimental results and may suggest that tangential forces are important in these systems.

In a further attempt to destabilize the semiconducting film using EHP, we prepared a PS/MEH-PPV/air gap capacitor device [11]. Hierarchical structure formation using a polymer-polymer-air trilayer system has been convincingly shown using PMMA/PS [14]. Figure 4(a) shows the PS/MEH-PPV bilayer surface after 22 hour of EHP exposure, and Figure 4(b) displays the remaining MEH-PPV film after selective dissolution of PS. Apparently, PS in the bilayer polymer configuration destabilised in a similar pattern to a film that was coated directly on the silicon electrode (Figure 3(a)) [14]. Again, however, the MEH-PPV film did not follow the expected destabilisation (Figure 4(b)). For shorter process times (7 hours), the underlying MEH-PPV film remained flat (roughness below 1 nm, data not shown), while for longer process times the film breaks up erratically, with no preferential feature length scale emerging (not showing features that would suggest an EHP mechanism). In addition, such films changed visually colour and black spots appeared, suggesting that MEH-PPV has degraded.

A further reason that PPV films could not be destabilized might be due to the high viscosity of these stiff, backbone-conjugated materials. The developing pattern wavelength itself does not directly depend on the flow behaviour of the polymer melt, but the viscosity ($\eta$) influences the characteristic time ($\tau$) for film destabilization strongly [13]. $\tau$ scales linearly with $\eta$ [25] and patterning times ranging from seconds [21] to several days [14] have been reported for polymers with varying $\eta$ values. For PS, the viscosity showed the characteristic thermoplastic flow behaviour, namely, a high
temperature dependence of $\eta$ for low shear rates, a low temperature dependence of $\eta$ for high shear rates, and a general decrease of $\eta$ with increasing shear rate. On the other hand, MDMO-PPV behaved like a solid, and no meaningful rheological data could be obtained; we also observed that the polymer partly decomposed for temperatures above 200°C during the measurement. From this we conclude that MDMO-PPV could not be melted in the accessible temperature range without decomposition, which is probably a signature of the rodlike polymer backbone and the substantial polymer entanglement resulting from the synthesis [26]. Since the $T_g$ of MDMO-PPV ($\sim-45^\circ C$) [18] is considerably lower than that of MEH-PPV ($\sim66^\circ C$) [19], we argue that the latter would also have a similar solid-like behaviour.

4. Final Remarks

Electrohydrodynamic instabilities have been used during the last ten years as a powerful method to pattern polymers with submicron and ~100 nm features using planar and prestructured electrodes. To this end, mainly insulating polymers of low viscosity at higher temperatures have been used. Here, we tested EHP for the semiconducting polymers MEH-PPV and MDMO-PPV, two of the initial benchmark materials in the emerging field of organic solar cells. While PS could be readily structured using EHP in the most simple and large-area compatible approach, the semiconducting polymer thin films could not be destabilized at all. At the moment, we suspect this to be a combination of the semiconducting nature and/or the viscosity of these materials. The application of even higher temperature or voltage stress to induce destabilization is limited, since semiconducting polymers are labile and their optoelectronic performance degrades quickly with the formation of structural and chemical defects. It might be possible to use alternating electric fields to avoid the tangential movement of charges at the surface and to induce film destabilization. Another approach would be to reduce the polymer viscosity, for instance, by using plasticizers. For example, exposing PS to toluene vapour lowered $T_g$ below room temperature and the viscosity was greatly reduced in the swollen film [25]. For possible solar cell applications, the more fundamental issues of creating destabilized features with dimensions well below ~100 nm and high aspect ratios need to be addressed in greater detail. The aspect ratio of the columnar structures is usually rather small (~0.05–0.2 [22]), meaning that the interfacial increase of the donor-acceptor area compared to the planar bilayer configuration in a solar cell would be limited. In principle, by using higher electric fields, both wavelengths can be decreased and the aspect ratio can be increased at the same time. However, due to the electrical breakdown of the capacitor device, this approach is clearly limited. Higher aspect ratios can be created by modulating the electric field distribution using topographically patterned masks. Possibly, this could also be achieved by combining flat electrodes with pre-structured polymer films [11]. In such cases, however, the spatial density of mask (or polymer) pillars is limited. This is because the effective electric field modulation at a fixed distance decreases with increasing density of mask features; a smaller air gap would have to be used in these cases, finally limiting the attainable aspect ratios again [27].

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Research Article

Electrochemical Characterization of Li$_4$Ti$_5$O$_{12}$/C Anode Material Prepared by Starch-Sol-Assisted Rheological Phase Method for Li-Ion Battery

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Li$_4$Ti$_5$O$_{12}$/C composite was synthesized by starch-sol-assisted rheological phase method using inexpensive raw material starch as carbon coating precursor. The Li$_4$Ti$_5$O$_{12}$/C powder was characterized using XRD, SEM, and TG techniques. The synthesized Li$_4$Ti$_5$O$_{12}$ crystallites are cohesively covered by conductive carbon from starch sol which leads to increased conductivity, and the particle size of Li$_4$Ti$_5$O$_{12}$/C is about 500 nm. The electrochemical performance of Li$_4$Ti$_5$O$_{12}$/C was characterized by galvanostatic charge/discharge and EIS methods, and the results show that the Li$_4$Ti$_5$O$_{12}$/C presents a high discharge capacity, high rate capability, and long cycle life. The capacity retention was at 87% (500 cycles at 1C) and 73.0% (2000 cycles at 20C) indicating promising high rate performance of Li$_4$Ti$_5$O$_{12}$/C as anode material for lithium ion battery.

1. Introduction

Lithium ion batteries have been promising to be used in electric vehicle (EV) or hybrid electric vehicle (HEV). Currently, graphite was widely used as anode material for commercial Li-ion batteries in cell phones, laptops, and cameras [1]. However, its performance has not been satisfactory for applications in power systems which require frequent high rate charge/discharge such as in EV or HEV [2, 3]. Owing to the low lithium intercalation voltage of approximately 0.1 V (versus. Li/Li$^+$), lithium metal is easily deposited on the surface of graphite anode forming dendritic lithium particularly during fast charge, resulting in safety issue [2, 4]. Spinel Li$_4$Ti$_5$O$_{12}$ has acquired much attention as alternative anode to graphite in Li-ion batteries, due to its excellent reversibility, zero-strain structure, and high lithium ion mobility in the Li$_4$Ti$_5$O$_{12}$ lattice [5~7]. Especially, the material has a high Li-insertion voltage operating at 1.55 V (versus. Li/Li$^+$), which could hinder dendritic lithium formation [1, 5]. However, the low electronic conductivity ($10^{-13}$ S cm$^{-1}$) of Li$_4$Ti$_5$O$_{12}$ has been the main obstacle limiting rate performance of the material [6].

To improve electronic conductivity and electrochemical performance of Li$_4$Ti$_5$O$_{12}$, extensive works have been taken to address the issue, and the common strategies are to reduce particle size [4, 5], coat conductive material on the surface of Li$_4$Ti$_5$O$_{12}$ [2, 6, 7], or dope with other metal ions [8, 9]. Reducing particle size can shorten lithium ion diffusion distance and enhance both the electrode capability of Li storage and kinetics; thus the rate capability and electrochemical performance of electrode material can be improved [10]. Various metal ions such as Ni$^{3+}$, Ca$^{3+}$, Ta$^{5+}$, V$^{5+}$, and Mg$^{2+}$ were employed in doping to substitute a small quantity of Li$^+$ or Ti$^{4+}$ to improve the electronic conductivity and rate capability [8, 11–13]. Coating conductive materials on the surface of Li$_4$Ti$_5$O$_{12}$ can improve surface conductivity and reduce contact resistance [2]. Conductive materials (carbon black, carbon nanotubes, graphene, Ag, Cu, and Sn [6, 10, 14, 15]) have been reported to be used to modify the electric conductivity of Li$_4$Ti$_5$O$_{12}$. 
Many kinds of carbon have been examined as coating source in synthesis of Li$_4$Ti$_5$O$_12$/C composite material with excellent electrochemical performance by conventional methods (e.g., solid state reaction and sol-gel methods) [3, 16, 17]. Starch is not only a good template in synthesis of various function materials with excellent electrochemical performance [18, 19] but also a good carbon source of coating to improve electric conductivity and reduce particle size [20]. Starch-sol-assisted rheological phase method has been developed to synthesize LiFePO$_4$/C composite [20]. However, to the best of our knowledge, synthesis of Li$_4$Ti$_5$O$_12$/C composite has not been reported previously by using this method. In this work, we have used inexpensive starch as carbon source and adopted rheological phase method to prepare high-capacity Li$_4$Ti$_5$O$_12$/C material. The composite was characterized by XRD, SEM, and TG techniques, and galvanostatic charge/discharge and EIS methods were used to examine electrochemical performance. The results show that carbon is cohesively coated on the surface of Li$_4$Ti$_5$O$_12$, and the composite has demonstrated high rate performance with good cycle stability when it was tested as anode in Li-ion batteries.

2. Experimental

2.1. Preparation of Li$_4$Ti$_5$O$_12$ and Li$_4$Ti$_5$O$_12$/C. The starting raw materials, TiO$_2$-anatase (Aldrich), LiOH (Aldrich), and starch, were analytical grade and used without further purification. Firstly, the soluble starch (3.0 g starch/0.02 mol Li$_4$Ti$_5$O$_12$) was mixed with appropriate amount of deionized water; the resultant mixture was heated initially at 110°C until the solution became transparent under stirring in oil bath. Then the starch sol was mixed with stoichiometric amount of LiOH and TiO$_2$ (molar ratio of Ti:Li = 4.2:5) under vigorous agitation to obtain a solid-liquid rheological body, which appears to be a mushy slurry. Finally, the mushy slurry was transferred to tubular furnace and heated at the rate of 15°C min$^{-1}$ to 850°C and sintered at 850°C for 4 h in N$_2$ atmosphere and then cooled to room temperature to obtain a grey Li$_4$Ti$_5$O$_12$/C powder. For comparison, pure Li$_4$Ti$_5$O$_12$ without carbon coating was synthesized by solid state method, which was prepared by mixing TiO$_2$-anatase and LiOH (mole ratio of Li:Ti = 4.2:5); the mixture was heated at 850°C for 4 h in air to obtain well-crystallized white powder Li$_4$Ti$_5$O$_12$.

2.2. Electrode Fabrication. The electrochemical characterization was carried out by galvanostatic charge and discharge using a two-electrode type of 2025 coin cell. The working electrode was prepared with 84 wt.% active material Li$_4$Ti$_5$O$_12$ or Li$_4$Ti$_5$O$_12$/C, 10 wt.% super-P-Li carbon black, and 3 wt.% CMC (carboxymethyl cellulose sodium), and 3 wt.% SBR (styrene butadiene rubber) dissolved into deionized water to form mixed slurry. The slurries were dispersed in a planetary mixer and then spread uniformly on an aluminum foil by using doctor blade. Finally, the laminates were dried under vacuum at 100°C for 24 h before electrochemical evaluation. Electrode disks were punched out of the laminates for the half-cell. The cells were assembled based on the configuration of Li/electrolyte/Li$_4$Ti$_5$O$_12$/C (or Li$_4$Ti$_5$O$_12$) with a liquid electrolyte (1.3 M LiPF$_6$ in a mixture (1:3 by mass) of ethylene carbonate and dimethyl carbonate), metallic Li foil was used as the counter electrode in 2-electrode cell, and Celgard 2320 was used as the separator; the coin cells were assembled in a glove box (Braun, [O$_2$] < 1 ppm, [H$_2$O] < 1 ppm) filled with pure argon.

2.3. Characterization. The samples of Li$_4$Ti$_5$O$_12$/C and Li$_4$Ti$_5$O$_12$ were characterized by X-ray diffraction (XRD) with D/Max-III A instrument using Cu Kα ($\lambda = 1.54056$ Å) as radiation source at a scanning rate 2° min$^{-1}$ for 2θ in the range of 10–80° to identify the structure and phase. The morphologies of the samples were investigated by scanning electron microscopy (SEM, Hitachi S4800) and transmission electron microscopy (TEM, JEOL JEM-2010). Thermogravimetric and differential thermal analyses (TG/TGA) were performed on a TG/TGA instrument (Pyris Diamond) using a heating rate of 5°C/min.

The charge/discharge behavior and rate capacity of the cell were tested at various current densities in the range 1.0–2.0 V using land battery test station at room temperature. The AC impedance spectrum was employed at a constant potential of 1.55 V against Li metal electrode, the voltage amplitude of the AC signal was 5 mV, and the frequencies were scanned from 100 kHz to 100 mHz by using an electrochemical workstation IM6ex (Zahner).

3. Results and Discussion

3.1. Physicochemical Characterization. Figure 1 presents the XRD pattern of the two samples. Figure 1(a) is for the pure Li$_4$Ti$_5$O$_12$ prepared by solid-state method, and Figure 1(b) is for the Li$_4$Ti$_5$O$_12$/C composite synthesized by starch-sol-assisted rheological phase method. According to JCPDS file...
Figure 2: SEM images of pure Li$_4$Ti$_5$O$_{12}$ (a) and Li$_4$Ti$_5$O$_{12}$/C composite (b). HRTEM image and the inset (c) clearly show the carbon coating on the Li$_4$Ti$_5$O$_{12}$ crystallites.

no. 49-0207, all the main peaks are investigated, which correspond to Li$_4$Ti$_5$O$_{12}$ without other impurity phases. The pattern could be indexed to spinel structure with Fd3m space group. The lattice parameter of “a” for Li$_4$Ti$_5$O$_{12}$/C composite is 8.359 Å, which is close to the reported value of the literature [21]. No carbon diffraction peak is observed in Figure 1(b) since the carbon generated from starch is in low content, and it is amorphous, which suggests that the addition of starch in the precursor does not influence the formation of spinel Li$_4$Ti$_5$O$_{12}$ crystal during heat treatment in the N$_2$ atmosphere. It was found through Scherrer equation that the average crystallite size was about 400–600 nm.

For comparison, Figure 2(a) shows the SEM image of pure Li$_4$Ti$_5$O$_{12}$ synthesized by solid-state method. The Li$_4$Ti$_5$O$_{12}$ crystallites have smooth sphere-like shape morphology, large blocks of particles at about 700 nm–1 μm in size. Figure 2(b) presents SEM image of the Li$_4$Ti$_5$O$_{12}$/C composite. It is seen that the carbon uniformly covers the surface of Li$_4$Ti$_5$O$_{12}$ and links together forming a carbon network. The particles of Li$_4$Ti$_5$O$_{12}$/C have a relatively less extent of agglomeration, and the average particle size is about 500 nm, which is consistent with the result from XRD calculation. Figure 2(c) shows the high-resolution TEM image of the Li$_4$Ti$_5$O$_{12}$/C composite; it is clearly seen that the Li$_4$Ti$_5$O$_{12}$ has sphere-like shape and the particle size ranges from 400 nm to 600 nm, consistent with the SEM observation in Figure 2(b). The inset HRTEM also shows a 5 nm thick carbon film coated on the surface of Li$_4$Ti$_5$O$_{12}$ crystallites. Due to the adhesive nature of starch sol, the carbon binds cohesively on the surface of reactants during calcination process, and the heating process further results in decomposition of the starch into a carbonaceous deposit on the surface of Li$_4$Ti$_5$O$_{12}$. The carbonization of starch precursor effectively restrains the growth of Li$_4$Ti$_5$O$_{12}$ crystallites into larger size.

In order to investigate the content of carbon in the composite Li$_4$Ti$_5$O$_{12}$/C, TG/DTA was carried out with 15°C min$^{-1}$ heating rate from 30°C to 1000°C in air. The profile of TG/DTA is given in Figure 3. A broad exothermic peak appears at about 580°C, which is owing to the combustion of the starch into carbonaceous species on the surface of Li$_4$Ti$_5$O$_{12}$, and the TG indicates that the content of carbon in the composite Li$_4$Ti$_5$O$_{12}$/C is about 5%. The weight loss after 700°C is due to the loss of lithium during the high temperature process [22, 23].
3.2. Electrochemical Characterization. The electrochemical performance of Li$_4$Ti$_5$O$_{12}$/C and Li$_4$Ti$_5$O$_{12}$ electrodes was examined in coin cells with metallic lithium as counter electrode. The cycle behavior of the cell was investigated in the range of 1.0–2.0 V. The initial charge-discharge curves of Li$_4$Ti$_5$O$_{12}$ (a) and Li$_4$Ti$_5$O$_{12}$/C (b) at various charge/discharge rates are compared as shown in Figure 4. At the current density of 0.2C, the initial discharge capacity of Li$_4$Ti$_5$O$_{12}$/C is 171.5 mAh g$^{-1}$, which is very close to its theoretical capacity (175 mAh g$^{-1}$), and at a 20C high rate, the capacity is 110 mAh g$^{-1}$. Furthermore, the profile shows very flat discharge/discharge plateau at around 1.55 V (versus. Li/Li$^+$). However, the initial specific capacity of pure Li$_4$Ti$_5$O$_{12}$ is only about 160 mAh g$^{-1}$ at 0.2C and 34.5 mAh g$^{-1}$ at 20C, the cell voltage is lower, and the polarization increases with the increasing of charge/discharge rate, as it is shown that the plateau voltage difference between charge and discharge is 60 mV at 0.5C and 440 mV at 20C.

With carbon-coated Li$_4$Ti$_5$O$_{12}$/C, the polarization was obviously smaller as it is seen that the voltage gap between charge and discharge plateau is narrower as shown in Figure 4(b), and the plateau appears more flat especially at high rates. This improved high-rate discharge ability of Li$_4$Ti$_5$O$_{12}$/C which may be attributed to the facts that (1) the carbon generated from starch was coated uniformly on the surface of Li$_4$Ti$_5$O$_{12}$ resulting in enhanced electrical conductivity and (2) the starch limited the growth of Li$_4$Ti$_5$O$_{12}$ crystallite size during calcination process, which benefits Li$^+$ diffusion with shortened distance in Li$_4$Ti$_5$O$_{12}$.

To investigate the rate capability, the Li$_4$Ti$_5$O$_{12}$ and Li$_4$Ti$_5$O$_{12}$/C have been cycled at various rates from 0.2C to 20C, and the results are presented in Figure 5. The charge and discharge cycles were taken for 5 cycles at each rate. As shown in Figure 5(a), the specific discharge capacity is gradually decreased with increased rates. However, the degree of capacity decay of Li$_4$Ti$_5$O$_{12}$/C sample is in less extent in comparison with the pure Li$_4$Ti$_5$O$_{12}$. At 20C the capacity remains 75% of that at 0.2C for Li$_4$Ti$_5$O$_{12}$/C, but for pure Li$_4$Ti$_5$O$_{12}$ the capacity is about 24% between 20C and 0.2C rates. Figure 5(b) shows the cycle performance of the Li$_4$Ti$_5$O$_{12}$/C electrode at 1C and 20C. The initial specific discharge capacity of Li$_4$Ti$_5$O$_{12}$/C is 168.6 mAh g$^{-1}$ at 1C, and it retains 87% after 500 cycles. At 20C cycle rate, the sample exhibits an initial discharge capacity of 110 mAh g$^{-1}$ in Figure 5(b), and there is no apparent capacity fading in the 1st–1000th cycles at 20C, indicating excellent rate performance of the material. After 2000 cycles, the capacity remains 73% of initial value. Compared to the literature data, the performance of the Li$_4$Ti$_5$O$_{12}$ and Li$_4$Ti$_5$O$_{12}$/C composites synthesized by starch-sol-assisted rheological phase method in this work is superior to those reported in [4, 6, 24–28]. This cycle performance can be ascribed to the fact that the Li$_4$Ti$_5$O$_{12}$ itself has a low volume expansion structure, and the cohesive carbon coating that resulted from the starch sol rheological phase method increased electrical conductivity in Li$_4$Ti$_5$O$_{12}$/C.

EIS technique has been considered as an effective way to identify diffusion phenomena in electronic and ionic conductors. To study conductivities of the Li$_4$Ti$_5$O$_{12}$/C and Li$_4$Ti$_5$O$_{12}$, EIS measurements were carried out at 1.55 V at
room temperature. The Nyquist plots of Li₄Ti₅O₁₂/C and Li₄Ti₅O₁₂ electrodes as well as the fitting results using an equivalent circuit are given in Figure 6. Both Nyquist plots are composed of a depressed semicircle in high frequency and a straight line in low frequency region in Figure 6(a). The solution resistance \( R_s \) of the cell from \( Z' \) axis interception at high frequency includes both electrolyte and electrode contact resistance. The charge transfer resistance \( R_{ct} \) is determined by the semicircle in the middle frequency range, which is mainly related to the electrochemical reaction at the electrolyte/electrode interface. The straight line in low frequency range is attributed to the Warburg impedance \( Z_w \), which is due to the solid-state diffusion of Li⁺ ions into the bulk of active material. It can be seen from Figure 6(a) that the Li₄Ti₅O₁₂/C electrode displayed a much lower impedance than that of pure Li₄Ti₅O₁₂ which is similarly observed in [16, 29].

The equivalent circuit model is depicted in Figure 6(b), where the \( R_s \) and \( R_{ct} \) are solution resistance and charge transfer resistance, respectively. CPE is the constant phase-angle element, involving double layer capacitance of the electrolyte-electrode interface. \( R_f \) is the surface polarization resistance, and the \( C_f \) is the surface capacitance. The equivalent circuit model is well suited to the fitting of experiment data, and the simulation results of the EIS data for the Li₄Ti₅O₁₂/C and Li₄Ti₅O₁₂ electrodes are given in Table 1. It can be seen that the \( R_s \) and \( R_{ct} \) are much smaller for the Li₄Ti₅O₁₂/C electrode \((R_s = 1.92 \, \Omega, \, R_{ct} = 24 \, \Omega)\) than for the pure Li₄Ti₅O₁₂ \((R_s = 3.33 \, \Omega, \, R_{ct} = 107 \, \Omega)\). This result indicates that the carbon coating has enabled easier charge transfer at electrode/electrolyte interface, and the overall battery internal resistance is also decreased. The electrode polarization resistance of Li₄Ti₅O₁₂/C \((R_f = 0.184 \, \Omega)\) is much smaller than that of Li₄Ti₅O₁₂ \((R_f = 1.08 \, \Omega)\), which implies that the carbon reduced the polarization. The results are also consistent with the charge/discharge curves shown in Figure 4. The two values of \( Z_w \) are almost identical; this is reasonable since the Li⁺ diffusion rate in Li₄Ti₅O₁₂ crystal lattice should be independent of carbon coating. The in situ starch sol coating in the synthesis of Li₄Ti₅O₁₂ significantly enhances the conductivity of Li₄Ti₅O₁₂ material. The adhesive conductive carbon coating facilitates electronic conductive paths in the Li₄Ti₅O₁₂ particle surroundings, that is considered a key factor in improving the discharge capacity, rate capability, and cycle life of the Li₄Ti₅O₁₂/C material.

4. Conclusions

With the starch precursor as carbon source, the carbon was cohesively coated on the surface of Li₄Ti₅O₁₂ crystallites in the Li₄Ti₅O₁₂/C composite by using starch-sol-assisted rheological phase method. This synthesis approach is relatively simple and effective to obtain the high electronic conductive material. The electrochemical characterization indicates that the Li₄Ti₅O₁₂/C exhibits high rate capability and long cycle life compared to pure Li₄Ti₅O₁₂ from solid-state method. The Li₄Ti₅O₁₂/C anode shows specific capacity of 168.6 mAh g⁻¹ at 1C and 110 mAh g⁻¹ at 20C, and the discharge capacity retains 73% after 2000 cycles at 20C cycle rate. The excellent rate capability and cycling performance of Li₄Ti₅O₁₂/C are mainly attributed to the cohesive carbon coating that enhances conductivity and to the intrinsic structural stability of Li₄Ti₅O₁₂ that allows reversible Li⁺ intercalation in charge-discharge process.

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Table 1: EIS simulation parameters of Li$_4$Ti$_5$O$_12$ and Li$_4$Ti$_5$O$_12$/C from the equivalent circuit.

| Sample          | $R_s$ (Ω cm$^2$) | $R_i$ (Ω cm$^2$) | CPE (S$^2$ cm$^{-2}$) | $W$ (S$^2$ cm$^{-2}$) | $R_f$ (Ω cm$^2$) | $C_f$ (S$^2$ cm$^{-2}$) |
|-----------------|------------------|------------------|-----------------------|-----------------------|------------------|-----------------------|
| Li$_4$Ti$_5$O$_12$ | 3.33             | 107              | $2.5 \times 10^{-4}$ | $2.55 \times 10^{-4}$ | 1.08             | 8.66 $\times 10^{-6}$ |
| Li$_4$Ti$_5$O$_12$/C | 1.92             | 24               | $2.45 \times 10^{-4}$ | $2.2 \times 10^{-4}$  | 0.184            | 2.72 $\times 10^{-6}$ |

Figure 6: EIS Nyquist plots of the Li$_4$Ti$_5$O$_12$ and Li$_4$Ti$_5$O$_12$/C electrode (a) and the corresponding equivalent circuit model that is used to fit the experimental data (b).

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Research Article

Influence of ITO-Silver Wire Electrode Structure on the Performance of Single-Crystal Silicon Solar Cells

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This study aimed to explore the effect of various electrode forms on single-crystal silicon solar cells by changing their front and back electrode structures. The high light penetration depth of the Indium Tin Oxide (ITO) and the high conductivity of the silver wire that were coated on the single crystal silicon solar cells increased photoelectron export, thus increasing the efficiency of the solar cell. The experiment utilized a sol-gel solution containing phosphorus that was spin coated on single-crystal silicon wafers; this phosphorus also served as a phosphorus diffusion source. A p-n junction was formed after annealing at high temperature, and the substrate was coated with silver wires and ITO films of various structures to produce the electrodes. This study proposed that applying a heat treatment to the aluminum of back electrodes would result in a higher efficiency for single-crystal silicon solar cells, whereas single-crystal silicon solar cells containing front electrodes with ITO film coated with silver wires would result in efficiencies that are higher than those achieved using pure ITO thin-film electrodes.

1. Introduction

Since the industrial revolution, energy has been harnessed from petroleum and coal; however, these two sources are exhaustible. Moreover, large amounts of carbon dioxide emitted from fuel-generating energy have resulted in serious consequences, such as global warming [1]. As a result of countries around the world also beginning to place importance on this issue, the Kyoto Protocol was passed on December 11, 1997. The Protocol requested that 37 industrial countries and European Nations follow the 1990 standards to achieve a 5% reduction in greenhouse gas emissions from 2008 to 2012, and it was the first step in addressing global warming. The Copenhagen Summit in December, 2009 determined the global carbon reduction target for 2012–2017 in an attempt to prevent an irreversibly damaged climate. Presently, clean, renewable energy is the best solution to the global warming issues addressed by these initiatives. Because solar power is currently the cleanest and most feasible solution, several researchers are studying the solar cell for its potential in addressing these problems [2, 3].

Currently, commercial single-crystal silicon solar cells use silver grid electrodes as front electrodes [4, 5], which have a light obscuration rate of approximately 5–9%. Moreover, when electrons itinerate to the surface of n-type semiconductors, they are susceptible to capture by dangling bonds. Silica is most commonly used as the passivation layer and the antireflection layer for single-crystal silicon solar cells [6]. However, because it lacks a conducting path for electrons, it cannot be use for electrodes. Therefore, this study employs ITO with high transmittance and conductivity characteristics [7–10] as the antireflection layer and conducting electrodes for single-crystal silicon solar cells. This study also explores the effect of adding high-conductivity silver wire on single-crystal silicon solar cells.

2. Experiment

2.1. p-n Diode Production

2.1.1. Specimen Pretreatment [11]. The substrate was a silicon wafer (p-type) with the dimensions of $156 \times 156 \text{ mm}^2$ (lattice directions 100, resistivity $0.5–3 \Omega\text{-cm}$, thickness $200 \mu\text{m}$) cut into silicon chips of size $1.5 \times 1.5 \text{ cm}^2$. The following cleaning steps were performed.

First, the silicon
This study used the sol-gel configuration. This study used the sol-gel method for modulating phosphorus chemicals. The solution contained 4.69 g of polyethylene glycol, 18 g of ethanol, 0.563 g of H$_2$O, and 0.555 g of phosphorus pentoxide; the solution was stirred in a magnetic rotator for two hours to ensure that the solution was completely mixed.

2.1.2. **Sol-Gel Configuration**. This study used the sol-gel method for modulating phosphorus chemicals. The solution contained 4.69 g of polyethylene glycol, 18 g of ethanol, 0.563 g of H$_2$O, and 0.555 g of phosphorus pentoxide; the solution was stirred in a magnetic rotator for two hours to ensure that the solution was completely mixed.

2.1.3. **Phosphorus Diffusion Process**. The sol-gel configuration was spin coated onto single-crystal silicon. The spin coating was divided into two phases. The first phase consisted of a rotation speed of 2,000 rpm for a total of 10 seconds and was then followed by a second phase that consisted of a rotation speed of 6,000 rpm for a total of 10 seconds. The silicon was placed on a heating plate, where it underwent two-stage heating. During the first stage, a temperature of 120°C was maintained for 15 minutes to evaporate the solvent; during the second stage, a temperature of 200°C was maintained for 20 minutes. Subsequently, the silicon was placed inside a high-temperature furnace at a temperature of 900°C for 10 minutes so that it could undergo phosphorus diffusion. Upon completion of diffusion, a BOE solution was used to etch the phosphate glass for five minutes.

2.2. **Electrode Production**

2.2.1. **Back Electrode Production**. A magnetron sputter was used to coat the back of single-crystal silicon (p-n structure) with aluminum. The sputtering power was 100 W, the work pressure was 5 × 10$^{-3}$ Torr, and the sputtering time was 30 minutes; these parameters resulted in an aluminum thickness of 500 nm.

2.2.2. **High-Temperature Sintering Process**. The single-crystal silicon was thermally annealed to increase the bond strength between the metal and silicon, to export electrons and to produce back surface field structure. The annealing was performed at 650°C for 20 minutes.

2.2.3. **Preparation of Transparent Conductive Film**. An RF magnetron sputter was used to produce a coat of ITO. The sputtering power was 80 W, the work pressure was 3 × 10$^{-3}$ Torr, and the sputtering time was 20 minutes. These parameters resulted in a 450 nm thick ITO film.

2.2.4. **Transparent Conductive Film Annealing**. Single-crystal silicon covered with an ITO electrode underwent a thermal annealing treatment to increase its light penetration depth and conductivity. The annealing temperature was maintained at 350°C for 30 minutes.

2.2.5. **Silver Wire Coating**. The front of single-crystal silicon was sputtered with silver wire using a magnetic sputter. The sputtering power was 50 W, the work pressure was 1 × 10$^{-3}$ Torr, and the sputtering time was 10 minutes; these parameters produced a silver thickness of 300 nm.

2.3. **Electrode Structure**. This study explored the effect of electrode structure on single-crystal silicon via changing the front and back electrode production processes. The study examined four different electrode structures: A, B, C, and D. Each structure contained a heat-treated aluminum electrode and an aluminum electrode that had not been heat-treated. Structure A used ITO film as the front electrode, structure B used annealed ITO film as the front electrode, structure C used ITO film covered with silver wire as the front electrode, and structure D used annealed ITO film covered with silver wire as the front electrode. A, B, C, and D are solar cell structures, as shown in Figure 1.

3. **Results and Discussion**

3.1. **Analysis of p-n Junction Properties**. The resistance of single-crystal silicon after phosphorus diffusion was approximately 40 Ω/sq. with a mixed concentration of 10$^{19}$ atom/cm$^{-3}$. The n-type silicon work function ($W_S$) was 4.07 eV, and the ITO work function of the experiment ($W_m$) was 3.8 eV, which met the conditions for metal-semiconductor ohmic contact ($W_m < W_S$).

The p-n junction solar cells are diodes that may be used to speculate the characteristics of single-crystal silicon.
solar cells through dark current analysis of various electrode structures. The dark current-voltage curve in conjunction may be used to calculate the diode ideality factor (n value) [12, 13].

The dark current-voltage curve was measured on various electrode structures. Logarithmic coordinates were obtained from the current to obtain Figures 2, 3, 4, and 5 and then to obtain diode ideality factors (n). As shown in structure A in Figure 6, the n value calculated from the aluminum electrode before the heat treatment was 2.33, whereas the diode ideality factor following heat treatment of aluminum electrodes was 2.44. For single-crystal silicon solar cells with structure B, the n value calculated from the aluminum electrode before heat treatment was 1.72; the diode ideality factor after heat treatment of the aluminum electrodes was 2.32. For single-crystal silicon solar cells of structure C, the n value calculated from the aluminum electrode before heat treatment was 1.96; the diode ideality factor after heat treatment of the aluminum electrodes was 2.76. For single-crystal silicon solar cells of structure D, the n value calculated from the aluminum electrode before heat treatment was 2.11; the diode ideality factor after heat treatment of the aluminum electrodes was 2.81.

By performing the aforementioned analyses, we discovered that for the four different electrode structures, the diode ideality factor (n) increased when the aluminum electrodes were heat treated. Such a phenomenon implied that applying heat treatment to an aluminum back electrode would enhance the performance of p-n junction single-crystal silicon solar cells.
### 3.2. Analysis of Optoelectronic Properties

The light source employed by this study was an AM1.5, which has an incident light source intensity of 100 mW/cm². The current-voltage curve for this lighting is shown in Figures 7 and 8.

#### 3.2.1. Effect of Front Electrode Structure on Optoelectronic Properties

This study first probed the optoelectronic properties of single-crystal silicon solar cells with two different structures: a front electrode using pure ITO film (structure A) and a front electrode using ITO film combined with silver wire (structure C). The data were compiled in Figures 9–11 according to Figures 7 and 8. Structure A resulted in an optoelectronic conversion efficiency of 7.98%, open-circuit voltage of 0.53 V, and short-circuit current of 29.748 mA/cm². In comparison, structure C yielded improved optoelectronic properties; the photoelectric conversion efficiency increased to 11.31%, the open-circuit voltage increased to 0.42 V and the short-circuit current increased to 39.214 mA/cm².

Subsequently, structures B and D with heat treatment applied to the ITO were compared. Structure D (covered with silver wire) resulted in an open-circuit voltage (0.47 V), short-circuit current (26.810 mA/cm²) and optoelectronic conversion efficiency (8.40%) that were all higher than the open-circuit voltage (0.33 V), short-circuit current (21.527 mA/cm²), and optoelectronic conversion efficiency (4.22%) that resulted from structure B.

These results show that ITO covered with silver wire promotes rapid export of photoelectrons and substantially improves its efficiency.

This study then examined the impact of ITO heat treatment on single-crystal silicon solar cells by comparing structure A with structure B and structure C with structure
D. The following results were obtained. The ITO film with structure A was used as the front electrode that exhibited higher efficiency; the optoelectronic conversion efficiency was 7.98%, the open-circuit voltage was 0.53 V and the short-circuit current was 29.748 mA/cm². In comparison, the ITO with structure B and a heat treatment of 350°C was used as the front electrode that exhibited higher efficiency; the open-circuit voltage increased to 0.59 V, the short-circuit current decreased to 23.150 mA/cm², and the optoelectronic conversion efficiency rose to 9.72%. The aluminum electrode of structure C produced an open-circuit voltage of 0.33 V, a short-circuit current of 21.527 mA/cm², and an optoelectronic conversion efficiency of 4.22% before the heat treatment was applied. After applying a heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.48 V, the short-circuit current decreased to 18.355 mA/cm², and the optoelectronic conversion efficiency rose to 5.27%. The aluminum electrode of structure D produced an open-circuit voltage of 0.42 V, a short-circuit current of 39.214 mA/cm², and optoelectronic conversion efficiency of 11.31% before the heat treatment was applied. After applying a heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.64 V, the short-circuit current decreased to 29.250 mA/cm², and the optoelectronic conversion efficiency rose to 11.62%. The aluminum electrode of structure C produced an open-circuit voltage of 0.47 V, a short-circuit current of 26.810 mA/cm², and optoelectronic conversion efficiency of 8.40% before the heat treatment was applied. After applying a heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.66 V, the short-circuit current decreased to 21.935 mA/cm², and the optoelectronic conversion efficiency rose to 8.62%. These results show that ITO causes single-crystal silicon solar cells to decrease in short-circuit current after undergoing a heat treatment, whereas the open-circuit voltage and optoelectronic conversion efficiency will both increase.

3.2.2. Effect of Heat Treating on the Aluminum Back Electrode. Figures 9–11 show that the aluminum electrode of structure A produced an open-circuit voltage of 0.53 V, short-circuit current of 29.748 mA/cm² and optoelectronic conversion efficiency of 7.98% before a heat treatment was applied. After applying a heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.59 V, the short-circuit current decreased to 28.664 mA/cm², and the optoelectronic conversion efficiency rose to 9.72%. The aluminum electrode of structure B yielded an open-circuit voltage of 0.33 V, a short-circuit current of 21.527 mA/cm², and an optoelectronic conversion efficiency of 4.22% before the heat treatment was applied. After applying the heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.48 V, the short-circuit current decreased to 18.355 mA/cm², and the optoelectronic conversion efficiency rose to 5.27%. The aluminum electrode of structure C produced an open-circuit voltage of 0.42 V, a short-circuit current of 39.214 mA/cm², and optoelectronic conversion efficiency of 11.31% before the heat treatment was applied. After applying the heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.64 V, the short-circuit current decreased to 29.250 mA/cm², and the optoelectronic conversion efficiency rose to 11.62%. The aluminum electrode of structure D produced an open-circuit voltage of 0.47 V, a short-circuit current of 26.810 mA/cm², and optoelectronic conversion efficiency of 8.40% before the heat treatment was applied. After applying the heat treatment of 650°C to the aluminum electrode, the open-circuit voltage increased to 0.66 V, the short-circuit current decreased to 21.935 mA/cm², and the optoelectronic conversion efficiency rose to 8.62%. These results show that ITO causes single-crystal silicon solar cells to decline in overall efficiency upon undergoing a heat treatment of 350°C.

3.2.3. Effect of Inside Impedance. The equivalent circuit for single-crystal silicon solar cells may be used to derive shunt resistance (Rsh), series resistance (Rs) and reverse saturation current (I0) [14]. Rsh proportionally affects open-circuit voltage, whereas Rs inversely affects the short-circuit current. Figures 10 and 12 show that there is a significant correlation between Rsh and open-circuit voltage. Additionally, as shown in Figure 12, Rsh significantly increased after applying heat treatment to the aluminum electrode. For structure A, Rsh was 299.99 kΩ before applying heat treatment to the aluminum electrode, and Rsh rose to 383.71 kΩ after heat treatment. For structure B, Rsh was 120 kΩ before applying heat treatment to the aluminum electrode, and Rsh rose to 342.85 kΩ after heat treatment. For structure C, Rsh was 209.53 kΩ before applying heat treatment to the aluminum electrode, and Rsh rose to 379.29 kΩ after heat treatment. For structure D, Rsh was 356.73 kΩ before applying heat treatment to the aluminum electrode, and
$R_{sh}$ rose to 478.24 kΩ after heat treatment. These results show that $R_{sh}$ increases after applying heat treatment to the aluminum electrode.

The increase in $R_{sh}$ implied a better insulation between the front electrode and back electrode, which would also correspond to a lower reverse saturation current. Reverse saturation currents for different structures are compiled in Figure 13, which shows that the $I_0$ of structure A was initially 4.43 µA, but then decreased to 2.39 µA after heat treatment. The $I_0$ of structure B was initially 12.52 µA, but then decreased to 6.07 µA after heat treatment. The $I_0$ of structure C was initially 9.76 µA, but then decreased to 3.59 µA after heat treatment. The $I_0$ of structure D was initially 7.31 µA, but then decreased to 2.42 µA after heat treatment.

Figure 14 illustrates the series resistance ($R_s$) for different structures. Figures 11 and 14 show that $R_s$ affects the short-circuit current ($J_{sc}$). After heat treatment, the $R_s$ of structure A rose from 30.12 mΩ to 32.41 mΩ, and $J_{sc}$ decreased from 29.748 mA/cm$^2$ to 28.664 mA/cm$^2$. After heat treatment, the $R_s$ of structure B rose from 27.68 mΩ to 31.83 mΩ, and $J_{sc}$ decreased from 21.527 mA/cm$^2$ to 18.355 mA/cm$^2$. After heat treatment, the $R_s$ of structure C rose from 23.28 mΩ to 25 mΩ, and $J_{sc}$ decreased from 39.214 mA/cm$^2$ to 29.250 mA/cm$^2$. After heat treatment, the $R_s$ of structure D rose from 24.78 mΩ to 29.26 mΩ, and $J_{sc}$ decreased from 26.810 mA/cm$^2$ to 21.935 mA/cm$^2$. These results show that $R_s$ grew, whereas $J_{sc}$ decreased after heat treatment of the aluminum electrode; this is due to the internal electronic transmission of the semiconductor encountering more obstacles, which thereby lowers the output current.

Moreover, in comparing structure A to structure C and structure B to structure C, we discovered that the electrode of ITO covered with silver wire results in a higher short-circuit current ($J_{sc}$) and a lower $R_s$ than those produced by the electrode of pure ITO. This indicates that adding the electrode of silver wire reduces the $R_s$.

4. Conclusion

This study used the sol-gel method to prepare a phosphorus diffusion solution, which was spin coated onto the single-crystal silicon substrate using a 900°C heat diffusion process to form a p-n junction. This substrate was then covered with an ITO layer of variable structure and silver wire as the electrode. The impact of heat treatment on electrode structure was also examined in the study and yielded the following results.

(1) The average efficiency of single-crystal silicon solar cells increased after applying a 650°C heat treatment to the aluminum electrode.
(2) The efficiency decreased after applying a 350°C heat treatment to the ITO for 30 minutes, which resulted mainly from a decreasing shunt resistance and an increasing series resistance. This led to a decrease in both open-circuit voltage and short-circuit current, thereby lowering efficiency.

(3) When comparing the composite electrode of ITO and silver wire with pure ITO, the lower resistance of silver wire substantially enhanced the short-circuit current; however, it also reduced some of the open-circuit voltages. Overall, ITO covered with silver wire increased the optoelectronic efficiency of ITO.

(4) This study showed that the optimal front electrode structure was ITO covered with silver, which resulted in an 11.62% efficiency when a heat treatment of 650°C was applied to the aluminum metal of the back electrode.

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Research Article

High-Efficiency CdS Quantum-Dots-Sensitized Solar Cells with Compressed Nanocrystalline TiO₂ Photoelectrodes

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Nanocrystalline TiO₂ films were fabricated on titanium substrates by compression method. The CdS quantum dots (QDs) were assembled onto the compressed TiO₂ layers, which serve as sensitizers. A maximum power conversion efficiency of 4.49% is achieved under 100 mW/cm² illumination. In this paper, we find that the compression can help increase the efficiency of the cell by increasing the absorption of the CdS QDs and improving the transportation of photogenerated electrons.

1. Introduction

Dye-sensitized solar cells (DSSCs) have established themselves as an alternative to conventional solar cells owing to their remarkably high power conversion efficiency, longtime stability, and low-cost production [1, 2]. The traditional DSSCs use organic dyes as sensitizers and the power conversion efficiency of 11% has been achieved using ruthenium dyes [3]. In addition to organic dyes, QDs [2, 4–7] can also serve as sensitizers of DSSCs. The use of semiconductor QDs as sensitizer has some advantages compared with the conventional systems. Firstly, the band gap of the nanocrystals can be adjusted by changing their size so that the absorption spectrum can be tuned to match the spectral distribution of sunlight [8]. Secondly, these QDs can use hot electrons or generate multiple charge carriers with a single photon [7].

In the past, QDs-sensitized solar cells usually use a sintered TiO₂ film as the photoelectrode. In such a device, a lot of energy and time is consumed in the fabrication of TiO₂ films; moreover, the power conversion efficiency (η) is mainly limited by the small amount of the QDs on the electrode surface and the electron recombination in the cell [2, 4]. In this study, the TiO₂ film was fabricated using the compression method developed by Anders Hagfeldt and coworkers [9–11]. This method can offer a simple and fast way to fabricate TiO₂ films; moreover, it was proved that the resulting films are porous and mechanically stable [9, 11]. We expect that this method would help resolve the problems mentioned above.

2. Experiment

In this paper, the TiO₂ paste used was composed of a mixture of commercially available nanosized P25 and ethanol solvent at a concentration of 20 wt%. After ultrasonication, the suspension was applied onto the Ti sheet by hand using doctor blading. Then, pressure (1500 kg/cm²) treatment was applied. CdS QDs were deposited to the TiO₂ film by S-CBD method [5]. A TiO₂ film was dipped into a CdCl₂ water solution (0.05 M) for 30 s, rinsed with water, and dipped for another 30 s into a Na₂S water solution (0.05 M). The two-step dipping procedure is termed as one S-CBD cycle and the incorporated amount of CdS can be increased by repeating the assembly cycles.

The performances of the photoelectrochemical (PEC) solar cells were studied in a three-armed cell with a platinum filament counter electrode and a saturated calomel electrode (SCE) as a reference. 0.1 M Na₂S solution serves as the electrolyte. Measurements were carried out under AM 1.5 G at 100 mW/cm² illumination.
3. Results and Discussion

Figure 1 shows the process and mechanism of this QDs-sensitized solar cell. A suitable pressure can make a well packing TiO$_2$ film; after a series of S-CBD cycles some CdS QDs can formed on the surface of TiO$_2$ film and in the gaps between TiO$_2$ particles. When it illuminates under the light, the CdS QDs can absorb the light energy and an electron from the molecular ground state $S^0$ is excited to an excited state $S^*$. The excited electron of the CdS QDs is injected into the conduction band of the TiO$_2$ particles, leaving the QDs molecule to an oxidized state $S^+$.

Figure 2 shows UV-vis absorption spectra of a bare TiO$_2$ film and TiO$_2$ films sensitized by CdS QDs with different S-CBD times. From Figure 2, we can see with the S-CBD cycles increasing, the UV-vis absorption shoulder and onset position become redshift, which indicating more visible light can be absorbed. By using the empirical equation proposed by Yu et al. [12], it is possible to estimate the sizes of CdS particles from the excitonic peaks of the absorption spectra. The mean diameter of CdS particles (10 S-CBD) was measured to be ca. 4.10 nm on the TiO$_2$ films, which indicating that the size of the CdS particles on the TiO$_2$ films is still within the scale of QDs.

To investigate the effect of compress and S-CBD cycles, different cells were made under the same condition. The most variation is the photoanode that we list in Table 1.

Figure 3 shows the photocurrent-voltage ($J$-$V$) characteristic curves of PEC solar cells as a function of S-CBD cycles. As seen in Figure 3, $J$-$V$ measurements were performed under the conditions of illumination and darkness, respectively. A summary of $J$-$V$ characteristics and the $\eta$ of cells with different photoanodes are presented in Table 1. Under illumination, it is found that $\eta$ can be markedly enhanced by the absorption of the CdS QDs. With the increase of S-CBD cycles to 10, generated photocurrent ($J_{sc}$) and open-circuit photovoltage ($V_{oc}$) are both enhanced, reaching their maximum of 5.91 mA/cm$^2$ and 1.17 V (versus SCE), respectively. An efficiency of 4.49% is achieved in such a cell. As the S-CBD cycles increases above 10, $V_{oc}$ and $J_{sc}$ both decrease. This trend can be explained as follows: when the amount of the CdS QDs due to the S-CBD cycles is too small, the absorbance of dye is insufficiently, which causes the low $J_{sc}$. When the amount of the CdS QDs is too large, the transmission process of the electrolyte inside the TiO$_2$ film’s pore is restricted, which causes the serious electron recombination of the cell and also reduces cell performance.
Table 1: Parameters obtained from the photocurrent-voltage ($J-V$) measurements of the PEC solar cells constructed using various electrodes.

| Cell  | Photoanode                      | $J_{sc}$ (mA/cm²) | $V_{oc}$ (V versus SCE) | $ff$   | $\eta$ (%) |
|-------|--------------------------------|-------------------|-------------------------|--------|-------------|
| Cell 1| Sample 1 (Ti/TiO₂ (pressure)/S-CBD10) | 5.91              | 1.17                    | 0.65   | 4.49        |
| Cell 2| Sample 2 (Ti/TiO₂ (pressure)/S-CBD5)  | 3.63              | 1.19                    | 0.68   | 2.94        |
| Cell 3| Sample 3 (Ti/TiO₂ (pressure)/S-CBD15) | 4.99              | 1.14                    | 0.60   | 3.41        |
| Cell 4| Sample 4 (Ti/TiO₂ (pressure)/S-CBD0)  | 0.16              | 0.87                    | 0.61   | 0.08        |
| Cell 5| Sample 5 (Ti/TiO₂ (heat after pressure)/S-CBD10) | 1.08              | 1.09                    | 0.58   | 0.68        |
| Cell 6| Sample 6 (Ti/TiO₂ (heat)/S-CBD10)   | 0.55              | 1.16                    | 0.70   | 0.44        |

Figure 4: SEM of (a) Ti/TiO₂ (compressed), (b) Ti/TiO₂ (heat treatment after compressed), and (c) Ti/TiO₂ (sintered).

Figure 5: EDS of (a) Ti/TiO₂ (compressed)/CdS(S-CBD10) and (b) Ti/TiO₂ (heat treatment after compressed)/CdS(S-CBD10). The inset lists their dates, respectively.
We also took $J$-$V$ curves under the dark condition (Figure 3(b)). Figure 3(b) displays that the onset of the dark current of cell 1 occurs at higher bias than cell 4, which indicates the electronic recombination in cell 1 is lower than cell 4.

Table 1 also lists the date of cell 5 and cell 6. The configuration of cell 5 and cell 6 is similar to cell 1; the only difference between them is the fabrication technical of TiO$_2$ films. The compressed TiO$_2$ film of cell 5 was posttreated by a heat treatment (450°C, 1.5 h) and the TiO$_2$ film of cell 6 was made by sintered (450°C, 1.5 h) directly. Surprisingly, both cell 5 and cell 6 show bad results compared with cell 1. Especially, the cell 5 (with $J_{sc}$ of 1.08 mA/cm$^2$, $V_{oc}$ of 1.09 V versus SCE, $ff$ of 0.58, and $\eta$ of 0.68%) performs worse than the cell 1 with respect to all cell parameters, especially in $J_{sc}$ response, the cell 1 is 5.5 times that of the cell 5. It seems that the additional heat treatment makes no improvement in cell performance. As mentioned in the introduction, the power conversion efficiency is mainly limited by the small amount of the QDs on the electrode surface when using a sintered TiO$_2$ film as the photoelectrode. Here, on the basis of the results presented above, we assume that the compression method may help solve this problem. To verify it, we first conducted scanning electron microscopy (SEM) studies of their morphology.

Figure 4 shows plane-view SEM images of a blank TiO$_2$ film of cell 1 (Figure 4(a)), cell 5 (Figure 4(b)), and cell 6 (Figure 4(c)). All of films display porous structures, creating holes in the film. However, big holes and aggregation between particles can be seen more obviously in Figures 2(b) and 2(c), implying that the surface area per unit volume of the sintered TiO$_2$ film is smaller than that of unsintered TiO$_2$ film. This is consistent with Lindström’s observation [11]. Kavan et al. [13] find the sintering can cause decrease in surface areas of coated TiO$_2$ films. They think this decrease is caused by the efficient filling of small pores in the aggregate by surface diffusion. At the same time, we assume that the heat treatment may induce some chemical changes, which may weaken the linkage between the TiO$_2$ particles and the CdS QDs (further researches are expected to certify it). And both of the phenomena we mentioned above limit the absorption of CdS QDs in cell 5 and cell 6, and thus decrease the $J_{sc}$. Support for this explanation is given by the EDS (Energy Dispersive X-ray Spectrometer) analysis. Figure 5 is the EDS spectrum of the cell 1 and cell 5, their values are listed inside the Figure 5. Although the EDS value cannot be considered very accurate, it offers us a simple and direct way to study the amount of the CdS QDs. The EDS results reveal the atomic ratio Cd/O of cell 1 and cell 5 is 7.4% and 2.3%, respectively. We can also change them into the volume ratio (CdS over TiO$_2$) of 22.8% and 7.1%, respectively; the details of the procedure can be found in the literature [5]. So we can find that the amount of CdS QDs of cell 1 is higher than that of cell 5. Noticeably, the above results verify our former explanation. In addition, as seen in Figure 4, the stacking of particles is enhanced by the pressure, which is conducive to the transportation of photogenerated electrons, and can also improve the cell performance.

At last, 4.49% is similar to that of the same structure solar cell (4.15%) [14]. In that literature, they use TiO$_2$ nanotubes as the photoanode. They think the crystalline nature of the nanotubes and the film geometry allows a fast and efficient transfer of the photogenerated electrons from CdS QDs to the Ti substrate. In our device, we have a large fine TiO$_2$ film which has a large surface area and the compress also can give a good cohesion between TiO$_2$ film and Ti sheet which can also improve the transportation of photogenerate electrons. There are also some important factors can affect our final result, as pressure, film’s height, and so on. This time we just use a fixed value to assess our devices. And we will make further researches to study it.

In summary, high efficiency has been obtained for QDs-sensitized TiO$_2$ PEC solar cells prepared by pressing (up to 4.49%). Through the comparison, we conclude that the compression can increase the absorption of CdS QDs and improve the transportation of photogenerate electrons. Overall, the compression method shows promise as a fast, simple, and effective method for the preparation of photoelectrodes in QDs-sensitized solar cells.

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Research Article
Nanoporous ZnO Photoelectrode for Dye-Sensitized Solar Cell

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Nanoporous and macroporous structures were prepared by using self-assembled monolayer (SAM) onto ZnO thin films in order to investigate the efficiency of dye-sensitized solar cells (DSSCs) produced using these films. Using SAM on ZnO thin films, it is obtained successfully assembled large-area, highly ordered porous ZnO thin films. Varying nanoporous radius is observed between 20 and 50 nm sizes, while it is 500–800 nm for macroporous radius. The solar conversion efficiency of 2.75% and IPCE of 29% was obtained using ZnO nanoporous/N719 dye/I−/I3− electrolyte, while macroporous ZnO given solar conversion efficiency of 2.22% and IPCE of 18%.

1. Introduction
The fundamental properties of II–VI compound semiconductor ZnO have been studied for many years, owing to its direct wide band gap (3.37 eV), large exciton binding energy, huge magnetooptical effect, chemical sensing, piezoelectric and ferroelectric properties, low toxicity, high infrared reflectivity, acoustic characteristics, high electrochemical stability, and excellent electronic properties [1–3]. Moreover, in the past few years, interesting much research has been reported connected with ZnO nanoscale structures with various morphologies such as nanoporous, nanowires, nanorods, and nanotubes obtained by different methods [4, 5]. Among these, porous ZnO nanostructures have some specific advantages such as a high specific surface area, chemical and photochemical stability, uniformity in pore size, shape-selective and rich surface chemistry [6]. According to the variation of porosity, physical and chemical properties of the samples can be controlled and changed. This makes porous structures a promising material in the field of separation, sensors, catalysis, bioscience, surface acoustic wave device, life science, photonic crystals, light emitting diodes, lasers application, and dye-sensitized solar cells (DSSCs) [7]. Since DSSCs based on nanostructure materials offer very low cost and relatively efficient photovoltaic energy conversion, it has attracted much attention in the last decade [8–12]. The highest efficiency, to the best of our knowledge, reported on DSSCs was 11% obtained on nanoporous TiO2 by using ruthenium complex dye, containing I−/I3− redox couple electrolytes, and platinum counterelectrode [13–15]. On the other hand, it has been achieved 0.3–6% conversion efficiency in ZnO nanosemiconductor materials [16, 17]. It is still thought of as an alternative to TiO2 because of its ease of crystallization and optical and structural quality.

Reports are available in the preparation of porous ZnO films by a variety of techniques including wet-chemical method, radio frequency magnetron sputtering deposition, electrochemical anodisation, etching surface by an organic and inorganic solvents, templating methods, and self-assembled monolayers (SAMs) which are ordered molecular assemblies of single layer formed via the adsorption of an active surfactant onto the surfaces [18–22]. One of the most common examples to the SAMs is alkene thiols such as hexanethiols and dodecanethiols. The SAMs of thiols have been studied widely in recent years, since they offer a rational and easy approach for fabricating interfaces with a well-designed composition, structural properties and thickness of the samples [23–25].

In this study, macro- and nanoporous ZnO structures were formed by using SAMs and DSSC performance is compared to each other in terms of efficiency. Detailed structural and optical characterizations were performed by using scanning electron microscopy (SEM), X-ray diffraction
(XRD), photoluminescence (PL), absorbance, and Raman spectroscopy.

2. Experimental

The ECD of ZnO thin films were performed on indium tin oxide (ITO) substrate in a solution of 0.05 M Zn(ClO$_4$)$_2$ and 0.1 M LiClO$_4$ into DMSO solvent. ITO substrates were cleaned in an ultrasonic bath with trichloroethylene, acetone, and methanol. A conventional three-electrode ECD system was used for deposition of ZnO thin films. ITO substrate was used as the working electrodes, Ag/AgCl as the reference electrode, and zinc plate as counter electrode. The growth temperature and deposition time were arranged at 100°C and 1 h, respectively. Deposition potential was kept constant during the growth process as $-1.1$ V which is determined from cyclic voltammogram (figure not shown). In order to obtain porous ZnO thin films, 0.05 M C$_6$SH-SAM and 0.05 M C$_{12}$SH-SAM were used with ethanol solution. After the growth of ZnO thin films, thin films were immersed in sulfuric acid (H$_2$SO$_4$) to obtain porous ZnO thin films, 0.05 M C$_{12}$SH-SAM and 0.05 M C$_6$SH-SAM were used with ethanol solution. After the growth of ZnO thin films, thin films were immersed in these solutions.

ZnO nano- and macroporous structure-based DSSCs were prepared by adsorbing cis-bis(isothiocyanato) bis(2,2-bipyridyl-4,40-dicarboxylato)-ruthenium (II) bis tetra butylammonium (N719) dye onto the surfaces of both kinds of samples. Substrates were first heated to 100°C before immersing them into 0.5 mM ethanolic dye solution for 30 min. The counter electrode was an Pt:F:SnO$_2$ substrate. The electrodes were separated by 20 nm polypropylene spacer and pressed together with binder clips. Electrolyte was introduced between the electrodes by capillary forces. The electrolyte consisted of 0.5 M tetrabutylammonium iodide, 0.05 M I$_2$ (Iodine), and 0.5 M 4-tertbutylpyridine in acetonitrile. Active electrode area was typically 0.25 cm$^2$.

The morphology of the nanostructures was analyzed using a JEOL-6400 SEM. The crystal structure was analyzed by XRD (Rigaku D/Max-IIIC diffractometer) with Cu-Ka radiation of 1.54 Å, within the 2θ angle ranging from 20 to 80. The PL measurements were conducted with the RF 5301 PC Shimadzu spectrofluorometer at room temperature. The Raman scattering measurements were performed using a micro-Raman Renishaw 2000 system with an excitation source of 514.5 nm at room temperature. Current-voltage (I-V) characteristics were recorded using a Keithley 175A digital multimeter using 0.01 V/s voltage ramp rate. The light source was a 250 W tungsten halogen lamp calibrated to 100 mW/cm$^2$ using a radiometer (IL1700, International). The absorbance and incident photon to current conversion efficiency (IPCE) were measured using a 1000 W Xe lamp and a monochromator.

3. Results and Discussion

3.1. Formation Mechanism of the Porous ZnO Structures by SAMs of C$_6$SH and C$_{12}$SH. SAMs are formed by immersing an appropriate substrate into a solution of an organic compound (surface-active material), possessing the ability to spontaneously form an ordered molecular layer on the substrate [26]. The driving force for the spontaneous formation of the two-dimensional assembly includes chemical bond formation of molecules with the surface and intermolecular interactions [27]. In this study, formation of different size of porous structures on ZnO semiconductors by SAMs of C$_6$SH and C$_{12}$SH was achieved. In detail, the constituting self-assembling of the C$_6$SH and C$_{12}$SH consists of three parts. The first part is the head group. It causes the exothermic process of chemisorption on the surface of the substrate. It has specific affinity for the substrate. Organization of C$_6$SH and C$_{12}$SH starts with the interaction between the head and the substrate by means of chemisorption and last to the thermodynamic equilibrium. The second part is the alkyl chain. It is responsible for the intermolecular distance, the molecular orientation, and the degree of order in the film. The third part is a functional group that constitutes the outer surface of the film. The principal driving force for the formation of these films is a specific interaction between the head group and the substrate surface [28]. Provided these interactions are strong, SAMs of C$_6$SH and C$_{12}$SH form stable films depending on waiting time into the C$_6$SH and C$_{12}$SH solution. Porous structures develop on ZnO due to solvent evaporation at room temperature. Ma and Hao indicated macroporous structures on gold using C$_{12}$SH solution due to solvent evaporation and water droplet condensation [29]. They showed that the rapid evaporation of the solvent decreases the air/liquid interfacial temperature, resulting in porous formation on substrate. Similarly, Cai et al. showed porous formation via SAMs and indicated that different solvent evaporation rates affect pore formation and pore size [30].

The structural properties of the porous ZnO, which is obtained via SAMs of C$_6$SH and C$_{12}$SH, are strongly influenced by the interactions between the functional groups comprising the alkyl chain and functional group of the molecules as it has been shown our previous study [25]. The porous formation is largely governed by molecular-substrate interactions, though lateral chain-chain interactions are also quite important. While porous formation on ZnO substrate occurs via SAMs of C$_6$SH and C$_{12}$SH, it was shown that films formed by short alkyl chains exhibit very different characteristics than those formed by long alkyl chains. The alkyl chain and functional groups are found to have a profound influence on the porous structure. This knowledge will undoubtedly prove useful in controlling the surface densities of the samples which are prepared via SAMs.

3.2. Structural and Optical Properties of the Porous ZnO. X-ray diffraction measurement of nano- and macroporous ZnO has been shown that the films have highly preferred $c$-axis orientation as shown in Figures 1(a) and 1(b), respectively. Another peak, ZnO (0004), can be seen in the Figure 1(a). The peak intensity of (0002) orientation of macroporous ZnO has shown the decrease by half that of the nanoporous sample. This might indicate that the bigger size
porous radius has a negative effect on the crystallization as expected.

Figures 1(c) and 1(d) show SEM images of the porous structures formed onto the ZnO thin films which were obtained by using SAMs technique. Figure 1(c) shows that the nanoporous ZnO thin films have successfully been obtained with porous radius around 50 nm as indicated on image and also uniform pore distribution can be seen through the surface. Average radius of the pores has been calculated by taking account of the pores having a maximum, minimum and average size as indicated in the Figure 1(c). Macroporous structures are shown in Figure 1(b). Perfect circular structures of the pores can be seen with an average pore radius of around ∼800 nm as it is indicated in the figure.

PL characterizations of the nano- and macroporous ZnO thin films are shown as an inset of the respective SEM images in Figures 1(c) and 1(d). PL measurements exhibit different spectra for both samples. Two luminescence bands are observed in nanoporous ZnO thin films. One of them is relatively weak deep level emission peaking at the 466 nm, while narrow UV emission is dominant emission peaking at about 362 nm which might be responsible for the recombination of the free excitons. As can be seen in the Figure 1(d), macroporous structure has a strong green emission peak located at ∼530 nm. This peak intensity is much higher compared to the intensity of the UV emission at ∼363 nm confirming that the decrease in the optical quality of the film [23–25]. Also, it can be speculated that the deep level defect formation causing the green emission in ZnO, which is believed commonly either oxygen vacancy or zinc vacancy [3], is more favorable in macroporous sample than the nanoporous sample.

The UV-visible absorbance and transmittance spectra of the ZnO nanoporous and macroporous are shown in Figures 2(a) and 2(b), respectively. The absorption spectra of the ZnO porous structures showed a strong absorption between 300 nm and 360 nm. The ZnO nanoporous sample has a good transparency with a visible light transmittance higher than 85%. However, the ZnO macroporous structures exhibited low transmittance at around 60%, which might be related to the pore density. These results clearly indicate that transparency of the ZnO nanoporous structure are related to surface morphology and roughness which affect the light scattering. Optical band gap of the ZnO nanoporous structures is calculated by using the following equation [30, 31]:

\[ \alpha \propto (h\nu - E_g)^{1/2}, \tag{1} \]

where \(\alpha\) is absorption coefficient. The direct band gap \(E_g\) was determined from \(\alpha^2\)-photon energy plot as shown in
Figure 2: (a) Absorbance. (b) Transmittance spectrum of ZnO nanoporous and macroporous structures at room temperature. (c) Absorption coefficient and energy plot for ZnO nanoporous. (d) Absorption coefficient and energy plot for ZnO macroporous.

Figure 3: Raman shift of nano- and macroporous ZnO.
A dominant and strong intensity peak at 438 cm\(^{-1}\) indicates the spectrum known as the optical phonon \(E_2\). The \(E_2\) mode corresponds to the band characteristic for the Wurtzite hexagonal phase of ZnO raman active branches, which is also one of the characteristics of ZnO nanostructures [34]. However, two short peaks at 331 and 380 cm\(^{-1}\) were shown to be as \(E_{2H}-E_{2L}\) (multiphonon) and \(A_{1T}\) modes, respectively. For the ZnO nanostructures, an extra Raman band at 582 cm\(^{-1}\) known to be related to the \(E_1\) mode because of the oxygen deficiency [35] indicates the presence of oxygen vacancies in the ZnO nanostructures.

Figure 4 shows the comparison of the current density-voltage \((J-V)\) characteristics of solar cells fabricated using the ZnO nano- and macroporous structures, under 100 mW/cm\(^2\) of AM-1.5 illumination. It can be seen from Figure 4 that there is a considerable increase in the short-circuit current density \((J_{sc})\) of the nanoporous ZnO DSSCs in comparison with the macroporous ZnO DSSCs. ZnO nanoporous structures provide both a larger surface area and a direct pathway for electron transport along the pore resulting in the observed higher \(J_{sc}\) [36, 37]. This is attributed to the superior light absorbing characteristics provided by the larger surface areas formed by the nanoporous structure.

The open-circuit voltage \((V_{oc})\) of the DSSCs shows that the higher solar conversion efficiency \((\eta)\) of 2.75% is reached for the nanoporous structures, while it is 2.22% for macroporous structure. A comparison of the different solar cell parameters for the both DSSCs is summarized in Table 1. It can be seen that the ZnO nanoporous structures shows \(\eta\) values that are almost 40% higher than that of the macroporous ZnO. This is a significant result showing the dependence of solar cell performance on surface morphology of the ZnO. Obtained larger surface area in nanoporous structure also results in improved dye adsorption [38-40]. Figure 5 displays the dependence of IPCE of the DSSCs on the different ZnO nanostructures. As expected, the ZnO nanoporous DSSCs show much higher IPCE, 29% at about 510 nm, compared to the macroporous ZnO, 18% at about 520 nm, mainly due to the aforementioned effect of larger surface area (and dye-loading).

Several attempts have been made to use ZnO photoelectrode in DSSCs. Redmond et al. [41] achieved an IPCE of 13% at 520 nm when using the Ru(dcbpyH2)2(NCS)2 dye. On the other hand, Keis et al. introduced IPCE values of 50–60% at 540 nm in photoelectrochemical (PEC) and an overall solar energy conversion efficiency of 2.0% was obtained under 56 mW/cm\(^2\) illumination with a solar simulator [42, 43] which is lower than the present study. Recently, solar energy conversion efficiency under 99 mW/cm\(^2\) illumination reaching 2.5% has been reported for then mercurochrome-sensitized ZnO PEC solar cell [44]. However, the efficiency is still moderate compared to solar cells based on TiO\(_2\).

### 4. Conclusion

In this study, using SAMs technique on ZnO thin films, we have successfully formed an assembled large-area, highly ordered porous ZnO thin films. Average pore radius is obtained around 50 nm, and 800 nm for nanoporous and macroporous ZnO structures, respectively. PL measurements show highly different emission characteristics according to the size of the porous radius: the samples having a macroporous show that green emission is the dominant, while, in the nanoporous sample the UV emission is dominant showing to better optical quality of the sample. In addition, the origin of an extra Raman band at 582 cm\(^{-1}\) was obtained due to oxygen deficiency of the ZnO nanostructures. The solar-to-electric energy conversion efficiency of 2.75% and IPCE of 29% was obtained by using the ZnO nanoporous/N719 dye/I\(_3^-/I_3^-\) electrolyte. In addition, Zn macroporous showed solar-to-electric energy conversion efficiency of 2.22% and IPCE of 18%. In conclusion, ZnO nanoporous electrodes have been studied into dye-sensitized solar cells where they
display reasonable light-harvesting efficiency, photovoltage, and relatively good fill factors than macroporous sample in the present study.

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| Morphology       | FF  | Efficiency | V<sub>oc</sub> | I<sub>s</sub>  | P<sub>max</sub> | V<sub>mp</sub> | I<sub>mp</sub> | Area (cm<sup>2</sup>) |
|------------------|-----|------------|--------------|-------------|--------------|--------------|-------------|---------------------|
| Nanoporous ZnO  | 0.494 | 2.750       | 0.745        | 0.672       | 0.247        | 0.520        | 0.476      | 0.25                |
| Macroporous ZnO | 0.796 | 2.222       | 0.748        | 0.336       | 0.200        | 0.580        | 0.345      | 0.25                |
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Research Article

PtRu/C Electrocatalysts Prepared Using Gamma and Electron Beam Irradiation for Methanol Electrooxidation

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PtRu/C electrocatalysts (carbon-supported PtRu nanoparticles) were prepared in a single step submitting water/2-propanol mixtures containing Pt(IV) and Ru(III) ions and the carbon support to gamma and electron beam irradiation. The electrocatalysts were characterized by energy dispersive X-ray analysis (EDX), X-ray diffraction (XRD), transmission electron microscopy (TEM), and cyclic voltammetry and tested for methanol electrooxidation. PtRu/C electrocatalyst can be prepared in few minutes using high dose rate electron beam irradiation while using low dose rate gamma irradiation some hours were necessary to prepare it. The obtained materials showed the face-centered cubic (fcc) structure of Pt and Pt alloys with average nanoparticle sizes of around 3 nm. The material prepared using electron beam irradiation was more active for methanol electrooxidation than the material prepared using gamma irradiation.

1. Introduction

Fuel cells convert chemical energy directly into electrical energy with high efficiency. However, the use of hydrogen as a fuel presents problems, principally with storage for mobile and portable applications [1–3]. Thus, there has been an increasing interest in the use of alcohols directly as fuel (Direct Alcohol Fuel Cell—DAFC). Methanol has been considered the most promising alcohol and carbon-supported PtRu nanoparticles (PtRu/C), normally with a Pt : Ru atomic ratio of 50 : 50, the best electrocatalyst [4]. However, the catalytic activity of PtRu/C electrocatalysts strongly depends on the method of preparation and it is one of the major topics studied in direct methanol fuel cells (DMFC) [4, 5]. PtRu/C electrocatalysts are produced mainly by impregnation and colloids methods. Although impregnation method is a simple procedure, the major drawback is the difficulty in controlling nanoparticle size and distribution. The colloidal methods have the advantage to produce very small and homogeneously distributed carbon-supported metal nanoparticles; however, the methodologies are very complex [4]. Lately radiation-induced reduction of metal ion precursors in solution has been described to prepare carbon-supported metal nanoparticles for fuel cell applications. Despite the complexity and cost of electron beam or gamma irradiation facilities, the methodologies used to prepare the electrocatalysts are easy to perform [6–10]. Le Gratiet et al. [6] prepared platinum nanoparticles submitting a K2PtCl4 salt dissolved in a CO-saturated water/2-propanol solvent to gamma irradiation. The reduction of platinum ions occurred by a combined effect of CO and radicals produced by radiolysis, leading to the formation of platinum nanoparticles of 2-3 nm diameter that were further impregnated on the carbon support. These catalysts were found to be effective for methanol or hydrogen electrooxidation. Oh et al. [7] prepared Pt-Ru alloy particles dispersed on various carbon structures in water/2-propanol using gamma irradiation, but no tests for DMFC were described using the obtained materials. Wang et al. [8] prepared Pt nanoparticles irradiating an aqueous solution of chloroplatinic acid in the presence of 2-propanol as a radical scavenger and sodium sulfonate as...
a surfactant. The synthesized Pt nanoparticles (2.5–4.0 nm) were further impregnated on multiwalled carbon nanotubes. The obtained material was tested on a single proton exchange membrane fuel cell operating with H2/O2, and the results showed that the electrocatalysts were very promising. Silva et al. [9] prepared PtRu/C electrocatalysts in a single step submitting water/ethylene glycol solutions containing Pt(IV) and Ru(III) ions and the carbon support to gamma irradiation at room temperature under stirring. The obtained carbon-supported PtRu nanoparticles showed mean particle sizes of 2.5–3.0 nm and were very active for methanol oxidation. Recently, Chai et al. [10] prepared Pt (80 wt%) supported on a mesoporous carbon support in a single step. The Pt salt was dissolved in a solution of water/2-propanol, and the carbon support was added to the solution. The mixture was irradiated at room temperature under stirring. The obtained material exhibited enhanced catalytic activity towards the oxygen reduction reaction (ORR). In this work, PtRu/C electrocatalysts were prepared using high dose rate electron beam and low dose rate gamma irradiation and were tested for methanol electrooxidation.

2. Experimental

PtRu/C electrocatalyst (20 wt%, Pt:Ru atomic ratio of 50:50) was prepared using H2PtCl6·6H2O (Aldrich) and RuCl3·1.5H2O (Aldrich) as metal sources, which were dissolved in water/2-propanol solution (25/75, v/v). After this, the carbon Vulcan XC72R, used as a support, was dispersed in water/2-propanol solution (25/75, v/v). After this, the mixture was dissolved in a solution of water/2-propanol, and the carbon support was added to the solution. The mixture was dispersed in the solution using an ultrasonic bath. The resulting mixture (dissolved metal ions and the carbon support) was treated in an ultrasound bath for 10 min, filtered, and transferred to the cavity (0.30 mm deep and 0.36 cm² area) of the working electrode. The quantity of electrocatalyst in the working electrode was determined with an accuracy of 0.0001 g using an analytical balance. In cyclic voltammetry experiments the current values (I) were normalized per gram of platinum (A g⁻¹). The quantity of platinum was calculated considering the mass of the electrocatalyst present in the working electrode multiplied by its percentage of platinum. The reference electrode was a RHE, and the counter electrode was a Pt plate. Electrochemical measurements were made using a Microquimica (model MQPG01, Brazil) potentiostat/galvanostat coupled to a personal computer and using the Microquimica software. Cyclic voltammetry was performed in a 0.5 mol L⁻¹ H2SO4 solution saturated with N2. Methanol oxidation was performed at 25°C using 1.0 mol L⁻¹ of methanol in 0.5 mol L⁻¹ H2SO4. For comparative purposes, a commercial PtRu/C E-TEK (20 wt%, Pt:Ru molar ratio 50:50, Lot # B0011117) was used.

3. Results and Discussion

Electron beam or gamma irradiation of a water solution containing metals causes the ionization and excitation of water, forming the species shown in (1) [14].

$$\text{H}_2\text{O} \rightarrow e_{aq}^{-}, \text{H}^{+}, \text{H}^{+}, \text{OH}^{+}, \text{H}_2\text{O}_2, \text{H}_2. \quad (1)$$

The aqueous solvated electrons, $e_{aq}^{-}$, and $\text{H}^{+}$ atoms are strong reducing agents and reduce metal ions down to the zero-valent state ((2) and (3)).

$$\text{M}^{+} + e_{aq}^{-} \rightarrow \text{M}^0, \quad (2)$$

$$\text{M}^{+} + \text{H}^{+} \rightarrow \text{M}^0 + \text{H}^{+}. \quad (3)$$

Similarly, multivalent ions, like Pt(IV) and Ru(III), are reduced by multistep reactions. However, $\text{OH}^{+}$ radicals could oxidize the ions or the atoms into a higher oxidation state and thus counterbalance the reduction reactions, (2) and (3). An $\text{OH}^{+}$ radical scavenger is therefore added to the solution, in this case 2-propanol, which reacts with these radicals leading to the formation of radicals exhibiting reducing power that are able to reduce metal ions ((4) and (5)) [14].

$$(\text{CH}_3)_2\text{CHOH} + \text{OH}^{+} \rightarrow (\text{CH}_3)_2\text{COH} + \text{H}_2\text{O}, \quad (4)$$

$$\text{M}^{+} + (\text{CH}_3)_2\text{COH} \rightarrow \text{M}^0 + (\text{CH}_3)_2\text{CO} + \text{H}^{+}. \quad (5)$$

In this manner, the atoms produced by the reduction of metals ions progressively coalesce leading to the formation of carbon-supported PtRu nanoparticles (PtRu/C electrocatalyst). The results of PtRu/C electrocatalysts preparation using electron beam and gamma irradiation are shown in Table 1. The water/2-propanol solution containing Pt(IV) and Ru(III) ions used in the preparation of PtRu/C electrocatalysts showed a dark brown color before the addition of the carbon support and irradiation. After irradiation and separation of the solid (PtRu/C electrocatalyst) by filtration, the reaction medium becomes colorless suggesting that all of
Table 1: Influence of electron beam and gamma irradiation on Pt:Ru atomic ratio and average crystallite size of the PtRu/C electrocatalysts (20 wt% of metals, nominal Pt:Ru atomic ratio of 50:50, water/2-propanol volumetric ratio of 25/75).

| Source       | Irradiation time (h) | Dose rate (kGy h\(^{-1}\)) | Total dose (kGy) | Pt:Ru atomic ratio (EDX) | Crystallite size (nm) |
|--------------|----------------------|-----------------------------|------------------|--------------------------|-----------------------|
| Electron beam| 0.05                 | 5760.0                      | 288              | 47:53                    | 2.9                   |
| Gamma \(^{60}\)Co | 6                   | 0.5                         | 3                | 45:55                    | 2.7                   |

The X-ray diffractograms of Pt/C and PtRu/C electrocatalysts prepared using electron beam and gamma irradiation.

![X-ray diffractograms](image)

Figure 1: X-ray diffractograms of Pt/C and PtRu/C electrocatalysts prepared using electron beam and gamma irradiation.

The Pt(IV) and Ru(III) ions were reduced. To confirm this assumption, a qualitative test using potassium iodide [15] did not detect Pt ions in the filtrates, which suggest that all Pt(IV) ions were reduced to metallic Pt. As no Pt ions were not detected in the filtrates, and the obtained Pt:Ru atomic ratios were similar to the nominal ones (Table 1), it was considered that both electrocatalysts were obtained with 20 wt% of metal loading. Using low dose rate gamma irradiation the total reduction of metal ions was observed only after 6 h of irradiation. On the other hand, only 3 min were necessary to observe the total reduction of the metal ions using high dose rate electron beam irradiation.

The X-ray diffractograms of Pt/C and PtRu/C electrocatalysts are prepared using electron beam, and gamma irradiation are shown in Figure 1.

The cyclic voltammograms in acid medium of the PtRu/C electrocatalysts are shown in Figure 3.

The cyclic voltammograms (CV) of both PtRu/C electrocatalysts do not have a well-defined hydrogen adsorption-desorption region (0–0.4 V) and show an increase of the current values in the double-layer region (0.4–0.8 V) when compared to the CV of Pt/C electrocatalyst [17]. The increase of current values in the double region was attributed to the capacitive currents and redox process of ruthenium oxides [17, 18]. However, comparing the CVs of both PtRu/C electrocatalysts, it is observed for the material prepared using electron beam irradiation a more defined hydrogen region when compared to the material prepared using gamma irradiation. On the other hand, the material prepared using gamma irradiation showed the double layer region more pronounced. This could suggest that the material prepared using electron beam irradiation has a surface more enriched in Pt while the material prepared using gamma irradiation has a surface more enriched in Ru.

The electrooxidation of methanol was studied by cyclic voltammetry in 1 mol L\(^{-1}\) methanol in 0.5 mol L\(^{-1}\) H\(_2\)SO\(_4\) (Figure 4).

The electrooxidation of methanol started only at about 0.45 V for the PtRu/C electrocatalyst was prepared using gamma irradiation, and the current values were lower than those observed for the PtRu/C electrocatalyst prepared using an electron beam. For the latter electrooxidation started at about 0.35 V, and the performance of this catalyst was very
similar to the commercial PtRu/C electrocatalyst from E-TEK. Studies have shown that the maximum activity for methanol oxidation at room temperature could be obtained using PtRu/C electrocatalysts with low Ru coverage [19–21]. The obtained results could be explained by the different dose rates of electron beam and gamma radiation and to different reduction potentials of Pt(IV) and Ru(III) ions. Using electron beam irradiation (high dose rate), the reduction of Pt(IV) and Ru(III) ions proceeds very quickly and enhances the probability of alloying, as confirmed by XRD measurements. Thus, the carbon-supported PtRu nanoparticles obtained using electron beam seem to have a more homogeneous distribution of Pt and Ru atoms on the nanoparticles surface. On the other hand, at low dose rate (gamma source) it seems that the Pt(IV) ions were reduced before the Ru(III) ions. In this case, Ru atoms deposit preferentially on the presupported Pt nanoparticles and the resulting carbon-supported PtRu nanoparticles have a Ru-rich surface. Another possibility is that Pt(IV) and Ru(III) ions were reduced with equal probabilities by radiolytic radicals, but a further electron transfer from the less noble metal atom, Ru, to the more noble metal ion, Pt(IV), could also result in the formation of carbon-supported PtRu nanoparticles with the surface enriched by Ru atoms, which could explain the low activity of this sample for methanol electrooxidation.

4. Conclusions
An active PtRu/C electrocatalyst for methanol oxidation was easily obtained in a single step within a few minutes using electron beam irradiation. The PtRu/C electrocatalysts showed the typical fcc structure of platinum and platinum
alloys with average particle sizes of 2.5 nm. At room temperature, the material prepared using electron beam irradiation has a similar methanol oxidation performance as that of a commercial PtRu/C electrocatalyst.

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Research Article

Virus-Assembled Flexible Electrode-Electrolyte Interfaces for Enhanced Polymer-Based Battery Applications

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High-aspect-ratio cobalt-oxide-coated Tobacco mosaic virus (TMV-) assembled polytetrafluoroethylene (PTFE) nonstick surfaces were integrated with a solvent-free polymer electrolyte to create an anode-electrolyte interface for use in lithium-ion batteries. The virus-assembled PTFE surfaces consisted primarily of cobalt oxide and were readily intercalated with a low-molecular-weight poly (ethylene oxide) (PEO) based diblock copolymer electrolyte to produce a solid anode-electrolyte system. The resulting polymer-coated virus-based system was then peeled from the PTFE backing to produce a flexible electrode-electrolyte component. Electrochemical studies indicated the virus-structured metal-oxide PEO-based interface was stable and displayed robust charge transfer kinetics. Combined, these studies demonstrate the development of a novel solid-state electrode architecture with a unique peelable and flexible processing attribute.

1. Introduction

The high energy, power densities, and excellent cycling stability of lithium-ion batteries have made them the preferred system for many portable electronic applications [1]. However, the liquid-based organic electrolytes commonly used in these systems are typically flammable and impose weight and size constraints that significantly limit their application [2]. To mitigate these issues, recent efforts have focused on the development of solid-state polymer-based lithium-ion electrolytes [3–7]. Polymer-based electrolytes can be readily processed to produce thin films that can be configured into a range of packaging shapes and coatings. In addition, the composition and self-assembly aspects of different block copolymer electrolytes can be combined and tuned to create electrolyte architectures tailored to enhance specific battery properties [8–11]. These polymer-based systems are also generally lighter and resistant to fire, making them safer and potentially more economical to use. However, a key obstacle in the use of polymer-based electrolytes is their reduced ionic conductivity and the high interfacial resistance that occurs between the electrode and the solid electrolyte. One method to enhance the charge transfer and rate performance in these systems is to create nanoscale three-dimensional electrode-electrolyte interfaces with increased surface areas and shortened Li⁺ diffusion lengths. Unfortunately, it is currently difficult to achieve such three-dimensional nanoscale electrode-solid electrolyte interfaces using traditional micromachining and casting methodologies [12, 13]. Thus, there is a need to identify novel methods to interface polymer electrolytes with three-dimensional electrode patterning capabilities.

Biological macromolecules encode unique properties that include self-assembly, genetic programmability, and diverse functionalities that have made them attractive building blocks for the development of nanoparticles and surface features [14, 15]. In particular, the uniform structure and assembly processes of virus particles have been used to develop battery electrodes and memory devices [13, 16–18]. For example, Nam et al. [13] demonstrated the use of M13 virus in the fabrication of a two-dimensional electrode-polymer...
electrolyte interface with enhanced battery applications. Previous efforts by our group have developed a novel combination of genetically engineered Tobacco mosaic virus (TMV) and electroless plating techniques to obtain near uniform inorganic coatings of Ni, Co, Pt, and silica on three-dimensionally patterned virus templates [18–21]. TMV produces a rod-shaped particle of 300 nm in length by 18 nm in diameter with a 4 nm wide central channel. Individual particles are assembled from ∼2130 identical 17.5 kDa coat protein subunits wrapped in a helical array around the virus’ single-stranded RNA genome. Using the known three-dimensional structure of TMV, we created novel coat protein mutants such as TMV1cys, which has a unique cysteine codon within the N-terminus of the coat protein open reading frame [18]. The positioning of the 1cys mutation directs the viral rod attachment via its 3’ end onto gold and other surfaces, including stainless steel and SiN [18, 22]. Combined with electroless plating methods, this process yields high aspect ratio surfaces composed of evenly coated virus templates containing confluent metal coatings as thin as 15 nm in thickness [18, 22]. These virus-assembled surfaces have been shown to enhance anode function in Ni-Zn and Li-Si battery systems [18, 23].

In this study, we demonstrate the production of virus assembled three-dimensional electrodes on nonstick polytetrafluoroethylene (PTFE) surfaces. The PTFE-assembled virus enhanced electroless coatings of cobalt and served as a robust template for the casting of a poly (ethylene oxide) (PEO) and poly (methyl methacrylate) (PMMA) based block copolymer that possesses a high transference number and excellent electrochemical properties [6, 7]. The resulting virus-assembled-polymer electrolyte system was easily peeled from the PTFE surface to produce a flexible electrode-electrolyte interface with stable electrochemical properties.

2. Materials and Methods

2.1. Virus Assembly, Patterning, and Coatings. A PTFE substrate was used for electrode release studies and SEM observations, while an Au-coated stainless steel substrate, in the form of a CR2032 coin disk, was used for electrochemical testing. To pattern the virus-structured PTFE surfaces, a parafilm mask with 1 cm² square opening was applied to a PTFE-coated substrate (Bytac Saint-Gobain Performance Plastics Poestenkill, NY). Virus-assembled surfaces were prepared by immersing the parafilm-patterned PTFE or Au-coated stainless steel discs in a solution of 0.2 mg/mL purified TMV1cys in 0.1 M pH 7 phosphate buffer and incubating overnight [18]. Cobalt-coated TMV electrodes were prepared using an electroless plating method as described in Rosyon et al. [18]. In summary, a 0.02 M solution of Na₂PdCl₄ in methanol was added to the virus solution to a final Na₂PdCl₄ concentration of 0.0015 M and allowed to incubate for 30 minutes. The virus solution was then removed and replaced with a cobalt electroless plating solution consisting of 0.033 M cobalt (II) sulfate, 0.008 M glycine, 0.5 M sodium tetraborate, and 0.175 M dimethylamine borane. The plating reaction was allowed to proceed for 10 minutes. The plating solution was then removed and the virus-coated surface washed with ethanol and dried overnight under vacuum before being moved into an argon-filled glovebox.

2.2. Analysis of Virus-Assembled Surfaces. Virus-assembled surfaces were visualized using a Hitachi field emission scanning electron microscope (FESEM) at 5–20 kV accelerating potential. X-ray photoelectron spectroscopy (XPS) was performed using a Kratos Axis 165 X-ray photoelectron spectrometer operating in hybrid mode, using monochromatic Alkα radiation (1486.6 eV) at 220 W. Survey spectra and high-resolution spectra were collected with a pass energy of 160 eV and 20 eV, respectively. All data were calibrated to the hydrocarbon contamination peak at 284.8 eV. Peak fitting was carried out using CASA XPS software (http://www.casaxps.com/). Peaks were fit with a 70% Gaussian, 30% Lorentzian peak shape after subtraction of a Shirley background. X-ray diffraction (XRD) pattern was recorded on a Bruker Smart1000 using CuKα radiation.

2.3. Preparation of the Polymer Electrolyte Membrane. The PEO-b-PMMA block copolymer (Molecular weight, 3000 : 500 for PEO : PMMA) (Polymer Source Inc.) was used as starting material to synthesize the diblock copolymer PEO-b-(PMMA-ran-PMMA-Li) [6]. Electrolyte films were prepared inside an argon glove box. The diblock copolymer PEO-b-(PMMA-ran-PMMA-Li) and LiBOB (Chemetall GmbH) were mixed in previously optimized compositions of 4 : 1 (PEO : LiBOB) and dissolved in anhydrous and degassed tetrahydrofuran (THF) [6, 7]. Electrolyte membranes were then prepared by solution casting onto the virus-assembled Bytac molds or virus-assembled Au-coated stainless steel disks resulting in translucent polymer films. The drying process lasted for several days under an argon atmosphere, followed by at least 48 hours of drying under vacuum. The resulting block copolymer has a low molecular weight and a high salt loading capacity, resulting in enhanced conductivity and transference numbers [7]. Cross-sectioning of the electrode-electrolyte interface was carried out in an argon glove box using a scalpel. Cross-sections were then transported inside a sealed vial under argon and quickly mounted on an FESEM sample holder followed by placement under vacuum to limit exposure to moisture.

2.4. Half-Cell Assembly and Testing. To demonstrate the electrochemical performance of the electrode-polymer electrolyte assembly, half-cells with Li metal as a counter electrode were assembled and tested in CR2032 coin cells. Solution casting of the block copolymer onto the virus assembled Au-coated stainless steel current collector was done as described above. Washers made from Celgard membranes were used as separators to standardize the area between the virus-assembled polymer electrolyte and the Li metal counter electrode. CR2032 coin cells were annealed at 75°C for at least 48 hrs before electrochemical testing. Cycling studies were performed at 80°C at various charge-discharge rates to study the dynamics of the half-cell’s capacity performance between 3 and 0.3 volts using an Arbin battery test station.
Figure 1: SEM analysis of cobalt-coated TMV1cys assembled surface at (a) 15 K and (b) 50 K.

Figure 2: Electroless cobalt coatings on (a) a patterned TMV1cys assembled PTFE surface and (b) a patterned unmodified PTFE surface. Parafilm masks were used to delineate 1 cm² square openings.

3. Results and Discussion

Previous studies demonstrated that TMV1cys binds in an oriented fashion to numerous surfaces including gold, stainless steel, and the polymer SU-8 [18, 21]. In this study, we demonstrate for the first time that TMV1cys is capable of similarly interacting with PTFE surfaces. Scanning electron microscopy (SEM) was used to study the 3 D TMV-templated cobalt oxide electrode structure. Figure 1 shows the SEM images of cobalt-coated TMV1cys at 15 k and 50 k magnifications. The rod-shaped virus is generally attached to the surface at one end and largely oriented away from the underlying PTFE. This configuration significantly increases reactive surface area by as much as 80-fold depending upon virus concentration [23]. The electroless plating process can be seen to provide very consistent and reproducible cobalt coatings that have a continuous metal layer ~20 nm in thickness. Attempts to assemble the unmodified wild-type virus onto the PTFE surface were unsuccessful (data not shown) indicating that the thiol groups on the TMV1cys virus are critical for assembly on PTFE surfaces. Based on this finding, we hypothesize that the electronegativity of the PTFE attracts
Figure 4: XPS spectrum of the (a) Co 2p3/2 and (b) O 1s peaks from a cobalt-coated TMV1cys surface.

Figure 5: Flexible polymer electrolyte-TMV1cys templated electrode half-cell. Virus assembly, cobalt coating, and polymer casting were done on a PTFE (Bytac) surface. The assembled half-cell was then peeled from the PTFE surface (picture taken inside an argon glovebox).

the sulfhydryl groups exposed at the end of the TMV1cys particle, thus allowing the virus to form a noncovalent association with the PTFE surface. Once assembled on the PTFE surface, TMV1cys functions as a robust template for the deposition of cobalt and other metals via electroless plating as previously described [18, 23]. The inability to produce uniform cobalt coatings on bare PTFE via electroless plating further demonstrates the importance of the assembled TMV1cys as a template for the attraction and deposition of metals on this surface (Figure 2).

Analysis of the cobalt-coated TMV1cys substrates was performed to determine the suitability of this material to function as a battery electrode. XRD patterns generated from the gold-coated stainless steel substrate and a similar substrate carrying the cobalt-coated TMV1cys are shown in Figure 3. The XRD pattern of the substrate alone shows the characteristic peaks for FeNi alloy that is the main component of the stainless steel and the peaks for the Au coating. The XRD pattern from the cobalt coated TMV1cys substrate clearly shows additional peaks at 30.6°, 36.7°, and 64.8° beside the peaks for stainless steel and Au. The positions of these new peaks are consistent with the XRD pattern of crystalline Co3O4. Similarly, high-resolution XPS scans for Co 2p show a ratio of ~10% metal and ~90% oxide with Co 2+ in the high spin state (Figure 4). This amount of oxide at the surface was consistent with the exposure of the metal surface to air. Combined these elemental studies indicate that the cobalt-coated virus surfaces are sufficient for electrode function.

Solution casting of the block copolymer electrolyte onto the cobalt-coated virus-assembled PTFE surface produced flexible translucent films as seen in Figure 5. Once the polymer film was dried, it could be easily peeled from the PTFE surface, leaving no residue from the virus-patterned electrode behind. Cross-sections of the electrode-electrolyte assembly were studied using SEM as shown in Figure 6. The lower-magnification SEM image (Figure 6(a)) clearly shows the flat released backside of the TMV1cys electrode that was peeled from the PTFE surface. The higher-magnification SEM image (Figure 6(b)) shows that the casting method permits the polymer electrolyte to fully intercalate within the three-dimensional TMV1cys electrode. The integration of the polymerized diblock copolymer provides the means for the simple release of the cobalt-coated virus electrode from the mold. The ability to cleanly peel off the cobalt-coated virus templates from the PTFE surface represents a unique means of patterning and producing flexible nanofeatured electrode-electrolyte interfaces.
To characterize the electrochemical properties of the TMV1cys enabled electrode-electrolyte interface, a coin-cell-assembled electrode was discharged (lithiation) and charged (delithiation) at various current densities at 80°C. Figure 7 shows the rate performance of the TMV polymer electrode and the charge/discharge curves at different current densities. The TMV1cys electrode shows a stable initial capacity of 0.7 mAh cm\(^{-2}\) at 0.06 C. Subsequent increases in the C rate by two-, five-, and ten-fold produce only incremental losses in electrode capacity (Figure 7(a)). In addition, restoration to the original C rate fully restored electrode capacity. These observations are consistent with good charge transfer kinetics at the electrode-electrolyte interface and indicate that the virus-assembled cobalt anode is in good contact with the block copolymer electrolyte and that this contact is stable throughout the experiment. Combined these data indicate that the virus-patterned cobalt oxide surfaces can be readily intercalated with a polymer electrolyte to produce robust high-surface-area solid-state electrode-electrolyte interfaces.

4. Conclusion

The demand for portable electronic devices is driving the development of compact power sources that combine enhanced energy storage characteristics with flexible battery architectures that are both safe and durable. Here, we describe the development of a novel virus-assembled electrode-electrolyte interface system that effectively produces tailored battery components of various form factors. This electrolyte-electrode interface leverages the use of a biologically driven
assembly process to produce a three-dimensional nanostructured conductive surface of high area. The TMV1cys assembly process also provides a robust surface for the deposition of electrode metals that could not otherwise be obtained on a PTFE surface. This electrode-electrolyte interface is further enhanced by the use of a low-molecular weight block copolymer electrolyte material that efficiently intercalates within the nanostructured electrode surface to produce a flexible electrode backing. Combined with the virus’ ability to assemble onto nonstick PTFE surfaces, this process represents a novel approach toward the production of peelable battery components that can address the size and shape constraints of device designs.

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Research Article

Synthesis of High-Quality Carbon Nanotube Arrays without the Assistance of Water

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Long and high-quality carbon nanotube (CNT) arrays have been synthesized through a chemical vapor deposition process. The Fe/Al2O3 on silicon was used as the catalyst, ethylene as the carbon source, and a gas mixture of Ar and H2 gases as the carrying gas. It is found for the first time that the high-quality and superlong carbon nanotube array can be improved by varying the content of hydrogen and carbon source.

1. Introduction

Carbon nanotubes (CNTs) have been extensively studied recently due to their unique structures and excellent mechanical, electrical, and chemical properties [1–14]. The special characteristics of carbon nanotubes arise from their atomic structures, number of walls, diameters, and lengths. For example, CNTs can be either metals or semiconductors depending on chirality and diameters [15, 16]. The alignment, diameter, and number of walls of a CNT significantly affect its mechanical and electrical properties and thus can impact a wide range of applications such as probe microscopy tips [17, 18], field emission devices [19], solar cells [7, 20], electromechanical devices [21], and structural composites [22, 23]. Recently, new technology based on the properties of very long aligned CNTs has been developing rapidly. The long nanotubes can be spun into fibers [20, 24–30] that are much stronger than all current structural materials. This will allow revolutionary advances in lightweight and high-strength applications. Intense research efforts have been undertaken to synthesize aligned long CNTs [31–38]; nevertheless, many limitations to the synthesis of very long aligned CNTs remain. Recently, the Iijima [38] group succeeded in growing 2.5 mm long aligned single-walled carbon nanotubes of high purity using the water-assisted chemical vapor deposition (CVD) technique. For the synthesis of very long aligned CNTs by CVD technique, the catalyst activity and lifetime are very important factors. The coating of the catalyst particles by amorphous carbon during CVD reduces the catalyst activity and stops the growth of CNTs. The successful synthesis of very long CNTs can be achieved if the catalyst activity and lifetime can be enhanced [38]. The control of the diameter and number of walls of CNTs represents one of the most basic issues in developing nanotube growth methods. There are some reports that the size of the catalyst used in thermal catalytic CVD can define the diameter of as-grown nanotubes [39–41]. This hypothesis has been supported by the observation that catalytic particles at the ends of CVD-grown CNTs have sizes commensurate with the nanotube diameters. The direct growth of long CNTs with a tunable number of walls by adjusting the catalyst thickness can be very promising due to its simplicity and the possibility of wide applications, which have not been yet reported. Accordingly, the developments of synthesis techniques that can control the length, diameter, alignment, and number of walls of the nanotubes have become the most important part of intense research. In the present study, we report the successful synthesis of 4 mm-long aligned CNTs by improving the flow ratio of hydrogen in the catalytic thermal CVD technique without waterassistance.
2. Experimental

The long aligned CNTs were synthesized using a conventional CVD technique. High-quality and long CNTs arrays were synthesized in an atmospheric pressure quartz tube furnace having the inner diameter of 47 mm. Ar with 30% H\textsubscript{2} was used as a carrier gas, and pure ethylene served as the carbon source. CNTs growth was performed in quartz tube furnace. High-purity ethylene (99.999%) was used as the source gas of carbon. Argon (99.999%) was used as a carrier gas, and hydrogen (99.999%) was introduced with Ar at 1 atm pressure.

Typical CVD growth was carried out at 750°C. During heating of the reactor, a steady flow of 300 sccm argon and 150 sccm H\textsubscript{2} was maintained into the chamber. After the growth temperature was attained, the total gas flow rate throughout the deposition process was maintained at 500 sccm. The optimum CVD condition was determined by balancing the relative levels of ethylene and H\textsubscript{2}. This sharp optimum condition was achieved for the growth of long aligned CNTs array by studying the dependence of the height of CNTs on ethylene-to-H\textsubscript{2} ratios, which has been reported earlier. In contrast to the report of Hata et al. \cite{38} where the catalyst was active only for 30 min, we achieved very long time (ca. 3 h) catalyst activity by optimizing the CVD parameters very accurately, which resulted in the growth of aligned CNTs as long as 4 mm. Si wafers (5 mm × 5 mm) with 1 μm thickness SiO\textsubscript{2} layer were used as the substrate for the growth of long CNTs. The Al\textsubscript{2}O\textsubscript{3} layer of thicknesses (10 nm) was used as a buffer layer between the Si/SiO\textsubscript{2} substrate and the catalyst. The Al\textsubscript{2}O\textsubscript{3} layer was deposited on the substrate on which thicknesses (1.0 nm) of Fe catalyst layer were deposited to grow CNTs. Both the catalyst and buffer layers were deposited by an e-beam evaporation technique. The progress was normally for few seconds (5 s for 1.0 nm thick Fe catalyst) to achieve catalyst thicknesses. The variation of Fe catalyst thickness was found to be most effective for controlling the number of walls of the aligned CNTs. The CNTs array with thicknesses up to 4.0 mm can be synthesized by improving the flow ratio of hydrogen to 30%. Amount of hydrogen is much larger than the 6% of total flow gas reported previously.

The microstructures of the deposited nanotubes were studied by scanning electron microscopy (SEM, Hitachi FE-SEM S4800 operated at 1 kV) and transmission electron microscopy (TEM, JEOL JEM-2100F operated at 200 kV). TEM samples were prepared by drop-casting CNT/ethanol solutions onto copper grids in the open air. Raman measurements were performed on Renishaw in Via Reflex with excitation wavelength of 514.5 nm and laser power of 20 mW at room temperature.

3. Results and Discussion

Long CNTs array up to 4.0 mm was synthesized through a chemical vapor deposition process. The experimental results showed that the synthesis of aligned CNTs is very sensitive to the CVD parameters. The large amount of hydrogen during CVD plays a critical role in synthesizing long aligned CNTs. The hydrogen acts like a weak oxidizer and helps to remove the amorphous carbon deposited on the catalyst surface during CVD. The deposition of amorphous carbon layer on the catalyst surface reduces the activity and lifetime of the catalyst; as a result long continuous CNTs cannot be synthesized. By optimizing the CVD parameters and the hydrogen concentration, it is possible to grow long CNTs with high reproducibility.

It seems that upon heating to CNTs growth temperature the Fe nanoparticles formation occurred from the thin Fe layer. These particles were then trapped in the holes of the Al\textsubscript{2}O\textsubscript{3} surface. The shape of the metal particles was also guided by the dip traps of the Al\textsubscript{2}O\textsubscript{3} surface, and the appropriate shape of the nanoparticles is very important for growing high-quality CNTs. The surface morphology of the Fe-coated Al\textsubscript{2}O\textsubscript{3} surface was also studied by scanning electron microscopy (SEM) (Figure 1(a)), and the Fe film breaks into uniform nanoparticles with an average size of about 14 nm. SEM images of the aligned CNTs deposited for 30 min by using Fe catalyst layer (1.0 nm) are shown...
in Figure 1(b), which indicates high-quality aligned CNTs without any defect. Figure 1(c) shows that the height of aligned CNTs is up to 4 mm.

Figure 2(a) compares the side views of the CNT arrays synthesized with different growth times of 10, 30, 60, 120, and 180 min. The CNTs are highly aligned with each other in the above arrays. The plot of variation of growth rate with time for 5 h deposition is shown in Figure 2(b). The average growth rate is ca. 0.54 μm/s in the first 3 hours. In the case of the aligned CNTs grown from a 1.0 nm thick Fe layer, the growth rate of CNTs was found to not vary for long time. The plot indicates gradual slow decrease in the growth rate of CNTs in the first 3 hours. The growth rate is directly related to the catalyst activity during deposition. The balance between the rate of amorphous carbon deposition on the catalyst surface and removal of that amorphous carbon layer by hydrogen is the factor keeping the catalyst active for a long time. However, with increasing time the rate of amorphous carbon deposition increases compared to the rate of removal, which results in the decrease in the catalyst activity; the growth rate decreases followed by saturation, and finally the growth terminates due to the termination of catalyst activity. Figures 3(a) and 3(b), respectively, show the transmission electron microscopy (TEM) and high-resolution TEM images of individual CNTs grown for 180 min. The outer diameters of CNTs are around 12 nm, and the inner diameters are about 9 nm. The CNTs synthesized by this approach are all multi-walled with the interlamellar distance of 0.34 nm.

The CNT arrays had been further investigated by Raman spectroscopy. The intensity ratios of D-band to G-band intensity ratio may reflect the structure integrity of the CNTs in the clean arrays. Figure 4(a) shows typical Raman spectra of the top surfaces in the CNTs arrays synthesized with the growth times of 10, 60, and 180 min. The D-band and G-band are located at about 1337 cm⁻¹ and 1571 cm⁻¹, respectively. Figure 4(b) carefully compares the ID/IG values from the top of the CNTs arrays. It was found that the ID/IG values remain almost unchanged with the increasing growth times for both cases. In other words, the synthesized CNTs show high structure integrity.

The growth mechanism of long CNT arrays had been previously investigated by several groups [28, 29]. For instance, Peng et al. [28] found that the catalyst played a critical role in the structure and size of the resulting CNTs during a chemical vapor deposition process. A tiny variation in the catalyst composition produced the so-called pearl-like CNT which was very important to the spinnability of the CNT array. In this work, we also believe that the high-performance catalyst should be responsive for the high-quality CNT array without the assistance of water. It was found that uniform nanoparticles could be produced upon
heating treatment of the catalyst in mimicking the synthetic condition. More efforts are underway.

4. Conclusions

In summary, we have developed a catalyst structure consisting of a thin Fe film on a dense Al$_2$O$_3$ buffer layer that was deposited on a Si substrate by e-beam evaporation. This catalyst lifetime was prolonged (>120 min), resulting in growth of CNT arrays that are 4 mm long. These CNT lengths are much longer than those reported previously. Increasing fraction of hydrogen was found to reduce the formation of amorphous carbon and to extend the catalyst life. Our synthesis of long CNT arrays requires a very low fraction of ethylene and large fraction of hydrogen.

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Research Article

Controlled Synthesis of Manganese Dioxide Nanostructures via a Facile Hydrothermal Route

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Manganese dioxide nanostructures with controllable morphological structures and crystalline phases were synthesized via a facile hydrothermal route at low temperatures without using any templates or surfactants. Both the aging duration and aging temperatures were the main synthesis parameters used to influence and control the rate of morphological and structural evolution of MnO2 nanostructures. MnO2 nanostructures comprise of spherical nanoparticulate agglomerates and highly amorphous in nature were formed at lower temperature and/or short aging duration. In contrast, MnO2 nanostructures of sea-urchin-like and nanorods-like morphologies and nanocrystalline in nature were prepared at the combined higher aging temperatures and longer aging durations. These nanostructures underwent notable phase transformation from δ-MnO2 to α-MnO2 upon prolonged hydrothermal aging duration and exhibited accelerated rate of phase transformation at higher aging temperature.

1. Introduction

One-dimensional manganese dioxide (MnO2) nanostructures such as nanorods, nanowires, and nanofibers have generated intense research interests over the past recent years due to their superior optical, electrical, catalytic, magnetic and electrochemical properties [1–3]. Such manganese dioxide nanostructures are of considerable importance in technological applications and have been intensively investigated as promising electrode materials in primary/secondary batteries and electrochemical capacitors due to their excellent electrochemical properties, low-cost, environmentally benign, and ease of preparation [4–7]. Various approaches have been used to fabricate manganese dioxide, such as self-reacting microemulsion [8], precipitation [9], room-temperature solid reaction [10], sonochemical [11], and hydrothermal methods [12]. The hydrothermal method is a powerful synthesis approach for synthesizing various forms of manganese oxides and affords various advantageous features including the use of mild synthesis conditions such as pH and temperature, and a wide range of precursors that can be used.

Various types of inorganic nanowires and nanorods have been synthesized with the aid of templates or catalysts. Templates are being used to confine the growth of crystals, while catalysts may act as energetically favorable sites for the adsorption of reactant molecules [13]. However, the introduction of templates or catalysts to a reaction system is often accompanied by drawbacks such as the need to prepare or select appropriate templates or catalysts. Besides, impurities in the final product may be difficult to be removed, thereby making the overall synthesis process more complicated and costly. As such, any synthetic method without the need to use any catalyst or template is more favorable for the preparation of low-dimensional nanostructures. Recently, a hydrothermal or solvothermal method has been employed to prepare one-dimensional nanoscaled materials, for example, α-MnO2, without the use of templates or catalysts [14]. This method is superior to traditional methods since no specific and expensive equipment is required for synthesizing nanostructured materials at low temperatures. The hydrothermal preparation of manganese dioxides involved mainly redox reactions of MnO4− and/or Mn2+ or the phase transformation of granular manganese dioxide precursors [15]. A common approach for the synthesis of single-crystalline α-MnO2 nanorods was based on the hydrothermal reaction of MnSO4 and KMnO4 [16]. DeGuzman et al. prepared fibrous α-MnO2 through redox reactions between KMnO4 and MnSO4.
[17]. However, some minor differences in the morphology of final products have been observed as specific reaction conditions were being altered slightly. Parameters such as temperature, time and capping molecules can influence the growth of nanocrystals under nonequilibrium kinetic growth conditions in the solution-based approach [18]. Henceforth, the controlled synthesis of manganese dioxide nanostructures with favorable surface morphology, phase structure, crystallinity, and high reproducibility remains a considerable challenge.

This paper reports the controlled synthesis of various MnO2 nanostructures via a facile and mild hydrothermal route without using any physical template and addition of any surfactant. Both δ-MnO2 and α-MnO2 nanostructures were synthesized based on the hydrothermal reaction of MnSO4 and KMnO4 in aqueous medium and at mild temperatures. Effects of hydrothermal synthesis conditions on the evolution of structural morphology and phase transformation of MnO2 nanostructures were investigated.

2. Materials and Methods

2.1. Synthesis of MnO2 Nanostructures. The synthesis of MnO2 nanoparticles was based on the method reported by Xiao et al. with some modifications [19]. MnO2 nanoparticles were synthesized by mixing aqueous solutions of KMnO4 and MnSO4 at ambient temperature and pressure, and the pH of solution mixture was adjusted to ~1 with concentrated HNO3. The aging temperatures were fixed at 25°C and 80°C, whereas the aging duration varied between 1 hour and 24 hours. The reaction product was collected by filtration, washed, and then air-dried at room temperature.

2.2. Characterization of MnO2 Nanostructures. The surface and structural morphologies of MnO2 samples were studied by scanning electron microscope (SEM) (JEOL Model JSM 6390LA) and transmission electron microscope (TEM) (JEOL Model JEM-1230). For SEM imaging, all MnO2 samples were pre-coated with a thin platinum layer using an Ion Sputtering Device (JFC-1100 E) in order to reduce the inherent charging effect. As for TEM imaging, MnO2 samples were first being dispersed well in ultrapure water by ultrasonication. 1 μL of the resulting dispersions were then drop-coated onto Formvar-covered copper grids and air-dried. The specific surface area and pore size distribution of MnO2 samples were determined using the nitrogen adsorption-desorption (BET) analyzer (Micrometrics ASAP 2010) at 77 K. The phases of MnO2 samples synthesized at different aging durations and temperatures were studied using a X-ray diffractometer (XRD) (Scintag) with Cu Kα radiation source.

3. Results and Discussion

A facile and mild hydrothermal route was being used to synthesize MnO2 nanostructures without the use of any template or surfactant. Figures 1 and 2 show SEM micrographs of MnO2 samples after being aged for various durations at temperatures of 25°C and 80°C, respectively. Both the aging duration and aging temperature were observed to have substantial effect on the shape and morphology of MnO2 nanostructures formed. At an aging temperature of 25°C, spherical agglomerates of MnO2 nanoparticles were obtained at aging durations of between 0 hour and about 4 hours (Figure 1). In absence of surfactant, MnO2 nanoparticles showed high tendency to aggregate and formed spherical agglomerates of various sizes [20]. However, upon prolonged aging for 8 hours, nanorod-like structures began to develop on surfaces of individual MnO2 nanoparticles. Upon aging for more than 24 hours, well-defined nanorods had developed around these spherical MnO2 nanoparticles. Upon aging for more than 24 hours, well-defined nanorods had developed around these spherical MnO2 nanoparticles to form sea-urchin-like MnO2 nanostructures.

As shown in Figure 2, MnO2 samples synthesized initially at 80°C comprised of mainly spherical agglomerates. However, such spherical agglomerates were no longer discernible but large aggregates of nanorod-like structures were observed upon being aged for 4 hours at 80°C. Well-defined and fully developed MnO2 nanorods were formed after being aged for extended durations of 8 and 24 hours at 80°C, respectively. The aging temperature was found to play a crucial role in accelerating the rate of evolution of MnO2 nanostructures from spherical agglomerates to aggregates of well-defined nanorods. On the contrary, no MnO2 nanostructure was formed at the aging temperature of below 20°C even after being aged for a week. At the aging temperature of 25°C, the rate of structural evolution was observed to be rather slow, with MnO2 nanostructures of sea-urchin-like shape formed only after being aged for 24 hours (Figure 1). However, at the elevated aging temperature of 80°C, distinctive and well-defined MnO2 nanorod-like nanostructures were clearly discernable after being aged for 4 or more hours (Figure 2). A higher aging temperature appeared to favor the growth of one-dimensional (1D) MnO2 nanostructures which could be attributed to the accelerated rate of decomposition of MnSO4 to form MnO2 at elevated temperatures. These nanorod-like nanostructures continued to grow in length due to their anisotropic nature and eventually led to the formation of nanowires [21]. Henceforth, hydrothermal synthesis conditions could be controlled and optimized was for the synthesis of MnO2 nanostructures of desired morphology and crystalline phase.

Figure 3 shows TEM micrographs of MnO2 nanoparticles synthesized at various aging durations at aging temperatures of 25°C. The evolution of nanorod-like nanostructures was observed to have initiated from the surfaces of MnO2 nanoparticles after being aged for 4 hours at 25°C. More distinctive and defined nanorod-like nanostructures had evolved from spherical MnO2 nanoparticles after being aged for 8 hours. Long and well-defined nanorods were observed to have developed on MnO2 nanoparticles forming sea-urchin-like nanostructures after being aged for 24 hours at 25°C.

As shown in Figure 4, TEM micrographs depicted the rapid evolution of well-defined nanorods from individual spherical MnO2 nanoparticles at elevated aging temperature of 80°C. Agglomerates of MnO2 nanoparticles were observed to have transformed rapidly into sea-urchin-like
Figure 1: SEM micrographs of MnO$_2$ samples synthesized at various aging durations at 25°C.

Figure 2: SEM micrographs of MnO$_2$ samples synthesized at various aging durations at 80°C.
Figure 3: TEM micrographs of MnO₂ samples synthesized at various aging durations at 25°C.

Figure 4: TEM micrographs of MnO₂ samples synthesized at various aging durations at 80°C.
nanostructures after being aged for 4 hours at 80°C. All MnO₂ nanoparticles were observed to have transformed completely into well-developed nanorods after being aged for 8 hours. No notable morphological changes of nanorods were observed after prolonged aging duration for up to 24 hours at 80°C. The diameter and length of well-defined MnO₂ nanorods ranged from 20 to 30 nm and 300 to 400 nm, respectively.

The hydrothermal synthesis route used in the present study had been shown to be a facile and mild synthesis approach for the synthesis of manganese dioxide nanostructures of desired morphology through judicious control of both the aging temperature and aging duration. In this synthesis approach, neither catalyst was needed to provide energetically favorable sites for the absorption of reactant molecules nor template was needed to direct the growth of nanorods. The driving force for the growth of MnO₂ nanorods during the synthesis process could be derived from the inherent crystal structure of MnO₂ material and its chemical potential in solution [22]. Based on our experimental observations, the formation mechanisms of MnO₂ nanostructures could entail the following processes. MnO₂ nanoparticles were initially produced by the redox reaction between MnSO₄ and KMnO₄. These MnO₂ nanoparticles would subsequently aggregate to form spherical agglomerates due to their high surface energies. During prolonged aging duration, MnO₂ nanoparticles would gradually transform into nanorods under the specific aging conditions. The gradual transformation of MnO₂ nanoparticles into nanorods could be attributed to their one-dimensional growth and anisotropic nature. Such processes obey the well-known “Ostwald Ripening” process, in which larger nanorods grow at the expense smaller ones because of differences in their surface energies. Similar formation mechanism of MnO₂ nanorods had been reported by Tang et al. with single-crystalline α-MnO₂ nanorods being successfully synthesized via a facile hydrothermal approach without any template and surfactant [23].

Figure 5 shows the effect of aging duration on the specific surface area and total pore volume of MnO₂ nanostructures synthesized at different aging temperatures of 25°C and 80°C. For MnO₂ samples synthesized at aging temperature of 25°C, both specific surface area and total pore volume were observed to increase with increasing aging durations. The specific surface area and total pore volume of MnO₂ samples increased from 91.1 m²/g and 0.225 cm³/g as prepared (or without aging) to 130.5 m²/g and 0.410 cm³/g, respectively, after being aged for 24 hours. Such increases could be attributed to microstructural changes associated with the gradual transformation of tightly packed nanoparticles into nanorod-like structures during aging. MnO₂ nanostructures of evolving nanorods should possess higher porosity as indicated by the increasing specific surface area and total pore volume with longer aging durations. In contrast, MnO₂ samples synthesized at aging durations between 0 and 8 hours at 80°C showed substantially higher values of specific surface area and total pore volume which were comparable or even higher than MnO₂ samples synthesized after being aged for 24 hours at 25°C (Figure 5). However, there was a notable decrease in both specific surface area and total pore volume for samples synthesized after being aged for 24 hours (Figure 5(b)) which could be due to aggregation and realignment of fully developed nanorods. These results showed that a substantially higher rate of transformation from spherical nanoparticles to nanorods occurred at elevated aging temperature, with complete transformation being achieved within 1 hour of aging duration at 80°C as compared to more than 24 hours at 25°C.

Figure 6 shows X-ray diffractographs of MnO₂ samples as prepared and after being aged for 24 hours at 25°C and 80°C. Both types of as-prepared MnO₂ samples without any post synthesis aging showed similar broad diffraction peaks which can be indexed to δ-MnO₂ phase albeit with low degree of crystallinity. The absence of other manganese dioxide diffraction peaks indicated the high purity of these as-prepared δ-MnO₂ samples. δ-MnO₂ samples synthesized at 25°C were observed to have only partially converted to α-MnO₂ phase as indicated by the 211 peak but have remained mostly amorphous in nature after being aged for 24 hours (Figure 6(a)). In contrast, samples of highly crystalline α-MnO₂ phase were obtained after being aged for 24 hours
at a higher aging temperature of 80°C. The presence of well-defined and sharp characteristic diffraction peaks of α-MnO₂ phase showed a complete phase transformation of δ-MnO₂ into α-MnO₂ phase upon aging at 80°C for 24 hours (Figure 6(b)). Such phase transformation could be associated with the simultaneous morphological transformation from spherical nanoparticles to nanorods during prolonged aging at 80°C. These results indicated that the rates of phase and morphological transformation were temperature dependent, with accelerated rate occurred at elevated aging temperature. Our findings concurred with observations reported by other researchers that a higher aging temperature would promote accelerated rate of transformation from δ-MnO₂ to α-MnO₂ phase and associated morphological changes from spherical nanoparticles to nanorods at elevated aging temperature. Both the morphological structure and crystalline phase of the MnO₂ nanostructures could therefore be modulated by varying the hydrothermal aging duration and temperatures.

4. Conclusion

In conclusion, MnO₂ nanostructures of spherical nanoparticulate agglomerates could be synthesized and transformed into nanostructures of sea-urchin-like or agglomerates of nanorods by a facile and mild hydrothermal route without the use of any templates or catalyst. MnO₂ nanostructures underwent accelerated rate of transformation from δ-MnO₂ to α-MnO₂ phase and associated morphological changes from spherical nanoparticles to nanorods at elevated aging temperature. The both morphological structure and crystalline phase of the MnO₂ nanostructures could therefore be modulated by varying the hydrothermal aging duration and temperatures.

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Research Article

Preparation and Characterization of Nano-Polyaniline Film on ITO Conductive Glass by Electrochemical Polymerization

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Polyaniline (PANI) films were synthesized on a conducting ITO glass by potentiostatic techniques to construct a low-cost counter electrode for dye-sensitized solar cell (DSSC). The compact layer, nanoparticles, nanorods- and fibrils were observed on the top of PANI films with different constant potentials by SEM. Then the conductivity test illuminated that a polyaniline film with the highest conductivity was electrodeposited at 1.0 V. Finally, the photoelectric measurement showed that the energy conversion efficiency of DSSC with the PANI electrode was increased with the potential decreasing. And the efficiency of DSSC with PANI counter electrode at 1.0 V was higher than that with Pt electrode, owing to the high surface area, high conductivity, and excellent catalytic activity of PANI electrode. Therefore, the PANI counter electrode with excellent catalytic performance is a potential substitute for platinized electrode to save cost of DSSC.

1. Introduction

Dye-sensitized solar cells have been one of the most promising photovoltaic devices due to their high conversion efficiency and low cost. In general, a DSSC comprises a dye-sensitized titanium dioxide (TiO2) electrode, an iodide/triiodide redox electrolyte solution, and a counter electrode (CE). The function of the CE is to transfer electrons from the external circuit back to the redox electrolyte and to catalyze the reduction of the triiodide ion [1–4]. Usually, a Pt electrode formed on a transparent conductive glass substrate is used as the CE for DSSC. Although Pt exhibits excellent catalytic activity for triiodide reduction and good electric conductivity, it is extremely expensive and has the problem of reserves for large-scale application [5, 6]. Meanwhile, the shape limitation will bring transport problem for the Pt counter electrode. Future large solar electric conversion systems will prefer materials abundantly available and easily handled. Therefore, it is necessary to develop cheap materials for CEs which also exhibit high electrical conductivity, good chemical stability, and good catalytic activity to the reduction of triiodide ion.

So far, inexpensive carbonaceous materials and conducting polymers such as graphite, carbon black, carbon nanotube, graphene, poly(3,4-propylenedioxythiophene) (PProDOT), and poly(3,4-ethylenedioxythiophene) (PEDOT) have been employed as the catalytic materials on indium tin oxide (ITO) glass for the CEs. Yen et al. reported a DSSC using the carbon nanotube composite plate had a cell efficiency of 4.73% [7]. Meanwhile, Saito et al. used chemically polymerized PEDOT on fluorine-doped tin oxide (FTO) glass as the counter electrodes and obtained a conversion efficiency of 3.6% under 100 mWcm−2 [8]. Fukuri and coworkers also developed PEDOT-based solid-state electrolytes, but the cell efficiencies were <1% [9].

Recently, the conducting PANI was applied as the counter electrode materials for DSSC to replace the platinized counter electrode due to its easy synthesis, high conductivity, and unique redox properties [10]. Among the previous research, the electrode with a layer of conductive PANI coated on
conducting glass is achieved by using several methods, such as the spreading of a solution of conductive polymer on the surface of the material followed by the evaporation of the solvent [11], the in situ chemical polymerization, and then deposition on the surfaces of various materials immersed in the polymerization solution [12] or the electropolymerization of the monomers on an electrode [13–16]. Wu et al. researched on using polyaniline [5] or polypyrrole [6] as counter electrode in DSSCs by the vertical dip-tugging method [17, 18]. Moreover, it was reported that PANI was deposited on the surface of ITO conductive glass by in situ chemical polymerization [19], and the energy conversion efficiency of DSSC with PANI counter electrode reached 2.64%. However, these methods were difficult to obtain a homogeneous and compact film owing to the insolubility or poor solubility of PANI in nearly all solvents [20].

Hence, the method of electropolymerization to prepare a compact film of polyaniline is considered in this study. The PANI counter electrodes are prepared on the surface of conducting glass by using potentiostatic techniques, and the photoelectric performances of DSSC with these electrodes are also discussed.

2. Experimental

2.1. Materials. Aniline (An, analytical grade from Sinopharm Chemical Reagent Co., Ltd) was purified by distillation under reduced pressure prior to usage. Analytical-grade reagents, H2SO4, iodine, lithium perchlorate, acetonitrile, ethanol, and acetone (Sinopharm Chemical Reagent Co., Ltd), were used without any pretreatment. All solutions were prepared from deionized water. ITO glasses from the Institute of Plasma Physics (CAS) were used as the substrates. Anhydrous lithium iodide, 4-tert-butylpyrididine, 1-methyl-3-propylimidazolium iodide, methoxypropionitrile were provided by Fluka Chemical Corporation. TiO2 colloid and dye N719 were the commercial product purchased from Solaronix SA (Switzerland). Pt counter electrode bought from Dyesol Limited was used in the contrast test.

2.2. Preparation of PANI Electrode. ITO glasses were rinsed with deionized water and immersed in ethanol ultrasonically for 15 min, then immersed in acetone ultrasonically for 15 min before PANI electropolymerization. One ITO glass sheet was used as working electrode for PANI deposition. The deposited area was 1 cm² with other area insulated by an adhesive tape coating. A Pt sheet was used as the counter electrode. PANI was electroposited on the surface of ITO glass by potentiostatic method at different constant potential (1.0 V, 1.1 V, 1.2 V, and 1.3 V) for 500 s in a 0.2 M aniline/0.5 M sulphuric acid (H2SO4) electrolyte solution. Finally, PANI-modified glass plate was immersed in 0.5 M H2SO4 statically in order to expel aniline monomer and oligomer PANI from the polymeric film and then rinsed with deionized water and ethanol for several times and dried in a vacuum at 60°C for 24 h.

2.3. Assembling of DSSCs. Nano-TiO2 colloid was dropped on the ITO glass plate by a doctor scraping technique to form a porous film. Then the TiO2 porous film was sintered by annealing at 450°C for 30 min. After cooling to 100°C, the TiO2 film was immersed in an ethanol solution of N719 dye for 24 h. Finally, DSSC was assembled by injecting a drop of electrolyte with I2, LiI, 1-methyl-3-propylimidazolium iodide, 4-tert-butylpyrididine in methoxypropionitrile into the aperture between the TiO2 porous film electrode and the PANI electrode.

2.4. Measurement and Characterization. The SEM image of the sample was performed using a scanning electron microscopy (LEO1550, GER). FTIR spectrum of the PANI was recorded in the range of 500–4000 cm⁻¹ using FTIR spectroscopy (Perkin Elmer 1760, USA). The conductivity of PANI was measured by manual four probe instrument (MP 1008, WENTWORTH, UK). Photocurrent-voltage characteristics of the DSSCs were obtained by a Keithley model 2400 digital source meter using an Oriol 91192 solar simulator equipped with AM 1.5 filter and intensity of 100 mW/cm². All the polymerization experiments were performed on an electrochemical workstation (CHI660A, CHInstrument, CHN).

3. Results and Discussion

3.1. Micromorphy of PANI. The photograph of the polyaniline films electrodeposited at 1.0, 1.1, 1.2, and 1.3 V on the ITO substrate is shown in Figure 1. It can be seen the emerald PANI thin layer is deposited on glass from Figure 1(a). With the electropolymerization potential rising, the thickness of PANI layer has a significant increase, and the color of PANI film becomes darker and darker. The PANI film electrodeposited at 1.3 V closes to dark green.

Figure 2 shows the SEM images of PANI electrodes at different constant potentials. At the low potential (1.0 V), PANI nuclei grow on the bare electrode to form a compact layer, which adhered to the electrode surface strongly. Increasing the potential to 1.1 V, some PANI particles with the diameter in range of 800 nm appear on the top of PANI nodules layer. As the potential shifts to more positive potential, the PANI particles disappear and the nanorods emerge on the surface of PANI film. At 1.2 V, PANI mainly forms the short nanorods with the length of 500 nm and with the diameter of 100 nm. At the higher potential, the PANI deposition shows an extensive fibril growth. At 1.3 V, PANI forms the long fibril with the length of 1 μm and with the diameter of 100 nm. This growth of PANI has no directional alignment, and the nanorods or fibrils are nonuniform. It is clearly demonstrated that the formation of thick, short nanorods is preferred at low polymerization potential while the development of thin, long fibrils is more easily achieved at higher potential.

3.2. Molecular Structure of PANI. Figure 3 shows the FTIR spectra of PANI deposited on conductive glass at different potentials. It is clear that the characteristic peaks of polyaniline (KBr, thin film, cm⁻¹) present near the wavenumbers of 1561 (C–C stretching mode for the quinoid ring), 1484 (benzenoid rings vibration), 1295 (C–N stretching mode),
1240 (C–N\textsuperscript{+} stretching vibration), 1077 (a vibration mode of the –NH\textsuperscript{+}= structure), and 792 (out-of-plane bending vibration of C–H) cm\textsuperscript{-1} \cite{21–27}. As the potential reaches to 1.3 V, the characteristic peaks of PANI become stronger, and the band intensity significantly enhances with the gradual rise of constant potential. It indicates that the PANI can be formed by potentiostatic polymerization at 1.0 V, and the more PANI obtained with the potential increasing can make the peak intensity of PANI stronger.

3.3. Conductivity Measurement. The conductivity of PANI deposited at different potentials is measured by four-point probes instrument and then calculated by this equation:

\[ \sigma = \frac{\ln 2}{\pi d} \cdot \frac{I_{12}}{U_{34}}, \]

where \( \sigma \) is the conductivity, \( d \) is the thickness of films, \( I_{12} \) is the current from 1 to 2 points, and \( U_{34} \) is the potential difference between 3 and 4 points. The schematic diagram of four-point probes method is shown in Figure 4.

Meanwhile, polymer thickness is estimated from the amount of charge \( Q_A \) according to the equation \cite{28}:

\[ d = \frac{Q_A M_w}{z F A \rho}, \]

where \( Q_A \) is the charge under the potentiostatic electrodeposition, \( M_w \) is molecular weight of aniline, \( z = 0.5 \) (number of electrons/aniline unit), \( A \) is area of the electrode, \( \rho \) is specific density of aniline, and \( F \) is Faraday’s constant. Table 1 presents the thickness and the conductivity of PANI layers at 1.0, 1.1, 1.2, and 1.3 V. It can be found that the film thickness of PANI increases with the polymerization potential increasing, but the conductivity of PANI decreases. As the potential is 1.0 V, the PANI film is very thin (about 33 nm). With the polymerization potential rising, the current density on the PANI/glass electrode is increased, leading to the thickness of PANI film increased also; however, the more thickness of PANI film results in the reduction of conductivity. The best conductivity of PANI deposited on conducting glass substrate reached to 62.59 S/cm when the polymerization potential was 1.0 V. Otherwise, the increase of current density on the PANI/glass electrode caused by the rise of potential leads to the peroxidation of PANI, and the conducting emeraldine form is converted into a nonconducting full-oxidation state. This transformation process among the different redox states of PANI also results in the reduction of conductivity.

3.4. Photoelectric Performance of PANI Electrode. The photoelectric properties of DSSCs with PANI/glass and Pt counter electrode are shown in Figure 5, and the results are also listed in Table 2. The result indicates that the PANI/glass electrode deposited at low potential (1.0 V) with the high conductivity and electrocatalytic activity can obtain the best
photoelectric conversion efficiency, which reaches to 5.68%. This photoelectric conversion efficiency of DSSC with PANI counter electrode is higher than that with Pt electrode (4.51%), and the fill factor (FF) is about 0.488, which is higher than Pt electrode (0.483). This improvement of the photoelectric efficiency is referred to the PANI electrode with the higher surface area, the higher electrocatalytic activity, and the lower charge-transfer resistance, compared to Pt electrode. Moreover, with the potential rising from 1.0 to 1.3 V, the FF has a little change, the Jsc of DSSC with PANI counter electrode has a remarkable decrease, and the energy conversion efficiency of DSSC with PANI electrode descends from 5.68% to 4.03%. It is due to the reduction of the conductivity and increase of the charge.
and Pt counter electrode.

| Sample      | Thickness $d$ [μm] | Conductivity $\sigma$ [S/cm] |
|-------------|--------------------|------------------------------|
| PANI-1.0 V  | 0.033              | 62.59                        |
| PANI-1.1 V  | 0.212              | 42.97                        |
| PANI-1.2 V  | 0.651              | 25.18                        |
| PANI-1.3 V  | 1.367              | 10.41                        |

Figure 5: Photoelectric properties of DSSCs with different PANI and Pt counter electrodes.

Table 2: Photovoltaic performance of DSSCs with different PANI and Pt counter electrode.

| Electrode     | $J_{sc}$ [mA/cm²] | $V_{oc}$ [mV] | FF  | Efficiency [%] |
|---------------|-------------------|---------------|-----|----------------|
| PANI-1.0 V    | 16.18             | 720           | 0.488 | 5.68           |
| PANI-1.1 V    | 15.56             | 740           | 0.476 | 5.48           |
| PANI-1.2 V    | 15.19             | 700           | 0.466 | 4.95           |
| PANI-1.3 V    | 13.49             | 620           | 0.482 | 4.03           |
| Pt            | 14.60             | 640           | 0.483 | 4.51           |

transmission impedance of PANI electrode with the polymerization potential increasing. As a result, the new PANI counter electrode prepared by electrochemical synthesis method has more predominant photoelectric performance than Pt electrode. Thus, the PANI counter electrode can replace the conventional Pt counter electrode for reducing the production cost of DSSC.

4. Summary

In summary, the PANI counter electrode was successfully fabricated on conducting ITO glass by electrochemical polymerization method. The results showed that the compact layers, nanoparticles, nanorods, and fibrils of PANI successively grew on the surface of glass as the potential rising from 1.0 V to 1.3 V. The conductivity and photoelectric test indicated that the optimal constant potential of PANI electrodeposition is 1.0 V. The energy conversion efficiency of DSSC with PANI counter electrode at 1.0 V reached 5.68%, which was higher than Pt electrode. This improvement corresponds to higher conductivity, higher electrocatalytic activity, and smaller charge-transfer resistance of PANI electrode. Therefore, the PANI counter electrode with the simple preparation procedure, low fabrication cost, and excellent catalytic properties will be a credible alternative of the counter electrode from DSSCs in the future.

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