Phosphorus $K_4$ Crystal: A New Stable Allotrope

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The intriguing properties of phosphorene motivate scientists to further explore the structures and properties of phosphorus materials. Here, we report a new allotrope named $K_4$ phosphorus composed of three-coordinated phosphorus atoms in non-layered structure which is not only dynamically and mechanically stable, but also possesses thermal stability comparable to that of the orthorhombic black phosphorus ($A17$). Due to its unique configuration, $K_4$ phosphorus exhibits exceptional properties: it possesses a band gap of 1.54 eV which is much larger than that of black phosphorus (0.30 eV), and it is stiffer than black phosphorus. The band gap of the newly predicted phase can be effectively tuned by applying hydrostatic pressure. In addition, $K_4$ phosphorus exhibits a good light absorption in visible and near ultraviolet region. These findings add additional features to the phosphorus family with new potential applications in nanoelectronics and nanomechanics.

Phosphorus is a common material for industrialized production of fertilizers, organophosphorus compounds, matches, and so on. Because of its high chemical reactivity and strong toxicity, this material has received little academic attention until recently1. In 2014 phosphorus was brought to spotlight by successfully fabricating the field-effect transistors (FETs) constructed using the exfoliated black phosphorus and finding their outstanding performance by two independent groups2,3, which was exactly 100 years after the discovery of black phosphorus4. Recently, monolayer black phosphorus (phosphorene), as a new member of two dimensional (2D) materials family, has attracted tremendous attention to explore the structural varieties5, electronic properties6, unique mechanical features7, and promising applications in electronics8 and gas sensors9.

Like carbon and boron, phosphorus displays fascinating structural variability. The stable allotropes of phosphorus at ambient conditions can be classified into three major categories: white phosphorus10, orthorhombic black phosphorus11, and various forms of red phosphorus12. White phosphorus is highly reactive with air and forms three crystal structures including $\alpha$-, $\beta$-, and $\gamma$-P4. Black phosphorus possesses an orthorhombic structure ($A17$, space group: $Cmca$) at ambient conditions and transforms to rhombohedral structure ($A7$, space group: $R3m$) at around 5 GPa12. While the $A7$ structure transforms to an $\alpha$-Po-type simple cubic three dimensional (3D) structure (space group: $Pm3m$) at a higher pressure of about 10 GPa13. Amorphous and crystalline forms of red phosphorus can be evolved from white phosphorus by heating it to higher than 20 °C13. The various forms of red phosphorus are based on tubular units of five- and six-membered rings12. Although a number of the allotropes of phosphorus have already been experimentally characterized or theoretically predicted11,12,14 the discovery of new phases of phosphorus has attracted continued attention15,16. Since the most stable phase of phosphorus is the three-coordinated layered black phosphorus, an interesting question then raises: can we find a stable 3D non-layered allotrope of phosphorus at ambient conditions which is also composed of only three-coordinated atoms?

Inspired by the unique geometry of $K_4$ structure where the coordination number of each atom is three, we explore the stability and properties of phosphorus in $K_4$ crystal structure. In fact, since the identification of $K_4$ geometry in mathematics in 200817, design and synthesis of pristine elemental substances in $K_4$ structure have been an attractive scientific topic. The nitrogen $K_4$ crystal named cg-N (cubic gauche) was synthesized from molecular nitrogen under pressure above 110 GPa using a laser-heated diamond cell18. The coordination polymers and metal-organic frameworks having the topology of $K_4$ crystal were reported in 1990s19. A $K_4$ crystal carbon was proposed to be a possible metallic allotrope of carbon in 200920. Unfortunately, a latter work suggested that $K_4$ carbon is dynamically unstable21. In 2010, Dai et al. proposed a boron $K_4$ crystal that is stable under...
ambient pressure\textsuperscript{22}. All these progresses, especially experimental realization of high pressure polymeric nitrogen in \( K_4 \) structure, make the exploration of \( K_4 \) phosphorus structure very promising.

In this work, the dynamical, thermal, and mechanical stabilities of \( K_4 \) phosphorus structure are confirmed, and its electronic, mechanical, and optical properties are studied based on a series of state-of-the-art calculations. The phonon vibrational modes at the first Brillouin zone center are also simulated to aid future experimental identification of the new phase of phosphorus from Raman and infrared (IR) spectroscopy.

### Results and Discussion

**Phosphorus \( K_4 \) structure.** Different from boron \( K_4 \) crystal, where each boron atom is \( sp^2 \)-hybridized, for the stable phosphorus allotropes at ambient conditions, due to the electron lone pair, each phosphorus atom is actually \( sp^3 \)-hybridized which displays a tetrahedral bonding character\textsuperscript{6}. To remain this bonding character, we conceive the idea of building a phosphorus \( K_4 \) structure using a building block composed of four phosphorus atoms in a tetrahedron configuration that is displayed in Fig. 1(a). In the building block, the central phosphorus atom is \( sp^3 \)-hybridized, and connects its three neighboring atoms, thus remaining the tetrahedral bonding character of phosphorus at ambient conditions. With equal bond lengths and bond angles, this structural unit is used to build the high symmetric \( K_4 \) crystal.

As shown in Fig. 1(c) and (d), the optimized structure of phosphorus \( K_4 \) is body-centered cubic with 8 atoms located at the 8a (0.206, 0.206, 0.206) Wyckoff position in the conventional unit cell. Compared to \( K_4 \) boron, \( K_4 \) phosphorus has a reduced symmetry with the space group symmetry of \( T5 \) (I2\_3, No. 199). The lattice parameters are optimized using both PBE and GGA-D2 functionals, and the results are presented in Table 1. The lattice parameters of black phosphorus are also calculated using the same level of theory and compared with the experimental data. We can see that \( K_4 \) phosphorus is slightly denser than the layered phosphorus structure of \( A_{17} \).

### Table 1

|       | \( a \)  | \( b \)  | \( c \)  | \( V_0 \) | \( \Delta E \) | \( E_g \) |
|-------|---------|---------|---------|---------|-------------|---------|
| \( K_4 \) | 5.37\textsuperscript{a} | 5.37 | 5.37 | 19.36\textsuperscript{a} | 0.02\textsuperscript{a} | 1.13\textsuperscript{a} |
|       | 5.32\textsuperscript{b} | 5.32 | 5.32 | 18.82\textsuperscript{b} | 0.04\textsuperscript{b} | 1.54\textsuperscript{b}(1.07\textsuperscript{b}) |
| \( A_{17} \) | 3.28\textsuperscript{a} | 11.22\textsuperscript{a} | 4.54\textsuperscript{a} | 20.89\textsuperscript{a} | — | — |
|       | 3.32\textsuperscript{b} | 10.43\textsuperscript{b} | 4.41\textsuperscript{b} | 19.09\textsuperscript{b} | — | 0.36\textsuperscript{b}(0.08\textsuperscript{b}) |
|       | 3.3133\textsuperscript{c} | 10.473\textsuperscript{c} | 4.374\textsuperscript{c} | 18.97\textsuperscript{c} | — | 0.335\textsuperscript{c} |

\textsuperscript{a}Our calculated results at the PBE level. \textsuperscript{b}Our calculated results at the GGA-D2 level. \textsuperscript{c}Experimental data. \textsuperscript{d}Our calculated results at the HES06 level.
However, the mass density difference is much smaller than that between diamond and graphite. This is because the electron-rich nature and the resultant Coulomb repulsion hinder phosphorus to form ultra-dense phases like diamond. As compared to black phosphorus, K₄ phosphorus has a higher symmetry, leading to an isotropic network. In the unit cell of K₄ phosphorus, only one kind of P-P bond exists and all bond angles are equivalent with a bond length of 2.24 Å/2.25 Å and a bond angle of 101.7°/102.2° at the PBE/GGA-D2 level.

### Energetic stability.

To investigate the energetic stability of K₄ phosphorus, total energy calculations are performed by using both the PBE and GGA-D2 functionals, respectively. For comparison, calculations are also carried out for black phosphorus. The results are listed in Table 1. The cohesive energy of K₄ phosphorus is found to be higher than that of black phosphorus by 0.02/0.04 eV/atom at the PBE/GGA-D2 functional, showing that K₄ phosphorus is thermodynamically metastable as compared to the most stable form of phosphorous allotropes, black phosphorus.

Table 2. Relative energies of hypothetical phosphorus structures with respect to the cohesive energy of K₄ phosphorus (ΔE in eV/atom).

| Structure | ΔE<sub>PBE</sub> | ΔE<sub>GGA-D2</sub> |
|-----------|-----------------|---------------------|
| srs        | 0.000           | 0.014               |
| acs-g      | 0.013           | 0.017               |
| bcu-f      | 0.021           | 0.023               |
| eta        | 0.026           | 0.024               |
| etd        | 0.028           | 0.027               |
| etf        | 0.030           | 0.032               |
| pbg        | 0.034           | 0.033               |

*ΔE<sub>PBE</sub> = E<sub>K₄</sub> - E<sub>black</sub>; ΔE<sub>GGA-D2</sub> = E<sub>K₄</sub> - E<sub>black</sub>.*<sup>2</sup>

The calculated cohesive energy (ΔE) of K₄ phosphorus is 0.07/0.01 eV/atom lower in energy than the simple cubic structure, and 0.08/0.14 eV/atom lower in energy than the two structures of white phosphorus (β-P₄, γ-P₄), indicating that the K₄ phosphorus phase is energetically more stable than these well-known structures. To understand the reason why K₄ phosphorus is relatively stable, we investigated its atomic configuration. We note that in the three-coordinated phosphorus K₄ structure each bond angle is 101.7°, as shown in Fig. 1(a), close to that of 101.6° in the P₄H₆ molecule that was found to be the most energetically favorable configuration among those of the three-coordinated phosphorus atoms<sup>3</sup>. Therefore, the favorable geometry results in a good energetic stability of the K₄ structure.

To search for 3D non-layered phosphorus allotropes with good energetic stability at ambient conditions, we have screened many possible 3-coordinated networks. Since the phosphorus K₄ structure has the srs topology, by using the network topology approach based on the RCSR database<sup>24</sup>, many candidate structures for 3D non-layered phosphorus with different network topologies can be obtained. Due to the limited computational resources, we only consider the so called uninodal (with only one type of vertex) structures. Actually there are 78 uninodal 3-coordinated nets in the RCSR database. Among them, there are 47 nets with too short non-bonded distance, and there are 1 net with the bond angles seriously deviated from that of P₄H₆ molecule (101.6°)<sup>9</sup>. So these 48 nets should be excluded. However, not all the 30 rest nets are suitable for forming 3D non-layered phosphorus structures. For example, the optimized structures of etb and utp configurations become 5-coordinated, and the optimized structure of utg configuration turns into phosphorus chains. Finally, only 14 nets including the srs net (K₄ phosphorus) are found to be the suitable candidate structures for 3D non-layered phosphorus allotropes in which every phosphorus atom retains the sp<sup>3</sup> hybridization character of tetra-phosphorus. The cohesive energy of these hypothetical phosphorus structures are calculated by using the PBE and GGA-D2 functionals, respectively. As listed in Table 2, we find that K₄ phosphorus is energetically most stable among the studied configurations.

### Dynamic stability.

To examine the dynamic stability of K₄ phosphorus, the lattice dynamics is studied by calculating its phonon dispersion using linear response method within density functional perturbation theory<sup>9</sup>, where the force-constant matrix is calculated through differentiation of the Hellmann-Feynman forces on atoms with respect to the ionic coordinates. As shown in Fig. 2(a), the absence of imaginary modes in the whole Brillouin zone confirms that K₄ phosphorus is dynamically stable. The primitive cell of K₄ phosphorus contains four atoms, leading to nine optical and three acoustic branches. The optical branches can be classified into two groups with a frequency gap between them. The low energy group consisting of three optical branches is predominantly bond-bending type in character, while the high energy group consisting of six optical branches is bond-stretching type. All the three acoustic branches of K₄ phosphorus are linearly dispersed near the Γ point in different directions, confirming the relatively strong covalent bonds between the phosphorus atoms along these directions, while in the structure of black phosphorus the dispersion relation of the TA<sub>1</sub> modes near the Γ point in the [100] and [010] directions<sup>26</sup> is almost in quadratic form that is contributed by the weak interlayer interactions.

### Thermal stability.

The thermal stability of K₄ phosphorus is examined by performing ab initio molecular dynamics (AIMD) simulations at 300 K with a large supercell (3 × 3 × 3). We find that no structure reconstruction occurs after heating for 8 ps with a time step of 1 fs, and the total potential energy remains almost constant during the simulation. These results suggest that the new structure is thermally stable at room temperature. The heat bath is then further elevated to 1000 K. As shown in the snapshot of atomic configuration of K₄ phosphorus at the
appears, and the fluctuation in total potential energy still remains almost unchanged. This implies that the elastic constants of the phosphorus derived from the strain-stress relationship by using the finite distortion method implemented in VASP. All the calculations were performed for black phosphorus, orthorhombic black phosphorus (a17), and the results are given in Table 3 as well.

To further investigate the thermal stability of black phosphorus, AIMD simulations were also carried out for the experimentally synthesized a17 structure for comparison. As shown in Fig. S1, the a17 structure can withstand the high temperature of 1000 K without any obvious structural reconstruction. However, when temperature of the heat bath is further increased to 1200 K, both the geometries of the K4 and a17 phases are destroyed. The results reveal that K4 phosphorus is thermally as stable as the a17 phase.

**Mechanical stability and properties.** To guarantee the positive definiteness of strain energy upon lattice distortion, the mechanical stability of K4 phosphorus is examined. In the linear elastic range the elastic constant tensor forms a 6 × 6 matrix with 21 independent components. For a simple cubic lattice, only C_{11}, C_{12} and C_{44} are independent. The linear elastic constants of a mechanically stable 3D cubic lattice have to obey the Born-Huang criteria: C_{11} > 0, C_{44} > 0, C_{11} > 2C_{12}, and (C_{11} + 2C_{12}) > 0^{28}. The elastic constants of the phosphorus K4 crystal are derived from the strain-stress relationship by using the finite distortion method implemented in VASP. All the elastic constants calculated with both the PBE and the vdW-corrected GGA-D2 functionals are listed in Table 3. These constants obey all of the Born-Huang criteria for simple cubic lattices, implying that the new structure is mechanically stable. For comparison, the calculations are also performed for a17, and the results are given in Table 3 as well.

The C_{11}, C_{22}, and C_{33} (C_{11} = C_{22} = C_{33}) elastic constants of K4 phosphorus directly relate to sound propagation along the crystallographic a, b, and c axes, respectively, and reflect the stiffness to the uniaxial strains along these directions. The calculated value of C_{11} is 223.6/232.2 GPa at the PBE/GGA-D2 level. While for the orthorhombic black phosphorus, C_{12} is significantly larger than C_{11} and C_{33} due to its structural anisotropy, but it is smaller than the C_{11} elastic constant of K4 phosphorus, indicating that the K4 structure is stiffer than black phosphorus for strains along the a, b, and c axis. The benefit from the high stiffness is that the K4 structure could avoid the sliding observed in the layered phosphorus structures under a shear stress. The resistance of sliding could make K4 phosphorus more suitable for nano-mechanical applications. The single crystal bulk moduli of the phosphorus K4 crystal and black phosphorus are calculated according to the formula of bulk modulus represented by single-crystal elastic constants. The result for black phosphorus is in good agreement with the value estimated from fitting the energy-volume relationship.

### Table 3. Calculated elastic constants (C_{ij} in GPa) and bulk moduli (B in GPa) of K4 phosphorus and orthorhombic black phosphorus (a17)

|       | C_{11} | C_{12} | C_{13} | C_{22} | C_{23} | C_{33} | B    |
|-------|--------|--------|--------|--------|--------|--------|------|
| K4    | 223.6  | —      | —      | 34.6   | —      | —      | 82.4 |
|       | 223.2  | —      | —      | 31.7   | —      | —      | 39.3 |
| a17   | 43.8   | 188.4  | 16.0   | 9.8    | 2.9    | 58.5   | 34.5 |
|       | 62.8   | 199.8  | 83.2   | 29.8   | 9.6    | 80.6   | 43.8 |

end of AIMD simulations at 1000 K (see Fig. 2(b)), after heating for 8 ps, no obvious distortion in the structure appears, and the fluctuation in total potential energy still remains almost unchanged. This implies that K4 phosphorus can withstand temperatures up to 1000 K, and this phosphorus phase is separated by high energy barriers from other local minima on the potential energy surface (PES) of elemental phosphorus.

However, it is worthy to note that the temperature for evaluating the thermal stability of a crystal structure may be overestimated by AIMD simulations due to using Canonical ensemble (NVT) during the simulations. To further investigate the thermal stability of K4 phosphorus, the AIMD simulations are also carried out for the experimentally synthesized a17 structure for comparison. As shown in Fig. S1, the A17 structure can withstand the high temperature of 1000 K without any obvious structural reconstruction. However, when temperature of the heat bath is further increased to 1200 K, both the geometries of the K4 and A17 phases are destroyed. The results reveal that K4 phosphorus is thermally as stable as the A17 phase.

Figure 2. Structure stability of K4 phosphorus. (a) Vibrational band structure of K4 phosphorus. (b) Total potential energy fluctuation of K4 phosphorus during AIMD simulation at 1000 K. The inset shows the atomic configurations (3 × 3 × 3 supercell) at the beginning and end of AIMD simulations at 1000 K.
Electronic properties.

We calculate the electronic band structure and corresponding total density of states (DOS) of $K_4$ phosphorus to study its electronic properties. The results are displayed in Fig. 3(a). At the PBE level, $K_4$ phosphorus is predicted to be an indirect band gap semiconductor with a band gap of 1.07 eV as the valence band maximum (VBM) and the conduction band minimum (CBM) lie at the different points along the $\Gamma$-H path. It is well-known that the PBE functional underestimates the fundamental band gaps of semiconductors, thus the band gap of $K_4$ phosphorus is corrected by calculations using the more accurate HSE06 functional. As shown in Fig. 3(a), although both of the functionals give very similar band dispersions, the band gap of $K_4$ phosphorus calculated with HSE06 functional is increased to 1.54 eV. Figure 3(b) shows the isosurfaces of the total valence electron density calculated using the PBE functionals. As mentioned above, the phosphorus atoms are $sp^3$ hybridized due to the three P-P covalent bonds and a lone electron pair, thus the valence electrons mainly distributed along the directions of the P-P bonds and the lone electron pair.

To investigate how the energy band gap of $K_4$ phosphorus changes with the applied hydrostatic pressure, we calculate the band gap as a function of the pressure, and plot the results in Fig. 3(c), which shows that the band gap of the $K_4$ phase decreases with volume compression under hydrostatic pressure. At the GGA level, the band gap decreases almost linearly from 1.13 eV to 0.49 eV as pressure increases from zero to 6 GPa. While at the GGA-D2 level, the band gap decreases from 1.07 eV to 0.30 eV within the same pressure range.

Optical properties.

We next explore the potential applications of $K_4$ phosphorus in optoelectronics. The imaginary part of dielectric function of $K_4$ phosphorus, which is directly related to its optical absorbance, is calculated at the HSE06 level. For comparison, calculation for diamond silicon is also carried out by using the same approach. According to photon energy, the spectrum is divided into three parts, namely the infrared, visible, and ultraviolet regions, respectively. As shown in Fig. 4, the imaginary part of the dielectric function of $K_4$ phosphorus reaches the maximum at 3.18 eV, and follows by a minimum at 4.21 eV. According to the calculated results, $K_4$ phosphorus exhibits much stronger optical absorption than diamond silicon in the visible range. The absorption
structure in experiment. Also an indirect band gap semiconductor, it possesses a direct band gap of 2.4 eV at P1 point in the Brillouin zone experimental conditions, such as elevated temperature and environment of the tetramers play a decisive role 35.

In summary, based on first principles calculations, we predict a 3D stable phosphorus phase different from those of black phosphorus. Currently, diamond silicon is still the leading material of solar cells 31. It has an indirect band gap of 1.1 eV and a large direct gap of 3.3 eV, making it inefficient for sunlight absorption 32. For K4 phosphorus, although it is also an indirect band gap semiconductor, it possesses a direct band gap of 2.4 eV at P1 point in the Brillouin zone (see Fig. 3(a)), which just lies in the middle part of the spectral range of visible light. As compared with diamond silicon, the smaller direct band gap of K4 phosphorus makes it a better solar absorber. Based on above analysis, we conclude that K4 phosphorus exhibits strong optical absorption in the visible and near ultraviolet region, making it a promising candidate for photovoltaics.

To provide a possible way to experimentally identify the K4 phase of phosphorus from other phosphorus allotropes using Raman and infrared (IR) spectroscopy, we simulate the vibration properties of the K4 structure at the \( \Gamma \) point. In order to verify the reliability of our calculations, we first perform the calculations for the A17 phase. From the analysis of the \( D_{2h} \) point group, the zone-center optical phonon modes of A17 can be classified into

\[
\Gamma_{\text{optical}}(A17) = B_{1u}(I) + B_{2g}(I) + A_u + 2A_g(R) + B_{1g}(R) + B_{2g}(R) + 2B_{3g}(R),
\]

where \( A_u, B_{1g}, B_{2g}, \) and \( B_{3g} \) are Raman-active modes (marked with “R”), \( B_{1u} \) and \( B_{2u} \) are infrared-active modes (marked with “I”), and the \( A_g \) mode is silent. The frequencies of the Raman-active modes calculated at GGA-D2 level are 185.2, 219.8, 347.6, 410.9, 426.2, and 454.1 cm\(^{-1}\). These calculated results of Raman-active modes are in good agreement with earlier studies calculated using the same functional 29.

The phosphorus K4 crystal belongs to \( T \) point group and possesses 4 irreducible representations, namely A, E, E’, and T, respectively. The symmetries of its optical phonons at the zone center can be represented by the irreducible representations of \( T \) point group:

\[
\Gamma_{\text{optical}}(K_4) = A(R) + E(R) + E'(R) + 2T(1 + R)
\]

All the optical phonons are Raman active, and the \( T^\flat \) and \( T^\sharp \) modes are also infrared active. The frequencies of the five Raman-active modes are 230.8 (\( T^\flat \) mode), 311.3 (A mode), 392.1 (E mode), 392.1 (E’ mode), and 403.9 (\( T^\sharp \) mode) cm\(^{-1}\), respectively. The eigenvectors of the vibrational modes, as calculated at the GGA-D2 level, are illustrated in Fig. 5. These simulated Raman vibration results would be helpful to identify the K4 phosphorus structure in experiment.

The satisfactory stability of K4 phosphorus due to its ideal configuration of three-coordinated phosphorus atoms implies the possibility of the existence of this phase. Although it might be challenging to synthesize the K4 crystal, there are some relevant experimental findings that are supportive of our prediction. As mentioned above, the counterpart of K4 phosphorus, the nitrogen K4 crystal, was theoretically predicted via first principles calculations in 1992 33, and subsequently synthesized in 2004 through the polymerization of the molecular form of nitrogen at temperature above 2000 K and pressures above 110 GPa 18. Similarly, the polymerization of the molecular form of phosphorus (tetrahedral P4 molecules) was reported by Katayama et al. 34, where the transformation from molecular phosphorus to polymeric phosphorus occurred via thermal collision of the tetrahedral units of P4 molecules, and then a 3D network of three-coordinated phosphorus was formed. During such process, the experimental conditions, such as elevated temperature and environment of the tetramers play a decisive role 35. Therefore, we conclude that the phosphorus K4 crystal could be synthesized through polymerization of molecular form of phosphorus at certain conditions as is the case with cg-N 18.

**Summary**

In summary, based on first principles calculations, we predict a 3D stable phosphorus phase. We show that the K4 phosphorus structure is energetically metastable. The energy difference between the K4 phase and black phosphorus is very small (0.04 eV/atom), similar to that between diamond and graphite. Compared to other stable phosphorus allotropes at ambient conditions, such as layered phosphorus structure (A17 and violet phosphorus) and the molecular forms of phosphorus (modifications of white phosphorus), K4 phosphorus is a non-layered 3D phase with covalent bonds in all three dimensions. We demonstrate that K4 phosphorus possesses special properties different from those of black phosphorus. K4 phosphorus is stiffer than black phosphorus, and has a much larger band gap of 1.54 eV, which can be effectively tuned by applying hydrostatic pressure. K4 phosphorus exhibits much better optical absorbance than diamond silicon from 1.50 to 3.66 eV, which may have potential applications in optoelectronics. In addition, all optical phonons of K4 phosphorus at the first Brillouin zone center are
Raman active and the T1 and T2 modes are also infrared active, which can be used to identify this new allotrope of phosphorus experimentally in the future. We hope that the present theoretical study would shed new lights on discovery of novel phosphorous materials and motivate experimental efforts in this direction.

Methods
First principles calculations are performed based on density functional theory (DFT) and the projector augmented wave (PAW) method as implemented in the Vienna Ab initio Simulation Package (VASP). The electronic exchange-correlation interaction is incorporated in Perdew-Burke-Enzerhof (PBE) functional. The plane-wave cutoff energy for wave function is set to 500 eV. Lattice parameters and atomic positions are allowed to fully relax within the conjugate gradient algorithm. Since van der Waals (vdW) interactions are important for phosphorus allotropes, the GGA-D2 functional with long-range dispersion is used for the refinement of geometry in addition to the PBE functional. The Heyd-Scuseria-Ernzerhof (HSE06) hybrid functional is then used for the high accuracy of electronic structure calculations. For structure optimization, the convergence thresholds are set at $10^{-4}$ eV and $10^{-3}$ eV/Å for total energy and force component, respectively. Monkhorst-Pack k-mesh of $9 \times 9 \times 9$ is adopted to represent the first Brillouin zone. Thermal stability is studied using the Canonical ensemble (NVT) ab initio molecular dynamics (AIMD) simulations with temperature controlled by Nosé thermostat. Phonon properties are calculated using the linear response method within density functional perturbation theory as implemented in the Phonopy code. Elastic constants are determined by the finite distortion formalism.

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**Author Contributions**

Q.W. designed the project; J.L. performed the calculations; J.L., S.Z., Y.G. and Q.W. analyzed the results and wrote the manuscript.

**Additional Information**

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