The Origins of Young Stars in the Direction of the Leading Arm of the Magellanic Stream: Abundances, Kinematics, and Orbits*

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Received 2018 October 26; revised 2018 November 29; accepted 2018 November 30; published 2019 January 24

Abstract

We explore the origins of the young B-type stars found by Casetti-Dinescu et al. at the outskirts of the Milky Way disk in the sky region of the Leading Arm of the Magellanic Stream. High-resolution spectroscopic observations made with the MIKE instrument on the Magellan Clay 6.5 m telescope for nine stars are added to the previous sample analyzed by Zhang et al. We compile a sample of 15 young stars with well-determined stellar types, ages, abundances, and kinematics. With proper motions from Gaia DR2 we also derive orbits in a realistic Milky Way potential. We find that our previous radial-velocity-selected Leading Arm candidates have substantial orbital angular momentum. The substantial amount of rotational component for these stars is in contrast with the near-polar Magellanic orbit, thus making it unlikely that these stars are members of the Leading Arm. There are four large-orbital-energy stars in our sample. The star with the highest orbital energy has an age shorter than the time to disk crossing, with a birthplace at 2.5 kpc and $R_{GC} \sim 28$ kpc. Therefore, the origin of this star is uncertain. The remaining three stars have disk runaway origins with birthplaces between 12 and 25 kpc from the Galactic center. Also, the most energetic stars are more metal-poor ($[\text{Mg}/\text{H}] = -0.50 \pm 0.07$) and have larger He scatter ($\sigma_{\text{He}/\text{H}} = 0.72$) than the inner-disk ones ($[\text{Mg}/\text{H}] = 0.12 \pm 0.36$, $\sigma_{\text{He}/\text{H}} = 0.15$). While the former group’s abundance is compatible with that of the Large Magellanic Cloud, it could also reflect the metallicity gradient of the Milky Way disk and the group’s runaway status via different runaway mechanisms.

Key words: stars: abundances – stars: early-type – stars: kinematics and dynamics

1. Introduction

Recent studies attempting to describe the nearest gas-rich galaxy interaction—the Magellanic Clouds with each other and the Milky Way (MW)—are based on a wide range of new and improved observations, including proper motions (Kallivayalil et al. 2013), gas content and abundances (e.g., Fox et al. 2014, 2018; Richter et al. 2018), and low-surface-brightness features obtained with observations with the Dark Energy Camera on the Blanco 4 m telescope (e.g., Mackey et al. 2018; Nidever et al. 2018) or other facilities (Besla et al. 2016). Some of these new models (e.g., Hammer et al. 2015; Pardy et al. 2018), while succeeding at capturing some aspects of the complex physics of this interaction, also leave outstanding questions, such as the total gas content of the Magellanic Stream and its relation to present-day gas content in the Clouds (see the discussion in Pardy et al. 2018), and the formation of the Leading Arm (LA; e.g., Putman et al. 1998; Nidever et al. 2010; D’Onghia & Fox 2016). Practically, no models have been able to reproduce the morphology of the LA, which has a width of about 60° and a filamentary structure organized in four different substructures in the sky (Veitzner et al. 2012; For et al. 2013). While most recent models struggle to bring Magellanic stream material as close as $\sim 25$ kpc from the Galactic center (GC) (e.g., Pardy et al. 2018), evidence from the gas (McClure-Griffiths et al. 2008, Richter et al. 2018) and from certain young stars at the edge of the MW disk (Casetti-Dinescu et al. 2014, hereafter CD14) suggests that Magellanic Stream material as close as $\sim 20$ kpc from the GC is present. Therefore, any further observational information regarding the LA is beneficial for understanding its formation.

In this contribution we focus on further analyzing the young stars proposed by CD14 to be formed at the edge of the MW disk from material in the LA as the stream reaches and interacts with the MW’s disk (Casetti-Dinescu et al. 2014, hereafter CD14) suggests that Magellanic Stream material as close as $\sim 20$ kpc from the GC is present. Therefore, any further observational information regarding the LA is beneficial for understanding its formation.

In a subsequent study, Zhang et al. (2017, hereafter Z17) used high-resolution spectroscopy to determine the chemical compositions and kinematics for eight OB-type stars selected from CD14. They identified five B-type stars that have RVs compatible with membership to the LA and found that the abundance of these stars—as reflected by Mg—is LMC-like,
and significantly more metal-poor than that of the stars that kinematically were MW disk members.

In this study we add nine more stars from the CD14 list. Specifically, we determine abundances of He, C, N, O, Mg, Si, and S using high-resolution spectra obtained with the Magellan Inamori Kyocera Echelle (MIKE; Bernstein et al. 2003) on the Magellan Clay 6.5 m telescope. Together with the stars from Z17 we assemble a well-classified sample of 15 young stars of which 7 have RVs compatible with LA membership. The sample was chosen such that typical disk stars are present in it for comparison purposes. To the spectroscopic data, we add the proper motions from the second data release (Gaia Collaboration et al. 2018) of Gaia Collaboration et al. (2016). Based on this, we present a comprehensive orbit and abundance analysis of these stars, as well as a discussion of their origin.

The observations are described in Section 2. The method and results of the measurements of stellar parameters, abundance, and distance are described in Sections 3 and 4. The 3D kinematics and orbit calculations are presented in Section 5. Finally, the possible origins of these stars are discussed in Section 6, followed by a summary in Section 7.

2. Observations and Data Reduction

2.1. Target Selection

Similar to Z17, five stars with heliocentric distances $D < 21$ kpc and RVs $\gtrsim 90$ km s$^{-1}$ were selected from the pilot, intermediate-resolution spectral study of CD14. Four stars located in similar regions and with RV $< 90$ km s$^{-1}$ were also included for the purpose of comparison. Thus, the present study targeted nine stars, distributed above and below the Galactic plane in LA regions LA I, LA II, and LA III (Nidever et al. 2008; Venzmer et al. 2012). We designate our stars as “CD14-A**”, “CD14-B**”, and “CD14-C**” for those at Magellanic coordinates $(\lambda_M, B_M) \sim (15^\circ, -22^\circ)$, $(\lambda_M, B_M) \sim (42^\circ, -8^\circ)$, and $(\lambda_M, B_M) \sim (52^\circ, 28^\circ)$, respectively. The detailed spatial distribution of the stars can be seen in Figure 1 of CD14, and we do not repeat it here. We list our stars in Table 1, which includes the current designation, the SPM4 identification number, equatorial coordinates, and $V$ magnitudes, as well as other observational details.

| ID     | SPM ID  | R.A. (2000) | Decl. (J2000) | V     | Start (UT) | Instrument   | Exp. (s) | S/N$^a$ |
|--------|---------|-------------|---------------|-------|------------|--------------|----------|---------|
| CD14-A07 | 0200040767 | 11:41:54.77 | −78:43:59.99  | 14.9  | 06:27:11   | MIKE—Blue    | 3600 x 1 | 47      |
| CD14-A13 | 0400496347 | 12:13:19.85 | −73:58:45.18  | 16.0  | 04:20:46   | MIKE—Blue    | 3600 x 2 | 50      |
| CD14-A18 | 0690042190 | 14:19:59.97 | −71:58:0.06   | 14.0  | 07:30:48   | MIKE—Blue    | 1800 x 1 | 62      |
| CD14-B04 | 2890060268 | 10:41:17.69 | −45:07:12.93  | 14.3  | 00:17:54   | MIKE—Blue    | 2100 x 1 | 70      |
| CD14-B07 | 2900036493 | 10:55:37.37 | −46:13:04.73  | 15.4  | 01:00:31   | MIKE—Blue    | 3600 x 1 | 60      |
| CD14-B16 | 2320047259 | 11:37:27.77 | −51:26:15.22  | 14.5  | 02:05:04   | MIKE—Blue    | 1200 x 2 | 70      |
| CD14-B18 | 2320058750 | 11:45:33.72 | −50:49:45.44  | 15.0  | 02:56:00   | MIKE—Red     | 3600 x 1 | 70      |
| CD14-B19 | 2930028829 | 12:04:39.79 | −46:23:50.86  | 13.5  | 04:00:22   | MIKE—Blue    | 900 x 1  | 68      |
| CD14-C06 | 5540048652 | 08:59:47.73 | −24:00:47.19  | 14.1  | 23:31:03   | MIKE—Red     | 1200 x 2 | 71      |

Notes.

$^a$ The observation was taken on 2016 April 30.

$^b$ The S/Ns were evaluated at 4000 and 5800 Å for the blue and red sides, respectively.
Table 1

| ID         | $T_{\text{eff}}$ (K) | log $g$ | log $\text{[Fe/H]}$ | $\nu$ sin $i$ (km s$^{-1}$) |
|------------|----------------------|---------|---------------------|-----------------------------|
| CD14-A07   | 15,000 ± 900         | 4.25    | 14.642 ± 204       | 4.13 ± 0.06                 |
| CD14-A13   | 14,000 ± 1500        | 4.23    | 12,525 ± 102       | 4.28 ± 0.04                 |
| CD14-A18   | 14,800 ± 800         | 4.40    | 14,452 ± 141       | 4.38 ± 0.05                 |
| CD14-B04   | 15,750 ± 900         | 4.26    | 15,975 ± 111       | 4.31 ± 0.04                 |
| CD14-B07   | 20,500 ± 950         | 4.56    | 20,211 ± 321       | 4.63 ± 0.05                 |
| CD14-B16   | 16,500 ± 750         | 4.30    | 15,874 ± 147       | 4.49 ± 0.03                 |
| CD14-B18   | 15,000 ± 900         | 4.28    | 14,735 ± 177       | 4.23 ± 0.05                 |
| CD14-B19   | 16,850 ± 700         | 4.15    | 16,612 ± 189       | 4.14 ± 0.04                 |
| CD14-C06   | 15,000 ± 600         | 4.12    | 14,254 ± 108       | 4.11 ± 0.03                 |

Notes.

a The parameter I values are measured by NLTE grids of BSTARS.
b The parameter II values are measured by an independent grid (LTE for B stars) and the analysis code of LINFOR.

2.2. High-resolution Spectroscopy and Data Reduction

High-resolution spectra of the nine target stars were obtained with the MIKE instrument on the 6.5 m Clay telescope in 2016 April. The setup gave a resolution of $R \sim 33,000$ and $R \sim 29,000$ for the blue (3200–5000 Å) and the red (4900–10000 Å) sides, respectively. The average seeing and airmass during the observations were 0.8 and 1.2. Table 1 summarizes the observational details, exposure times, and the achieved S/Ns of the spectra in the range of 4000 and 5800 Å.

The standard Carnegie Python distribution\(^{11}\) (Kelson et al. 2000; Kelson 2003) routines for MIKE were used for data reduction, including order identification, wavelength calibration, flat-field correction, background subtraction, one-dimensional spectra extraction, and flux calibration. The RVs of our targets were measured after continuum rectification, by cross-correlation with a synthetic template whose temperature

\(^{11}\) http://code.obs.carnegiescience.edu
and gravity is similar to those of the targets. The standard algorithm described by Tonry & Davis (1979) was adopted, and the synthetic template was taken from the synthetic spectral library of Munari et al. (2005). Two telluric oxygen A bands, the (0,1) transition at 6884 Å, and the (0,0) transition band at 7621 Å, were also adopted to correct the RV zero-point, where the molecular data were taken from the High-resolution TRANsmission molecular absorption (HITRAN) database.12 Examples of a portion of the spectra are shown in Figure 1 and the obtained heliocentric RVs are listed in Table 4.

3. Measurements

We used the same methodology described in Z17 to measure the stellar parameters, to determine the elemental abundances, and to estimate the corresponding uncertainties. Here, we briefly describe the procedures, but more details can be found in Z17 and references therein.

3.1. Stellar Atmosphere Parameters

The stellar model atmospheres we employed for our target stars are interpolated from comprehensive grids of metal line-blanketed, non-local thermodynamic equilibrium (NLTE), plane-parallel, hydrostatic model atmospheres of B-type stars (Lanz & Hubeny 2007, BSTAR2006), with opacity sampling, which is a simple Monte Carlo-like sampling of the superline cross sections and solar abundance. For all our targets, the solar-metallicity ($Z/Z_\odot = 1$) grid is adopted. However, for CD14-A07, CD14-A18, CD14-B04, and CD14-B18, which show strong helium lines, an additional super-solar-metallicity ($Z/Z_\odot = 2$) grid was also employed for comparison purposes.

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12 https://www.cfa.harvard.edu/hitran/
We have used the spectrum synthesis program SYNSPEC, developed by Ivan Hubeny & Thierry Lanz. The automatic NLTE treatment mode of SYNSPEC is employed, in which the program automatically decides which levels are treated in NLTE, and assigns proper NLTE populations to the lower and upper levels of a given transition. The input model atoms and ions (Lanz & Hubeny 2003) are provided by TLUSTY. Balmer series from H\textsubscript{9} to H\textsubscript{10}—except H\textsubscript{0}, which is blended with Ca I H line—and nine He I lines in ranges of 4000–6000 Å are included in the spectrum synthesis procedure for the measurements of stellar parameters. The best fitted stellar parameters are found by a $\chi^2$ test, and the uncertainties estimated from the $\chi^2$ statistics were multiplied by three to obtain the final error, because the errors propagated from the data reduction procedure, such as sky subtraction and the normalization, cannot be neglected (Moni Bidin et al. 2012, hereafter MB12). The synthetic spectra with the best-fit stellar parameters are shown in Figure 2.

Similar to Z17, we also perform a separate, independent measuring process using different grids of model atmospheres and analysis code (specifically, that in MB12) for the stellar parameters, in order to test the reliability of our stellar parameter results. The grid of model spectra is computed with Lemke’s version of the LINFOR program (developed originally by Holweger, Steffen, and Steenbock at Kiel University), which is based on local thermodynamic equilibrium (LTE) model atmospheres of ATLAS9 (Kurucz 1993). In this analysis, Balmer series from H\textsubscript{9} to H12, four He I lines in ranges of 4000–5000 Å, were fitted simultaneously. For more details on the grids of model atmospheres and analysis code, we refer the reader to MB12 and references therein.

Comparisons of the atmospheric parameters between the two independent analyses are also presented in Table 2 and Figure 3. For convenience, we label our default spectral analysis process described in the previous subsection and the separate one à la MB12 as “SA I” and “SA II,” respectively. The two sets of parameter results reveal an overall good agreement, apart from one problematic case—CD14-A13. The SA II temperature of CD14-A13 is lower than the cooler end of the BSTARS grid at 15,000 K. Although model atmospheres could be extrapolated from the grid, this is only valid when $T_{\text{eff}}$ and $\log g$ values are not far from the limits of the grid, that is, small compared to the grid steps. Therefore, the parameter values of SA I are the best choice given the possible grid ranges of BSTARS. With the exclusion of CD14-A13, the mean difference is a mere 330 K in $T_{\text{eff}}$ and −0.015 dex in $\log g$. This offset in $T_{\text{eff}}$ seems systematic, although it is negligible compared to the errors. Rotational velocities are in good agreement (mean difference of 2 km s$^{-1}$). For consistency, we will adopt SA I results for the subsequent abundance analysis.

### 3.2. Abundances and Uncertainties

Using the stellar parameters obtained with the stellar model atmospheres and synthesis code described in the previous subsection, we proceed to determine the elemental abundances of carbon, nitrogen, oxygen, neon, magnesium, silicon, and sulfur by fitting observed lines with synthetic spectra. We only determine the iron abundance for CD14-B07 and CD14-B19, as there are no clear, unambiguous iron lines that can be detected for the other sample stars. For each element, synthetic spectra are produced using the stellar parameters measured with the SA I procedure and with various abundance values. These synthetic spectra are then compared with the observed ones. A $\chi^2$ test was employed to establish the best fit for detected absorption lines. The mean abundance of each element is derived from all lines that can be clearly detected. For stars in which feature lines for certain elements cannot be detected clearly, the feature that is at the position of the theoretical absorption line was fitted, and the maximum value for this element that could fit the spectrum is considered as the upper limit for its abundance. During the procedure, the solar composition from Grevesse & Sauval (1998) is adopted. The atomic line data of elements are the same as Z17, that is, they are mainly selected from Kurucz & Bell (1995, K95) and the NIST Atomic Spectra Database.

The overall abundance uncertainty is estimated by summing in quadrature the uncertainties from the following sources:

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13 http://nova.astro.umd.edu/Synspec49/synspec.html
14 http://nova.astro.umd.edu/Tlusty2002/tlusty-frames-data.html
15 http://www.sternwarte.uni-erlangen.de/~ai26/linfit/linfor.html
16 http://www.nist.gov/pml/data/asd.cfm
Uncertainties from the observational error and the fitting procedure.

(2) Uncertainties from atomic data. In the present study, an uncertainty of 10% in $\log g$ was also adopted to explore its effect on the abundance, which results in an error of 0.02 dex on average.

(3) Errors in the continuum rectification. In the worst case, the error in continuum rectification was estimated to be 5%, which results in a change of the abundance of up to 0.05 dex.

(4) Uncertainties propagated from errors of the stellar parameters.

(5) The scatter in the abundance determinations from different lines gives another estimate of the uncertainty.

More details can be found in Z17. Synthetic spectra with error ranges that are generated for the best-fit parameters of sample stars are shown in Figures 4–6, using representative C II, N II, O II, Mg II 4481 Å, Si II, and S II as examples.

3.3. Distances and Ages

Gaia DR2 provides parallaxes for ~1.7 billion stars in the MW. However, the relative errors are larger than 25% for most of our sample stars. For this reason, we instead adopt the following spectroscopic distances for the stars.

The present sample is contaminated by horizontal-branch (HB) stars (see Section 4.2). For the true main-sequence (MS) stars of the sample, we determine their age via comparison of their position in the temperature–gravity plane with PARSEC isochrones (Bressan et al. 2012). The uncertainty in this estimate is given by the error box associated with the stellar parameters. The isochrones also return an estimate of the absolute magnitude $M_V$ of each star. For the HB stars, we adopt the HB evolution tracks (PARSEC) to estimate the masses and luminosities, and possibly the evolution time after helium flash.

Finally, we calculate the true distance moduli ($m - M$)$_0$ from the apparent $V$ magnitudes of the SPM4 catalog, de-reddened with the Schlegel et al. (1998) maps as corrected by Bonifacio et al. (2000), and assuming a standard reddening law with $R = E(B-V)/E(B-V) = 3.1$. We have verified that the results are not affected within errors by the use of the Schlafly & Finkbeiner (2011) calibration, and that the distance moduli obtained using the 2MASS $J$ and $K$ magnitudes are consistent with those obtained in $V$, with no systematics.

4. Spectroscopic Results

4.1. Abundances

The measured stellar parameters and abundances are summarized in Tables 2 and 3 and shown in Figure 9. For stars CD14-A07, CD14-A18, CD14-B04, and CD14-B18, whose detected abundance patterns are all super-solar, the abundance values derived by the metal-rich model atmospheres...
(Z/Z_⊙ = 2) are also shown (see yellow symbols in Figure 9). Abundances derived using the metal-rich model atmosphere appear slightly, but systematically, lower than those obtained using the solar-metallicity model atmosphere; however, values are consistent within their 1σ errors. For a better analysis, in the following discussion we will adopt the results from metal-rich model atmospheres for these four stars. Except for CD14-B07 and CD14-B16, He and C abundances of the stars are super-solar or near-solar, within uncertainties. Because of rapid v sin i, S II lines combined with He I 4143 Å cannot be separated clearly for most of the targets. Only an upper limit of S abundance can be estimated by fitting line wings of He I 4143 Å. We proceed now to discuss individual stellar results.

4.1.1. CD14-A07, CD14-A18, CD14-B04, and CD14-B18

Among the four metal-rich stars, CD14-A07, CD14-B04, and CD14-B18 have large rotational velocities. Weak lines cannot be detected due to the broadening effects of rotation. Only strong lines such as He I 5587 Å, Mg II 4487 Å, and C II 4266 Å can be detected and used in the abundance determination. The measurement of stellar parameters and of elemental abundances is made more difficult by having so few spectral features. Therefore, we do not have sufficient information for a reliable abundance determination for these three stars. Also, during the line fitting, rotational broadening implies that the line profile is insensitive to the variance of the abundance, which leads to relatively large uncertainties of the abundances for the three stars.

Although no usable iron lines can be detected for [Fe/H] measurement, we assume their [Fe/H] is about 0.3 dex because the model atmospheres used for these four stars is Z/Z_⊙ = 2, while [Mg/Fe], [Si/Fe], and [S/Fe] show no obvious enhancements.

4.1.2. CD14-B07 and CD14-B19

Most of the absorption lines used in our chemical analysis are detected for these two stars because they are not sensibly affected by the broadening effects of rotation. Strong N II 3995 Å can be detected in the spectra of CD14-B07 and CD14-B19, implying a significant nitrogen enhancement. The abundance of Mg for the two stars is sub-solar, while Si and S abundances are super-solar and sub-solar for CD14-B07 and CD14-B19, respectively. The Fe III line at 4421 Å can be detected for the two stars, from which we estimated their [Fe III/H] is about 0.42 ± 0.12 dex and 0.05 ± 0.18 dex, respectively. Besides these lines, one phosphorus line P II 4588 Å can be detected in the spectra of CD14-B07, yielding an abundance of [P II/H] = 1.23 ± 0.10 dex.

4.1.3. CD14-A13, CD14-B16, CD14-C06

There are no He and C enhancements in CD14-A13. The abundance uncertainties of the star are relatively large because
### Table 3
Element Abundance Results of the Targets

|                  | CD14-A07 | CD14-A13 | CD14-A18 | CD14-B04 | CD14-B07 | CD14-B16 | CD14-B18 | CD14-B19 | CD14-C06 |
|------------------|----------|----------|----------|----------|----------|----------|----------|----------|----------|
| log ε (He)       | 11.45 ± 0.17 | 10.93 ± 0.40 | 11.31 ± 0.18 | 11.39 ± 0.18 | 10.40 ± 0.13 | 9.88 ± 0.15 | 11.27 ± 0.22 | 11.43 ± 0.12 | 11.51 ± 0.17 |
| log ε (C)        | 9.31 ± 0.18 | 8.15 ± 0.28 | 8.94 ± 0.17 | 9.12 ± 0.23 | 7.84 ± 0.12 | 8.06 ± 0.16 | 9.07 ± 0.20 | 8.95 ± 0.14 | 9.03 ± 0.20 |
| log ε (N)        | ...       | ...       | ...       | ...       | 8.66 ± 0.13 | <7.82      | ...       | 8.29 ± 0.15 | ...       |
| log ε (O)        | ...       | ...       | ...       | ...       | 8.38 ± 0.20 | <8.83      | ...       | 9.34 ± 0.13 | ...       |
| log ε (Mg)       | 8.09 ± 0.15 | 7.93 ± 0.23 | 7.92 ± 0.15 | 7.87 ± 0.19 | 7.23 ± 0.15 | 7.07 ± 0.18 | 8.01 ± 0.23 | 7.25 ± 0.11 | 7.44 ± 0.15 |
| log ε (Si)       | 8.01 ± 0.18 | 7.50 ± 0.26 | 8.00 ± 0.16 | 7.62 ± 0.18 | 7.89 ± 0.16 | 7.82 ± 0.18 | 7.88 ± 0.32 | 7.18 ± 0.15 | 7.77 ± 0.17 |
| log ε (P)        | ...       | ...       | ...       | ...       | 6.68 ± 0.10 | ...        | ...       | ...       | ...       |
| log ε (S)        | <7.98     | <7.33     | 7.73 ± 0.15 | <7.43     | 7.40 ± 0.13 | 7.17 ± 0.20 | ...       | 6.98 ± 0.15 | <7.53     |
| log ε (Fe)       | ...       | ...       | ...       | ...       | 7.92 ± 0.12 | ...        | ...       | 7.60 ± 0.18 | ...       |
| [He/H]           | 0.52 ± 0.17 | 0.00 ± 0.40 | 0.38 ± 0.18 | 0.46 ± 0.18 | −0.53 ± 0.13 | −1.05 ± 0.15 | 0.34 ± 0.22 | 0.50 ± 0.12 | 0.58 ± 0.17 |
| [C/H]            | 0.78 ± 0.18 | −0.38 ± 0.28 | 0.41 ± 0.17 | 0.59 ± 0.23 | −0.69 ± 0.12 | −0.47 ± 0.16 | 0.54 ± 0.20 | 0.42 ± 0.14 | 0.50 ± 0.20 |
| [N/H]            | ...       | ...       | ...       | ...       | 0.14 ± 0.13 | <−0.10      | ...       | 0.37 ± 0.15 | ...       |
| [O/H]            | ...       | ...       | ...       | ...       | −0.45 ± 0.20 | <0.00       | ...       | 0.51 ± 0.13 | ...       |
| [Mg/H]           | 0.51 ± 0.15 | 0.35 ± 0.23 | 0.34 ± 0.15 | 0.29 ± 0.19 | −0.35 ± 0.15 | −0.51 ± 0.18 | 0.43 ± 0.23 | −0.25 ± 0.11 | −0.14 ± 0.15 |
| [Si/H]           | 0.46 ± 0.18 | −0.05 ± 0.26 | 0.45 ± 0.16 | 0.07 ± 0.18 | 0.34 ± 0.16 | 0.27 ± 0.18 | 0.33 ± 0.32 | −0.37 ± 0.15 | 0.22 ± 0.17 |
| [P/H]            | ...       | ...       | ...       | ...       | 1.23 ± 0.10 | ...        | ...       | ...       | ...       |
| [S/H]            | <0.65     | <0.00     | 0.40 ± 0.15 | <0.10     | 0.07 ± 0.13 | −0.16 ± 0.20 | ...       | −0.35 ± 0.15 | <0.2      |
| [Fe/H]           | ...       | ...       | ...       | ...       | 0.42 ± 0.12 | ...        | ...       | 0.05 ± 0.18 | ...       |

**Notes.**

a Solar compositions are taken from Grevesse & Sauval (1998).

b With $Z/Z_\odot = 2$ model atmosphere.

c $\log \varepsilon (X) = \log \frac{N_X}{N_H} + 12$. 
of the large error bars propagated from the stellar parameters. CD14-B16 appears depleted in He and C, and its Mg abundance is similar to the average value of the LMC. Similar to the four metal-rich stars mentioned above, CD14-C06 shows He and C enhancements. The remainder of its abundance values are slightly sub- or near-solar, within the uncertainties.

### 4.2. Stellar Classification

The log g values of the present targets are larger than 4.0 dex in the $T_{\text{eff}}$ range of 15,000–21,000 K, in which hot HB and MS B stars occupy overlapping regions of the $T_{\text{eff}} - \log g$ diagram. This can be seen in Figure 7, where the targets are plotted in the $T_{\text{eff}} - \log g$ plane along with HB and MS evolutionary tracks (PARSEC; Bressan et al. 2012, 2013; Chen et al. 2014). In this figure, the stars are separated into two groups according to their metallicities, that is, CD14-A07, CD14-A18, CD14-B04, and CD14-B18 belong to $Z/Z_0 = 2$ (red symbols), while the remaining five stars belong to $Z/Z_0 = 1$ (black symbols). To classify each of these stars, abundance and $v \sin i$ info are also considered. A star-by-star analysis reveals the following:

1. **CD14-A13, CD14-B19, and CD14-C06.** The three stars lie above the ZAMS of solar metallicity. Except for CD14-B19, the other two stars are found between the canonical zero-age and terminal-age HB (ZAHB and

Figure 7. The location of target stars with respect to PARSEC evolution tracks (Bressan et al. 2012, 2013; Chen et al. 2014) in the temperature–gravity plane. The black and red symbols/lines represent $Z/Z_0 = 1$ and $Z/Z_0 = 2$, respectively. The zero-age horizontal-branch (ZAHB) and terminal-age horizontal-branch (TAHB) tracks are shown with solid lines, while the zero-age main-sequence (ZAMS) stars are represented by dashed lines. The upward and downward triangles connected by a dotted line represent the location of CD14-A13 with $T_{\text{eff}}$ and log $g$ from SA I and SA II, respectively. (a) HB evolution tracks are overplotted, corresponding to models with $M = 0.490 M_\odot$, $0.495 M_\odot$, and $0.505 M_\odot$ for $Z/Z_0 = 1$, and $M = 0.475 M_\odot$, $0.480 M_\odot$, and $0.490 M_\odot$ for $Z/Z_0 = 2$ (bottom left to upper right). Nodes on the tracks indicate the evolution times. The time interval is 10 Myr. (b) Similar to (a), but MS evolution tracks are overplotted, corresponding to models with $M = 3.0 M_\odot$, $3.6 M_\odot$, $4.2 M_\odot$, $4.8 M_\odot$, and $5.4 M_\odot$ for $Z/Z_0 = 1$, and $M = 3.4 M_\odot$, $4.0 M_\odot$, $4.6 M_\odot$, $5.2 M_\odot$, and $5.8 M_\odot$ for $Z/Z_0 = 2$ (right to left).

### Table 4 Distance, Age, and Radial Velocities

| ID           | (m–M)$_0$ (mag) | Age (MS Myr) | HB (Gyr) | RV (km s$^{-1}$) |
|--------------|----------------|--------------|----------|-----------------|
| CD14-A07     | 14.54 ± 0.30   | 13 ± 12      | ...      | 94 ± 5          |
| CD14-A13     | 15.20 ± 0.47   | 45 ± 40      | ...      | −2 ± 7          |
| CD14-A18     | 14.76 ± 0.21   | 69 ± 20      | ...      | −12 ± 6         |
| CD14-B04     | 13.05 ± 0.28   | 10 ± 8$^d$   | ...      | −20 ± 5         |
| CD14-B07     | 14.21 ± 0.22   | 6 ± 4        | >3       | 136 ± 5         |
| CD14-B16     | 12.62 ± 0.50   | >3           | 132 ± 4  | 136 ± 5         |
| CD14-B18     | 11.62 ± 0.40   | >3           | 9 ± 5    | 7 ± 5           |
| CD14-B19     | 14.67 ± 0.23   | 9 ± 2        | >3       | 81 ± 3          |
| CD14-C06     | 14.13 ± 0.23   | 32 ± 6       | >3       | 111 ± 7         |
| CD14-C06     | 14.33 ± 0.18   | 59 ± 8       | >3       | 99 ± 6          |

Notes:

$^a$ $A_V$ is calculated assuming an extinction-to-reddening ratio $A_V/E(B−V) = 3.1$. $E(B−V)$ is from Schlafly & Finkbeiner (2011).

$^b$ The heliocentric radial velocity values calculated in this work.

$^c$ PARSEC isochrones with metallicities of $Z/Z_0 = 2$ are used.

$^d$ $M_V$ of CD14-A18 is assumed to be at the same temperature from the ZAMS.
TAHB, respectively), where post-helium-flash stars spend 99% of their helium-burning lifetime. MB12 found that the He abundance of HB stars strongly increases with $T_{\text{eff}}$ in the range $15,000 \text{ K} < T_{\text{eff}} < 20,000 \text{ K}$, and reaches the minimum value of $\log \frac{\text{He}}{\text{H}} \sim -2.5$ around $15,000 \text{ K}$ because of the maximum efficiency of diffusion at this temperature. There might, however, be an offset of $-0.5$ dex for field HB stars compared to the cluster measurements (Mont Bidin et al. 2009). The He abundances of the three stars are all super-solar, which suggests that they are not HB after all. We then check the evolutionary state of CD14-B19. If CD14-B19 were an HB star, it would have evolved off the TAHB phase, as shown in Figure 7, its He abundance would show depletion to some degree. However, its super-solar [He/H] value ($0.50 \pm 0.12$) indicates otherwise. Also, from Figure 7, one can see that CD14-A13 and CD14-C06 will pass through an evolution phase from $T_{\text{eff}} = 15,000 \text{ K}$, $\log g = 4.3$ to $T_{\text{eff}} = 20,000 \text{ K}$, $\log g = 4.0$ in $\sim 10 \text{ Myr}$ if they are HB stars, whereas the time interval is $\sim 100 \text{ Myr}$ if they are MS stars. Therefore, it is more probable to observe these two stars if they are in their MS evolutionary phase. In addition, both CD14-A13 and CD14-C06 show rapid rotation ($v \sin i > 50 \text{ km s}^{-1}$), which are seldom found in evolved objects (Figure 8 of Behr 2003). We conclude that these three stars are in the MS phase.

2. CD14-B07. The locus of the star is far below the ZAMS, and above the region between ZAHB and TAHB. Considering the abundances we see that: (1) both its magnesium and helium are depleted; (2) iron abundance is enhanced; (3) phosphorus is strongly enhanced. These features are found in hot HB stars in globular clusters as well as in field stars (e.g., Behr et al. 1999; Moehler et al. 1999; Behr 2003) because of chemical diffusion. That is, iron abundance may be enhanced, while magnesium and helium show depletion because of the radiative levitation of most metal species, and gravitational settling of helium (Behr et al. 1999). On the contrary, MS stars are usually not affected by atmospheric diffusion. Combining these results with the slow rotation, we classify CD14-B07 as a star with $0.49 M_\odot$ that has just evolved past the post-helium-flash phase.

3. CD14-B16. This star lies slightly above the solar-metallicity ZAMS. From Figure 7, it can be seen that the time intervals CD14-B16 spends in HB and MS tracks are comparable. Although the $v \sin i$ of the star is $55 \pm 10 \text{ km s}^{-1}$—which is larger than the upper limits ($35 \text{ km s}^{-1}$) of field HB stars given by Behr (2003)—the value is still acceptable for an HB star if a 2σ uncertainty is considered. Moreover, its low He abundance ($\text{[He/H]} = -1.05 \pm 0.15$) implies that it could be a post-helium-flash star evolving off the HB. In addition, its abundance patterns are similar to those of CD14-B07 (see Figure 9). The star is therefore marked as an HB object.

4. CD14-A07, CD14-A18, CD14-B04, and CD14-B18. These four metal-rich stars are all in the post-helium-flash region. Although CD14-A18 lies slightly below the metal-rich ZAMS, its $v \sin i$ of $65 \pm 10 \text{ km s}^{-1}$ and super-solar abundance patterns suggest it is not an HB star. Moreover, if it is actually in the MS phase, from Section 3.3, it is a metal-rich star located at $R_\odot = 6.26 \text{ kpc}$, which agrees well with the Galactic radial metallicity gradient (e.g., Magrini et al. 2009; Sales Silva et al. 2016). Therefore, we mark this target as a MS star. The other three stars, CD14-A07, CD14-B04, and CD14-B18, lie on the metal-rich ZAMS. They all rotate very fast, with $v \sin i > 100 \text{ km s}^{-1}$, and are therefore very young MS stars.

In summary, we find that the sample is contaminated by two HB objects, CD14-B07 and CD14-B16. CD14-A18 is an MS star because of its fast $v \sin i$ and super-solar He abundance values, although it lies below the ZAMS models. Therefore, the seven objects in the MS evolutionary phase are further considered in the following discussion.

4.3. Radial Velocities

The results of our RV measurements are given in Table 4. The error bar of CD14-B18 is relatively large because its absorption lines are broadened by large $v \sin i$ values. A comparison between our RV measurements and the first-epoch values of CD14 is shown in the same table. The two sets of results exhibit good agreement. The RV differences will be taken as ours—CD14 hereafter, with errors being the quadratic sum of the uncertainties of the two measurements. All stars show variations lower than the average 0.6σ error, and small compared to the huge variations typical of massive B-type binary stars. The only exception is star CD14-B19, whose variation of $\Delta \text{RV} = 15 \text{ km s}^{-1}$ is significant at the 2σ level.

From Section 4.2, it can be seen that seven stars are classified as B-type MS stars. It is known that the fraction of close binaries among B-type stars is greater than 50% (e.g., Kouwenhoven et al. 2007). To check the fraction of binary detections given our two-epoch observations and a detection threshold of $\Delta \text{RV} = 20 \text{ km s}^{-1}$, we adopt the same model as the one in Z17 (and references therein), and find that the probability of our null detection out of seven stars is 14% (8.7%) if the underlying binary fraction is 50% (60%). We did not find any binaries in the sample of Z17. If the current and the Z17 samples are combined, the fraction of null detection turns out to be 2.0%(0.8%).
Considering that the \(\Delta RV\) value of CD14-B19 is close to the detection threshold adopted here, and the rather low null detection fraction obtained in this calculation, it is possible that CD14-B19 is a binary, and the adopted threshold of \(\Delta RV = 20\) km s\(^{-1}\) is a rather conservative value.

### 4.4 Distances and Ages

The results of distance and age estimates are also presented in Table 4 and Figure 8. We used super-solar-metallicity isochrones with \(Z/Z_\odot = 2\) for CD14-A07, CD14-B04, and CD14-B18. The results show that the three targets are extremely young stars (<13 Myr) located at heliocentric distance \(d = 4\) to \(d = 8\) kpc. Star CD14-A18 would also require the same treatment, but its location in the \(T_\text{eff}-\log g\) plane is slightly below the the metal-rich zero-age MS (ZAMS) model (Bressan et al. 2012, 2013) even if the \(1\sigma\) error of \(\log g\) is considered. Assuming it is a MS star, its theoretical magnitude is determined at the temperature from the metal-rich ZAMS. This results in a heliocentric distance of 8.5 kpc for CD14-A18.

### 5. 3D Kinematics and Orbit Determination

#### 5.1 Proper Motions

With our measured RVs and spectroscopic distances, and with proper motions extracted from the Gaia DR2 catalog it is possible to determine complete 3D kinematic information. We do so for the stars in the present study and in Z17. One distant star from CD14, namely CD14-A19, is also included in the following calculation, although no high-resolution spectroscopic information is available. The proper motions and parallaxes from the Gaia DR2 catalog are listed in Table 5:
| ID       | Gaia ID           | RV (km s⁻¹) | Distance (kpc) |plx (mas) | μₐ (mas yr⁻¹) | Δμₐ (mas yr⁻¹) | μ₉ (mas yr⁻¹) | Δμ₉ (mas yr⁻¹) | l (deg) | b (deg) | μₐ (mas yr⁻¹) | Δμₐ (mas yr⁻¹) |
|----------|-------------------|-------------|----------------|----------|---------------|---------------|---------------|---------------|---------|---------|-------------|---------------|
| CD14-A05 | 5226326178138343168 | 133 ± 6     | 19.95 ± 2.96   | 0.025 ± 0.033 | -4.32 ± 0.07 | 1.74 ± 0.05   | 295.75        | -11.78        | -4.66 ± 0.04 | -0.06 ± 0.07 |
| CD14-A07 | 520092659627780352  | 94 ± 8      | 8.09 ± 1.20    | 0.072 ± 0.036 | -8.52 ± 0.06 | -0.51 ± 0.05  | 299.44        | -16.33        | -8.04 ± 0.05 | -2.85 ± 0.07 |
| CD14-A08a| 5224454951082866048 | 48 ± 8      | ...            | 0.202 ± 0.034 | -5.69 ± 0.06 | -0.07 ± 0.06  | 298.79        | -13.92        | -5.46 ± 0.04 | -1.60 ± 0.07 |
| CD14-A11a| 52249859480348416   | 71 ± 4      | 13.80 ± 3.57   | 0.010 ± 0.028 | -2.91 ± 0.05 | 0.05 ± 0.04   | 298.96        | -12.88        | -2.84 ± 0.04 | -0.66 ± 0.06 |
| CD14-A12b| 526633985831636480  | 33 ± 7      | 10.47 ± 1.01   | 0.090 ± 0.025 | -7.21 ± 0.05 | -1.35 ± 0.04  | 298.81        | -12.09        | -6.67 ± 0.04 | -3.03 ± 0.05 |
| CD14-A13l| 584149623121376512  | -2 ± 6      | 8.95 ± 0.86    | 0.003 ± 0.040 | -6.17 ± 0.07 | 1.44 ± 0.06   | 300.26        | -11.30        | -6.32 ± 0.06 | 0.49 ± 0.07  |
| CD14-A15  | 5838307413790168960 | 162 ± 6     | 10.96 ± 1.62   | 0.017 ± 0.033 | -4.88 ± 0.06 | 1.15 ± 0.05   | 300.74        | -11.82        | -4.99 ± 0.05 | 0.52 ± 0.05  |
| CD14-A18  | 579716195829770496  | -20 ± 6     | 4.07 ± 0.60    | 0.077 ± 0.030 | -6.24 ± 0.04 | -3.93 ± 0.04  | 309.74        | -10.28        | -7.21 ± 0.05 | -1.57 ± 0.03 |
| CD14-A19  | 5793052064945812480 | 234 ± 6     | 30.20 ± 15.30  | 0.101 ± 0.065 | -2.18 ± 0.19 | -1.62 ± 0.11  | 309.35        | -14.09        | -2.65 ± 0.22 | -0.61 ± 0.03 |
| CD14-B02b| 5414441420674317184 | 170 ± 4     | 11.48 ± 1.11   | 0.034 ± 0.029 | -3.65 ± 0.04 | 1.47 ± 0.04   | 275.29        | 10.56         | -3.84 ± 0.01 | -0.89 ± 0.06 |
| CD14-B03c| 5368186306521663248 | 227 ± 5     | 15.14 ± 0.71   | 0.093 ± 0.040 | -4.10 ± 0.06 | 1.92 ± 0.06   | 277.72        | 13.06         | -4.50 ± 0.02 | -0.48 ± 0.08 |
| CD14-B04b| 5369285353939460480 | 132 ± 5     | 6.95 ± 0.67    | 0.128 ± 0.038 | -4.70 ± 0.06 | 3.61 ± 0.06   | 280.11        | 11.91         | -5.86 ± 0.02 | 0.85 ± 0.08  |
| CD14-B07b| 5363434938821203712 | 19 ± 4      | 3.34 ± 0.87    | 0.238 ± 0.048 | -4.44 ± 0.08 | 0.20 ± 0.07   | 282.90        | 12.10         | -4.08 ± 0.04 | -1.77 ± 0.10 |
| CD14-B14b| 537367743735021760  | 202 ± 5     | 20.89 ± 2.02   | -0.003 ± 0.036 | -5.46 ± 0.05 | 1.42 ± 0.05   | 288.78        | 12.20         | -5.62 ± 0.04 | -0.45 ± 0.07 |
| CD14-B16b| 536973919404171520  | 9 ± 5       | 2.11 ± 0.43    | 0.072 ± 0.035 | -7.82 ± 0.05 | 2.04 ± 0.04   | 291.36        | 9.76          | -8.07 ± 0.03 | -0.29 ± 0.05 |
| CD14-B18b| 536992529252766000  | 92 ± 3      | 8.59 ± 0.83    | 0.032 ± 0.038 | -5.37 ± 0.05 | 0.59 ± 0.04   | 292.43        | 10.69         | -5.34 ± 0.04 | -0.81 ± 0.05 |
| CD14-B19b| 614325809937783424  | 96 ± 7      | 6.70 ± 0.65    | 0.081 ± 0.049 | -4.11 ± 0.06 | 3.32 ± 0.04   | 294.58        | 15.72         | -4.66 ± 0.05 | 2.49 ± 0.04  |
| CD14-C06b| 5654350872120057856  | 96 ± 6      | 7.34 ± 0.71    | -0.008 ± 0.045 | -0.52 ± 0.07 | 0.01 ± 0.05   | 249.94        | 14.26         | -0.34 ± 0.01 | -0.40 ± 0.08 |

Notes:
1. Spectroscopic distance determined in the present study and Z17.
2. Sample stars from Z17.
3. Sample stars from CD14.
there are nine stars from the current study, eight from Z17 and one from CD14. However, for the velocity and orbit analysis, we will discard stars CD14-B07 and CD14-B16 from this study, because they are HB stars, and star CD14-A08 from Z17 because it has an ambiguous classification. We are therefore left with 15 stars for the kinematic and orbit analysis.

5.2. Velocity Components in the Galactic Cylindrical Coordinate System

We first calculate the rectangular velocity components with respect to the GC by adopting a value of 238 km s$^{-1}$ (Schönrich 2012) for the local standard of rest ($V_{LSR}$) and a solar motion of (+9.58, +10.52, +7.01) km s$^{-1}$ (Tian et al. 2015) in ($U$, $V$, $W$), which are defined in a right-handed Galactic system with $U$ pointing toward the GC, $V$ in the direction of rotation, and $W$ toward the north Galactic pole (Dehnen & Binney 1998). The kinematic parameters, including distance and velocities, are then transformed to the Galactic cylindrical coordinate system:

$$
\phi = \arctan \left[ \frac{D \cos(b) \sin(l)}{R_{\odot} - D \cos(b) \cos(l)} \right],
$$

$$
v_{\text{GC}} = \Pi \hat{r}_{\odot} + \Theta \hat{\phi} + Z \hat{z},
$$

where ($l$, $b$) are the Galactic longitude and latitude, $R_{\odot}$ = 8.34 kpc (Reid et al. 2014) is the adopted distance to the GC, $\Pi = V_{\odot} \equiv \dot{R}$, $\Theta = V_{\odot} \equiv R_{\odot}$, and $Z = V_{\odot} \equiv \dot{z}$.

The errors in the velocity components are estimated by Monte Carlo sampling assuming a multivariate Gaussian distribution. For each star, 1000 sets of variables used in the velocity calculation are generated from a six-dimension multivariate Gaussian distribution, $N(\mu, \Sigma)$, where $\mu = [\text{ra}, \text{dec}, \mu_l, \mu_b, D, v_{\text{rad}}]$ are observed values, and $\Sigma$ is the covariance matrix derived from the observational errors and the correlation coefficients provided by the Gaia DR2 catalog.

The velocity components are listed in Table 6. Along with these components, we list the heliocentric RV, which was used in Z17 and CD14 as the criterion for LA membership.

Specifically, stars with RV $> 100$ km s$^{-1}$ were considered members; in this Table we highlight these values in bold. A total of seven stars fit this criterion. It is illuminating to inspect the cylindrical velocity components of these seven stars: all of them have large $V_{\phi}$, indicating a strong rotation component. This is unexpected for potential LA members, since the Clouds have a nearly polar orbit (Kallivayalil et al. 2013). It is the proper motions of these stars that now challenge the notion that these are LA members.

We plot the velocity components in Figure 10. Together with our target stars, we plot a set of Galactic disk O/B stars. These stars are from the DR5 catalog of the Large Sky Area Multi-Object Fiber Spectroscopic Telescope (LAMOST, also called the Guo Shou Jing Telescope) survey (Cui et al. 2012). Only stars with velocity errors $<50$ km s$^{-1}$ from this catalog (Liu et al. 2018) are shown. We use this sample of O/B stars as representative for the thin-/thick-disk kinematics. From Figure 10 and Table 6 it can be seen that eight stars have velocity components well within the range of thin-/thick-disk stars: low $V_R$ and $V_\phi$ with $V_z$ $\sim$ 120–250 km s$^{-1}$. These stars are CD14-A11, A13, A15, A18, B02, B04, B18, and C06. Three of these stars—specifically CD14-A15, B02 and B04—were previously assumed to be LA members according to their RVs. The remaining seven stars from our sample have peculiar velocities with respect to the canonical thin/thick disk, i.e., with large $V_R$ or $V_\phi$ or $V_z$. These high-velocity stars are highlighted in Figure 10 with blue and red symbol contours; the red contour is for the most energetic stars among these high-velocity stars (see the next subsection). We will next focus on the orbits of our 15 stars to further investigate their origins.

5.3. Orbit Integration and Parameters

Orbit integrations were done via galpy, a python package for dynamical calculations (Bovy 2015). In the calculations, the bulge, disk, and dark-halo contributions are modeled by a power law with exponential cutoff, a Miyamoto–Nagai disk (Miyamoto & Nagai 1975), and a Navarro–Frenk–White
Table 7 and plotted in Figure 11. Total specific orbital energy $E$ and $L_z$ orbital angular momentum are shown in Figure 12. $E$ and $L_z$ of the LMC under the same potential are also shown in this plot. The parameters used in the calculations of $E$ and $L_z$ of the LMC are: $D = 49.97 \pm 0.19$ kpc (Pietrzynski et al. 2013), $RV = 262.2 \pm 3.4$ km s$^{-1}$ (van der Marel et al. 2002), $\mu_L \cos \delta = 1.910 \pm 0.020$ mas yr$^{-1}$, and $\mu_\delta = 0.229 \pm 0.047$ mas yr$^{-1}$ (Kallivayalil et al. 2013).

The orbital parameters of our target stars indicate that a large number of them fit well within the thin-/thick-disk populations, in agreement with the velocity results. The seven high-velocity stars mentioned in the previous subsection are CD14-A05, A07, A12, A19, B03, B14, and B19. Four of these stars are also the most energetic ones in our sample, with large pericentric radii, ($R_{peri} > 12$ kpc), and high orbital energies. These are CD14-A05, A19, B03, and B14. We caution that CD14-A19 has large errors due to its large distance error; nevertheless, its orbital energy seems large even within the quoted error (see Figure 12). The remaining three high-velocity stars have pericentric radii less than $\sim 10$ kpc. From Figure 12, they appear to have less rotation ($L_z$) at a given energy when compared to thin-disk stars (the dense gray plume of objects in Figure 12). In this regard, they fit with the low-rotation tail of the thick disk. Overall, none of our stars have orbital parameters compatible with an association with the LMC (see Figure 12).

6. Origin of the Stars

6.1. Orbits

Although in the previous subsections we touched upon the origin of these stars as indicated from velocities and orbital elements, here we further explore the orbits of the stars in the context of their lifetimes. We separate our 15 stars in the three groups already hinted at in the previous subsections. The groups are: (1) Group 1 (G1): four stars with large velocity components and high orbital energies: CD14-A05, A19, B03, B14; (2) Group 2 (G2): three stars with large velocity components, but moderate orbital energies: CD14-A07, A12, B19; and (3) Group 3 (G3): the remaining eight target stars, with no kinematical peculiarities with respect to the Galactic thin-/thick-disk population.

The ages are adopted as the integration time. We show the orbits in Figure 13: the left panels show the $X$-$Y$ Galactic plane, and the right panels show the $R_{GC} - Z$. The top row shows stars in G1, with each star labeled; the middle row shows stars in G2, and the third row shows stars in G3. In Figure 14 we show the $R_{GC}$ (left panels) and the total velocity $V_{tot}$ (right panels) as a function of time. G3 orbits (bottom rows) indicate in-plane, more-or-less circular orbits with $Z$ excursions up to 2.5 kpc. These are typical thin-/thick-disk-like orbits. Their birthplace is not far from the Galactic plane, and within 5–12 kpc from the GC. G2 orbits are qualitatively very similar to G3 orbits; total velocity appears to be somewhat larger. Therefore, G2 stars indicate a disk origin as well. The most probable origin of G2 stars and perhaps some G3 stars is in the thin disk, and then scattered into more thick-disk-like orbits via dynamical interactions of binary stars in young associations, i.e., disk runaway type (e.g., Brown 2015; Irrgang et al. 2018).

We also show times of flight to reach the Galactic plane ($t_{flight}$, listed in Table 7) versus the ages of the stars in Figure 15. With the exception of one star, CD14-B14, $t_{flight}$ of
Table 7
Orbit Parameters

| ID            | $Z_{\text{max}}$ (kpc) | $e$     | $R_{\text{age}}$ (kpc) | $R_{\text{peri}}$ (kpc) | Age (Myr) | $t_{\text{flight}}$ (Myr)$^3$ |
|---------------|-------------------------|---------|-------------------------|-------------------------|-----------|-------------------------------|
| CD14-A05$^b$  | 4.05$^{+0.64}_{-0.2}$  | 0.22$^{+0.05}_{-0.07}$ | 27.39$^{+8.77}_{-7.35}$ | 17.40$^{+22.99}_{-19.39}$ | 83$^{±7}$  | 2$^{±12}$                     |
| CD14-A07      | 3.45$^{+0.63}_{-0.2}$  | 0.36$^{+0.04}_{-0.05}$ | 13.00$^{+4.23}_{-2.60}$ | 6.15$^{+1.13}_{-0.99}$   | 87$^{±8}$ | 15$^{±14}$                    |
| CD14-A11$^b$  | 3.34$^{+0.82}_{-0.2}$  | 0.22$^{+0.25}_{-0.15}$ | 13.35$^{+3.92}_{-2.39}$ | 8.32$^{+4.53}_{-2.04}$   | 65$^{±6}$ | 22$^{±16}$                    |
| CD14-A12$^b$  | 8.80$^{+0.27}_{-0.10}$ | 0.36$^{+0.06}_{-0.05}$ | 20.20$^{+3.97}_{-2.63}$ | 9.58$^{+0.63}_{-0.72}$   | 35$^{±5}$ | 21$^{±11}$                    |
| CD14-A13      | 2.14$^{+0.37}_{-0.2}$  | 0.07$^{+0.03}_{-0.05}$ | 10.05$^{+1.57}_{-0.98}$ | 8.70$^{+0.80}_{-0.48}$   | 69$^{±20}$ | 16$^{±16}$                    |
| CD14-A15$^b$  | 2.26$^{+0.29}_{-0.10}$ | 0.5$^{+0.14}_{-0.11}$   | 10.07$^{+0.78}_{-1.08}$ | 3.14$^{+0.91}_{-1.08}$   | 70$^{±20}$ | 13$^{±15}$                    |
| CD14-A18      | 0.72$^{+0.08}_{-0.02}$ | 0.15$^{+0.02}_{-0.01}$ | 6.68$^{+0.06}_{-0.61}$  | 4.94$^{+0.23}_{-0.28}$   | 10$^{±8}$ | 14$^{±10}$                    |
| CD14-A19$^c$  | 7.88$^{+3.59}_{-0.4}$  | 0.16$^{+0.01}_{-0.04}$ | 25.80$^{+5.26}_{-9.57}$ | 19.59$^{+13.80}_{-15.13}$ | 76$^{±34}$ | 26$^{±26}$                    |
| CD14-B02$^b$  | 2.16$^{+0.21}_{-0.08}$ | 0.24$^{+0.08}_{-0.04}$ | 13.54$^{+0.87}_{-0.77}$ | 8.22$^{+2.07}_{-1.66}$   | 75$^{±8}$ | 10$^{±8}$                     |
| CD14-B03$^b$  | 3.45$^{+0.15}_{-0.14}$ | 0.18$^{+0.01}_{-0.01}$ | 17.56$^{+1.69}_{-1.57}$ | 12.24$^{+0.90}_{-1.09}$ | 70$^{±10}$ | 9$^{±5}$                      |
| CD14-B04      | 2.10$^{+0.16}_{-0.2}$  | 0.27$^{+0.05}_{-0.06}$ | 9.90$^{+0.37}_{-0.30}$  | 5.71$^{+0.96}_{-0.81}$   | 6$^{±4}$  | 14$^{±9}$                     |
| CD14-B11$^b$  | 4.38$^{+0.42}_{-0.02}$ | 0.36$^{+0.02}_{-0.02}$ | 38.11$^{+5.72}_{-5.72}$ | 18.02$^{+2.10}_{-1.48}$ | 70$^{±10}$ | 99$^{±15}$                    |
| CD14-B18      | 1.63$^{+0.17}_{-0.13}$ | 0.29$^{+0.05}_{-0.07}$ | 9.47$^{+0.32}_{-0.43}$ | 5.18$^{+0.16}_{-0.59}$   | 5$^{±2}$  | 10$^{±3}$                     |
| CD14-B19      | 3.33$^{+0.28}_{-0.11}$ | 0.23$^{+0.05}_{-0.05}$ | 8.35$^{+0.26}_{-0.98}$  | 5.17$^{+0.84}_{-0.96}$   | 32$^{±6}$ | 18$^{±11}$                    |
| CD14-C06      | 1.84$^{+0.17}_{-0.11}$ | 0.28$^{+0.06}_{-0.06}$ | 13.45$^{+0.72}_{-0.72}$ | 7.54$^{+1.30}_{-1.24}$   | 59$^{±8}$ | 11$^{±14}$                    |

Notes.
$^a$ The time of flight to reach the Galactic plane.
$^b$ Sample stars from Z17.
$^c$ Sample star from CD14.

G1, G2, and G3 stars are shorter than or comparable to their ages within errors. This also implies a disk origin (Irrgang et al. 2018). Two stars in G3 have typical thin-disk velocities (and thus origin), with no need to invoke any dynamical interactions; these are CD14-A13 and CD14-A18.

G1 stars appear to have originated near the Galactic plane (except CD14-A19, which has large error bars). They were also born at $R_{\text{GC}} \sim 20$ kpc. The fact that they were born near the Galactic plane and have substantial rotation implies a disk origin as well. Thus, their origin appears unrelated to the gaseous material from the LA, which should follow the Clouds’ orbit about the Galaxy, or a nearby polar orbit. Another possible origin, still related to the Magellanic system, was proposed by Boubert et al. (2017). They proposed a runaway origin from the LMC for some hyper-velocity stars, and computed a model to determine the spatial distribution and kinematics of such stars. Our G1 stars do not fit this scenario either. This is because of the G1 stars’ large $V_{\phi}$, an aspect also reflected in their large $\mu_\phi$ proper motions (see our Table 5 values of the order of $\sim 4$ mas yr$^{-1}$) compared to the predictions ($\sim 2$ mas yr$^{-1}$) in Boubert et al. (2017), their Figure 3).

If G1 stars also have a disk runaway origin, then their birthplace is far out in the disk. Specifically, the plane crossing is between 15 and 30 kpc (see Figure 13). Is this unusual? Among the 15 B-type, hyper-velocity stars studied by Irrgang et al. (2018) in light of Guaita DR2 proper motions, only 2 have an unambiguous disk runaway origin at radii 15–20 kpc, namely HV57 and B485. Only one star in their sample, namely HV53, which is also the youngest at $\sim 18$ Myr, has an LMC runaway origin. The remaining stars in their sample are either runaways from near the solar circle, or have proper-motion uncertainties that are too large for a definite conclusion. The total velocities of our G1 stars are well within the range of disk runaway stars, albeit at the lower end: see our Figure 14 and Table 1 in Irrgang et al. (2018; see also Brown et al. 2018). Therefore, it does appear that our G1 stars have been born rather far away in the disk, but there is some overlap with the findings of Irrgang et al. (2018).

Besides their birthplace in the outer disk, another intriguing aspect of our G1 stars is that they all originate from a similar Galactic longitude $l \sim 350$ deg.

This may be a selection effect due to our choice of potential LA members in a given area on the sky and with RV $> 100$ km s$^{-1}$.

To conclude, the kinematical evidence points to a disk runaway origin for most of the stars in our sample; the most energetic ones were born in the outskirts of the Galactic disk at 12–28 kpc from the GC. Only one star, CD14-B14, has an age shorter than the time to reach the plane, indicating a birthplace some 2.5 kpc above the disk. It is hard to imagine Galactic OB associations exist at $\sim 2.5$ kpc above the plane, and at a distance of $\sim 28$ kpc from the GC. Could this star have an extragalactic origin? If so, from what system does it originate, since an association with the Magellanic system seems unlikely?

6.2. Abundances

We then investigate abundances as functions of angular momentum and energy for our target stars. As in our previous study, we choose three clusters to chemically represent the Galactic disk, the LMC, and the SMC. These clusters are NGC 4755, NCG 2004, and NCG 330, respectively, with abundances from Trundle et al. (2007); here we use the average abundance of Mg for B-type stars in these clusters.

We inspect He, C, and Mg because these are best determined for the entire set of stars. Figures 16 and 17 show the abundance ratios for He, C, and Mg as functions of angular momentum and energy, respectively. G1 stars are highlighted by surrounding red squares. There are only three stars, because CD14-A19 does not have high-resolution spectroscopic observations.
In the bottom panels of Figures 16 and 17 we also indicate the average $[\text{Mg}/\text{H}]$ abundance for our three representative samples; the Galactic disk (red dashed line), the LMC (blue dashed line), and the SMC (yellow dashed line). The color-coded shaded areas correspond to $1\sigma$ ranges around the averages.

At high energy and angular momentum Mg is low. Our three G1 stars have an average $[\text{Mg}/\text{H}] = -0.50 \pm 0.07$, which is compatible with LMC’s Mg abundance. The remaining stars have an average $[\text{Mg}/\text{H}] = 0.12 \pm 0.36$. They have a large scatter owing to the fact that they split into two groups at $[\text{Mg}/\text{H}] \sim 0.4$ and $[\text{Mg}/\text{H}] \sim -0.2$. Therefore, on average, our most energetic stars are formed from the least Mg-enriched material, which happens to be compatible with LMC’s abundance. However, this low Mg abundance can simply reflect the abundance gradient in the disk (see, e.g., McEvoy et al. 2017, and references therein).

C abundance shows a large scatter over the entire range of energy and angular momentum. He abundance, in contrast, is rather tight ($\sigma_{[\text{He}/\text{H}]} = 0.15$), with super-solar values for stars with low energy and angular momentum, and it becomes very scattered ($\sigma_{[\text{He}/\text{H}]} = 0.72$) at high energy and angular momentum. This scatter in He, and possibly C, may reflect different mechanisms that produce runaway stars. The two main mechanisms that produce disk runaway stars are: the binary supernova process when the runaway is the secondary star in a binary system where the most massive star undergoes a supernova explosion, and the dynamical ejection process when multibody interactions take place in a cluster/association. As discussed in McEvoy et al. (2017; see also Przybilla et al. 2008a), both these processes may affect the surface chemical composition of the runaway star. However, McEvoy et al. (2017) concluded that abundance
Figure 13. Orbits of the target stars in (a) the Galactic $X$–$Y$ plane and (b) in the $R_{GC}$–$Z$ plane. The filled star symbols and open circle symbols represent current positions and positions of the birthplace, respectively. The shades depict the error range of the orbits. The Sun’s location is represented with a dotted circle, while the LMC’s location is represented with a dotted square. In the upper panel of (a), the filled circles along the trajectories mark the crossing time of the Galactic plane. The hatched areas in panels (a) indicate the observed sky span of the LA (Nidever et al. 2010) projected onto the $X$–$Y$ plane. In panels (b) the hatched areas show the $R_{GC}$ range of the LA estimated in McClure-Griffiths et al. (2009).
Figure 14. Similar to Figure 13, but for variations of $R_{GC}$ and $V_{tot}$ with integration time. The hatched area in plot (a) indicates the $R_{GC}$ range of the LA estimated in McClure-Griffths et al. (2009).
patterns of C, N, Si, and Mg do not differ in B-type runaway stars from their counterparts in open clusters, instead they attribute the scatter of these elements to the different birthplaces of the runaway stars. Increased He abundance can also be due to mixing between surface and interior in these fast-rotating B-type stars. To summarize, the scatter in He (and possibly C) seen in the most energetic stars can be interpreted at best as a confirmation of their runaway status.

7. Summary

We have measured chemical abundances and RVs for nine B-type stars in the region of the Magellanic LA, based on high-resolution spectra taken with the MIKE instrument on the 6.5 m Clay telescope. Among the nine stars, two were found to be HB stars based on the stellar parameters, and were excluded from the analysis. We combined the remaining seven young stars with the samples in Z17 and CD14 for a total of 15 stars. For these 15 stars, 3D velocities (using Gaia DR2 proper motions) and orbits were determined. We found that none of the RV candidates to LA membership have orbits compatible with a Magellanic origin. This is owing to their large proper motions along Galactic longitude, as measured by Gaia DR2. To this end, the fact that these stars are not members of the LA indicates that the abundance of the LA is mainly SMC-like, as consistently determined from the gas studies (Lu et al. 1998; Fox et al. 2018; Richter et al. 2018).

The majority of the stars in our sample originate as: (1) typical thin-disk stars (and are chosen as such in our observing program for comparison purposes), and (2) runaway disks with approximately thick-disk-like kinematics. The ages of these stars (with one exception) are comparable to or larger than their times since disk crossing, thus confirming their runaway status. However, the most energetic star in our sample, CD14-B14, has an age (70 Myr) shorter than the time since disk crossing, indicating a birthplace some 2.5 kpc above the Galactic plane and at ~28 kpc from the...
GC. With an orbital energy comparable to that of the LMC, but an orbital angular momentum much larger, it is unclear what the origin of this star is. The next three most energetic stars have birthplaces ranging from 12 to 25 kpc from the GC. These four most energetic stars also appear to have been born at a similar Galactic longitude of ~350 deg; this may be an observational selection effect.

The three most energetic stars that have abundance measurements have low Mg abundance, with an average comparable to the Mg abundance in the LMC (as represented by cluster NGC 2004). This may also simply imply that they formed from less enriched material on the outskirts of the Galactic disk. The He abundance is tight and super-solar for low orbital-energy stars, while for high orbital-energy stars it is very scattered. This large He abundance scatter may reflect the different mechanisms that produce runaway stars.

We thank Dr. Z.-X. Lei and Dr. C. Liu for discussions of the stellar evolution, and the anonymous referee for helpful comments. This project was developed in part at the 2018 *Gaia*-LAMOST Sprint workshop, supported by National Science Foundation of China (NSFC) grants 11333003 and 11390372. L.Z. acknowledges supports from NSFC grants 11773033, 11573030, and 11673030. D.I.C. acknowledges partial support from NASA grant 80NSSC18K0422. C.M.B. acknowledges support from NASA grant 80NSSC18K0422. V.K. acknowledges support from RSF grant No. 18-12-00213.

References

Bressan, A., Marigo, P., Girardi, L., Nanni, A., & Rubele, S. 2013, in EJP Web

Bressan, A., Moni Bidin, C., Girardi, T. M., et al. 2014, *ApJ*, 784, L37

Casetti-Dinescu, D. I., Moni Bidin, C., Girardi, T. M., & van Altena, W. F. 2012, *ApJ*, 753, 123

Chen, Y., Girardi, L., Bressan, A., et al. 2014, *MNRAS*, 444, 2525

Cui, X.-Q., Zhao, Y.-H., Chu, Y.-Q., et al. 2012, *RAA*, 12, 1197

Dehen, W., & Binney, J. J. 1998, *MNRAS*, 298, 387

D’Onghia, E., & Fox, A. J. 2016, *ARA&A*, 54, 563

For, B.-Q., Staveley-Smith, L., & McClure-Griffiths, N. M. 2013, *ApJ*, 764, 74

Fox, A. J., Barger, K. A., Wakker, B. P., et al. 2018, *ApJ*, 854, 142

Fox, A. J., Wakker, B. P., Barger, K. A., et al. 2014, *ApJ*, 787, 147

Gaia Collaboration, Brown, A. G. A., Vallenari, A., et al. 2018, *A&A*, 616, A1

Gaia Collaboration, Prusti, T., de Bruijne, J. H. J., et al. 2016, *A&A*, 595, A1

Grevesse, N., & Sauval, A. J. 1998, *SSRv*., 85, 161

Hammer, F., Yang, Y. B., Flores, H., Puech, M., & Fouquet, S. 2015, *ApJ*, 813, 110

Irrgang, A., Kreuzer, S., & Heber, U. 2018, *A&A*, 620, A48

Kalifivayalil, N., van der Marel, R. P., Besla, G., Anderson, J., & Alcock, C. 2013, *ApJ*, 764, 161

Kelson, D. D. 2003, *PASP*, 115, 688

Kelson, D. D., Illingworth, G. D., van Dokkum, P. G., & Franx, M. 2000, *ApJ*, 531, 137

Kounouhoven, M. B. N., Brown, A. G. A., Portegies Zwart, S. F., & Kaper, L. 2007, *A&A*, 474, 77

Kurucz, R. 1993, *ATLAS9 Stellar Atmosphere Programs and 2 km/s Grid* (Cambridge, MA: Smithsonian Astrophysical Observatory), 13, Kurucz CD-ROM No. 13

Kurucz, R., & Bell, B. 1995, in Atomic Line Data. ed. R. L. Kurucz & B. Bell (Cambridge, Mass.: Smithsonian Astrophysical Observatory), 23, Kurucz CD-ROM No. 23

Lanz, T., & Hubeny, I. 2003, in ASP Conf. Ser. 288, Stellar Atmosphere Modeling, ed. I. Hubeny, D. Mihalas, & K. Werner (San Francisco: CA: ASP), 117

Lanz, T., & Hubeny, I. 2007, *ApJS*, 169, 83

Liu, C., Cui, W., Liu, C., et al. 2018, ApJ, submitted

Lu, L., Sargent, W. L. W., Savage, B. D., et al. 1998, *AJ*, 115, 162

Mackey, D., Koposov, S., Da Costa, G., et al. 2018, *ApJ*, 858, L21

Magrini, L., Sestito, P., Randich, S., & Galli, D. 2009, *A&A*, 494, 95

McClure-Griffiths, N. M., Pisano, D. J., Calabretta, M. R., et al. 2009, *ApJS*, 181, 398

McClure-Griffiths, N. M., Staveley-Smith, L., Lockman, F. J., et al. 2008, *ApJ*, 673, L143

McEvoy, C. M., Dufton, P. L., Smoker, J. V., et al. 2017, *ApJ*, 842, 32

Miyamoto, M., & Nagai, R. 1975, *PASJ*, 27, 533

Moebeil, S., Sweigart, A. V., Landsman, W. B., Heber, U., & Catelan, M. 1999, *A&A*, 346, L1

Moni Bidin, C., Moehler, S., Piotto, G., Momany, Y., & Recio-Blanco, A. 2009, *A&A*, 498, 737

Moni Bidin, C., Villanova, S., Piotto, G., et al. 2012, *A&A*, 547, A109

Munari, U., Sordo, R., Castelli, F., & Zwitter, T. 2005, *A&A*, 442, 1127

Navarro, J. F., Frenk, C. S., & White, S. D. M. 1997, *ApJ*, 490, 493

Nidever, D. L., Dey, A., Olsen, K., et al. 2018, *AJ*, 156, 134

Nidever, D. L., Majewski, S. R., & Butler Burton, W. 2008, *ApJ*, 679, 432

Nidever, D. L., Majewski, S. R., Butler Burton, W., & Ngira, L. 2010, *ApJ*, 723, 1618

Pardy, S. A., D’Onghia, E., & Fox, A. J. 2018, *ApJ*, 857, 101

Pietrzyński, G., Grażycki, D., Gieren, W., et al. 2013, *Natur*, 495, 76

Przybilla, N., Fernanda Nieva, M., Heber, U., & Butler, K. 2008a, *ApJL*, 684, L103

Przybilla, N., Nieva, M. F., Heber, U., et al. 2008b, *A&A*, 480, L37

Putman, M. E., Gibson, B. K., Staveley-Smith, L., et al. 1998, *Natur*, 394, 752

Reid, M. J., Menten, K. M., Brunthaler, A., et al. 2014, *ApJ*, 783, 130

Richter, P., Fox, A. J., Wakker, B. P., et al. 2018, *ApJ*, 865, 145

Sales Silva, J. V., Carraro, G., Anthony-Twarog, B. J., et al. 2016, *AJ*, 151, 6

Schunker, E., & Finkbeiner, D. P. 2011, *ApJ*, 737, 103

Schlegel, D. J., Finkbeiner, D. P., & Davis, M. 1998, *ApJ*, 500, 525

Schönrich, R. 2012, *MNRAS*, 427, 274

Tian, H.-J., Liu, C., Carlin, J. L., et al. 2015, *ApJ*, 809, 145

Torry, J., & Davis, M. 1979, *AJ*, 84, 1511

Trundle, C., Dufton, P. L., Hunter, L. et al. 2007, *A&A*, 471, 625

van der Marel, R. P., Alves, D. R., Hardy, E., & Suntzeff, N. B. 2002, *AJ*, 124, 2639

Vennin, M. S., Kerp, J., & Kalberla, P. M. W. 2012, *A&A*, 547, A12

Zhang, L., Moni Bidin, C., Casetti-Dinescu, D. I., et al. 2017, *ApJ*, 835, 285

Zhang, et al.