Phase diagram, stability and electronic properties of an Fe–P system under high pressure: a first principles study†

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Fe–P binary compounds have attracted much attention, particularly under high pressure, since they are the constituents of the Earth’s core. However, most studies focus on the single stoichiometry of Fe–P binary compounds at high pressure, and their whole phase diagram and relative stabilities have been unexplored thus far. Herein, first principles swarm structure predictions are performed to find stable structures of Fe–P compounds with various Fe₇–₈P₆ (x = 1–4 and y = 1, or x = 1 and y = 2) compositions. Then, their phase diagram and relative stabilities are reliably determined based on predicted structures. Specifically, the FeP, Fe₂P and Fe₄P compounds are found to be stable in the pressure range of 0–400 GPa. The Fe₆P compound decomposes into Fe₅P and Fe₄P above 214 GPa. Fe₇P becomes unstable above 82 GPa. Notably, two new phases (i.e. C2/c-structured Fe₆P and Cmcm-structured Fe₇P) are found to be more stable than the previously reported phases. In addition, the XRD pattern of the predicted Cmcm-structured Fe₆P matches the experimental patterns, and we are awaiting future experimental confirmation. Electronic band calculations show that the Fe–P binary compounds are metallic, with a pronounced Fe 3d component crossing the Fermi level. Cmcm-structured Fe₆P is ferromagnetic. Our study not only provides useful information for the further study of Fe–P binary compounds but also for the determination of the Earth’s core components.

1. Introduction

Geophysical and geophysical studies show that light elements (i.e. hydrogen, oxygen, carbon, sulfur, phosphorus and silicon) play important roles in the Earth’s core.1–2 Early models of the Earth’s core consisted of a considerable amount of schreibersite, cohenite, and troilite.3 The Fe content of the Earth’s core is rather rich (90.67%).4 Several phases of iron compounds, such as iron–sulfur, iron–silicon, and iron–oxygen, with light elements under high pressure have been determined by a large number of experimental and theoretical studies.4–7 Notably, the phosphorus content of the Earth’s core was estimated to be around 0.20 wt%, and up to 4 wt% of phosphorus was found to be soluble in crystalline iron at 23 GPa in the Fe–S–P system.8,9 Since the structures of Fe–P compounds are similar to that of Fe–S, iron phosphides should have a stable composition in the cores. As a consequence, investigation of the structures and properties of iron-phosphorus (Fe–P) compounds is worthy of attention.

To date, extensive studies have been carried out on Fe–P binary compounds, and although most of these studies were based on single stoichiometry, the relative stabilities of the Fe–P binary compounds remain undetermined, particularly at high pressure. For example, several stoichiometric compounds (i.e. FeP, Fe₂P, Fe₃P, Fe₄P, FeP₂ and FeP₃) were separately investigated. FeP is known as mellite, which is a mineral from Northwest Africa 1054 aca-pulcoite. The cubic polymorph with the P2₁3 space group symmetry (AlAu₄-type structure) was determined using single-crystal X-ray diffraction data (XRD).10 Under high pressure, the FeP phase is relatively stable compared to FeP + Fe.11 The tetragonal structure Fe₃P was discovered from schreibersite at ambient pressure.12 Its possible high-pressure polymorph is either a P4/mnc or Pnam structure, as predicted by theoretical calculations, and the pressure-induced phase transition from I to P4/mnc has been observed at 64 GPa.13 For the Fe₄P compound, it is a terrestrial garnet peridotite with implications for the Earth’s core. A hexagonal structure of Fe₃P with P6₃m2m symmetry exists at ambient conditions, and the Co₃Si-type structure (space group Pnam) is found to be more stable above 8.0 GPa (ref. 14) or 16.8 GPa,15 as observed through high-pressure and high-temperature X-ray diffraction experiments. Based on first principles calculations,
another P3m-structured Fe2P was reported at 153 GPa.16 Several studies on FeP have been carried out because FeP is also an important member in the planetary cores. FeP crystallizes in the regular MnP-type structure (space group Pnma) at ambient conditions and undergoes a structural phase transition to a simple cubic FeSi-type phase (space group P213) at 87.5 GPa.20,21 FeP2 crystallizes in the regular marcasite structure (space group Pnmm) at ambient conditions with no unpaired spins in the low-spin configuration.19 Under high-pressure and high-temperature conditions, a new phase of FeP2 with Immc symmetry was observed at 70 GPa.22 Recent studies show that Fe–P binary compounds (i.e. FeP, Fe2P, FeP, FeP2 and FeP3) are stable at ambient pressure,23 in which all the structures came from reported compounds. By all appearance, it is necessary to establish the whole phase diagram of Fe–P binary compounds under high pressure.

Recently, the first principles structural prediction method has become an irreplaceable tool to discover new materials and determine their stable structures.23–24 For example, a recent breakthrough in the superconductivity field is the observation of remarkably high superconductivity (with $T_c$ up to 203 K),25 in which H3S is the main stoichiometry of the superconducting phase.26 This observation was achieved by direct investigation on the theoretical prediction of high-$T_c$ superconductivity in compressed solid H2S.27 A new clathrate sodalite-like BN structure was predicted with a high energy density and large hardness.28 The reaction of iron with xenon has been predicted in the Earth’s inner core environment through theoretical calculations, which provides a solution to the missing Xe paradox.29 Numerous unusual stoichiometric compounds with exotic physical or chemical properties have been discovered under high pressure.30–37

Herein, the most stable structures of the Fe–P binary compounds (Fe3P, Fe2P, FeP, FeP and FeP2) are reliably determined using the first principles calculations in combination with the swarm structural search technology.21–26 The relative stabilities of these considered compounds are reliably established. Two new phases (C2/c-structured Fe2P, and Cmcm-structured Fe3P) are found to be more stable than the previously reported phases. Cmcm-structured Fe3P meets the Stoner criterion and is ferromagnetic. Our study provides some useful information for the investigation of the constituents of the Earth’s core.

2. Computational details

To search the stable structures for the given chemical compositions, structural predictions of the Fe–P system were performed at 0, 50, 100, 200, 300, and 400 GPa using the Crystal structure AnaLYsis by the Particle Swarm Optimization (CALYPSO) code.23–29 The CALYPSO method can successfully find stable structures just depending on chemical compositions and has demonstrated success in predicting the high-pressure structures of various systems.40–44 The ab initio structural relaxations and electronic calculations are based on the density functional theory (DFT) within the Perdew–Burke–Ernzerhof (PBE) of the generalized gradient approximation (GGA),45 as implemented in the Vienna ab initio simulation package (VASP) code.46 The electron–ion interaction was described using the all-electronic projector-augmented wave (PAW) method47 with $3s^23p^63d^14s^1$ and $3s^23p^6$ valence electrons for the Fe and P atoms, respectively. Hubbard $U$ is necessary to correct the effective on-site electron–electron interaction of the transition metal, with valence electrons on the d and f orbits.48,49 Hubbard $U$ parameters with $U = 2$ eV and $J = 0.95$ eV were applied to the Fe atoms in the Fe2P compounds.46 It successfully reproduces the experimental magnetic moments of Fe atoms.46 As a consequence, we chose Hubbard $U$ with 2 eV at a fixed value of 0.95 eV (i.e. $U_{eff} = U - J$) to calculate the properties and fine structure optimization. $U = 2.2$ eV was verified to be suitable to elemental iron by the linear response approach.50 The cutoff energy of 700 eV and Monkhorst–Pack27 k-meshes with a grid spacing of $2\pi \times 0.03$ Å$^{-1}$ were used to ensure that the total energy calculations converged well. The phonon calculations were carried out using the finite displacement approach to determine the dynamical stability of the predicted structures, as implemented in the Phonon code.51 To study the interatomic interaction, the integrated crystal orbital Hamilton populations (ICOHPs) were calculated as implemented in the LOBSTER package.52

3. Results and discussion

3.1. Phase stabilities of Fe–P compounds

Herein, we mainly considered Fe-rich phosphorus compounds (i.e. Fe3P, Fe2P, FeP, FeP and FeP3) since the Fe content of the Earth’s core is estimated to be around 90.67% and P content is about 0.17%.4 In this study, the pressure ranges up to 400 GPa are considered. The predicted structures of Fe–P binary compounds were used to evaluate the formation enthalpy ($\Delta H$) relative to the elemental solids Fe and P according to eqn (1), in order to explore the thermodynamic stabilities of Fe–P compounds.

$$\Delta H(\text{Fe}_m\text{P}_n) = [H(\text{Fe}_m\text{P}_n) - mH(\text{Fe}) - nH(P)]/(m + n).$$

(1)

where $H = U + PV$ is the enthalpy of each composition and $\Delta H$ is the enthalpy of formation per atom. Herein, $U$, $P$, and $V$ are the internal energy, pressure, and volume, respectively. The relative stabilities of the considered compositions at the selected pressures of 0, 50, 100, 200, 300, and 400 GPa, with $\Delta H$ evaluated per atom, are shown in Fig. 1a. The stable phase lies on the convex hulls (the global stability lines). As can be seen, all the considered compositions are stable at 0 GPa, which is well in agreement with the recent theoretical study.41 It should be noted that the Fe2P compound is not included in the study by Jing Wu et al.21 $P2_13$-structured FeP2 has been found from mellinite and is confirmed to be stable at ambient pressure.49 The known FeP, FeP, Fe3P, FeP, FeP, and FeP3 structures40,47,48,55,56 at ambient pressure were well reproduced in our calculations, which validates our structure searching method applied to the Fe–P system. In addition, our calculated lattice parameters for Fe3P (space group I4, 8 formula units per cell) are $a = b = 9.13$ Å and $c = 4.40$ Å, which are in excellent agreement with the experimental values at $a = b = 9.10$ Å and $c = 4.46$ Å.47 These lattice parameter calculations support the validity of the pseudo potentials adopted in this study.
Two new phases (i.e. C2/c-structured Fe₄P and Cmcm-structured Fe₃P) were discovered by comparing previous structures and our predicted structures, and their structural characters are shown in Fig. 2. The predicted C2/c-structured Fe₄P (4 formula units) is more stable than the phase with P2₁/3 symmetry above 37 GPa (Fig. S3a†). Its lattice parameters are \( a = 7.342 \, \text{Å}, \quad b = 3.772 \, \text{Å}, \quad c = 5.627 \, \text{Å} \) at 200 GPa. The structure contains two inequivalent Fe atoms occupying 8f (0.297, 0.337, 0.850) and 8f (0.900, 0.671, 0.551) sites and one equivalent P atom sitting at the 4e (0.000, 0.170, 0.750) position (Fig. 2a). To distinguish the two Fe atoms at different sites, we labelled the inequivalent Fe atoms with two different colors (Fig. 2c). In this structure, each P atom is in 12-fold coordination with the Fe atoms (Fig. 2b), forming a tetrakaidecahedron. The adjacent polyhedrons are interconnected through common edges or co-vertices. The distance of the two nearest Fe–Fe is 2.11 Å, which is smaller than the Fe–Fe distance of 2.48 Å in bulk Fe with Im3m symmetry under ambient conditions and 2.44 Å in Fe₅S₄₁ at 100 GPa.²⁵ However, this structure is nonmagnetic, as will be discussed later. To determine the interatomic interaction in the C2/c-structure Fe₄P, the integrated crystal orbital Hamilton populations (ICOHPs) for the Fe–P and Fe–Fe.

Pairs at 200 GPa were calculated and are shown in Table S1.† The values of ICOHP between the Fe–P and Fe–Fe are –1.80 and –4.97 eV, respectively. The ICOHP value of the Fe–Fe pair is very close to –4.90 eV in Fe₅S₄₁.²⁵ Apparently, the Fe–Fe interaction plays an important role in determining the structural stability of the compound. Further analysis indicates that the major contribution of the Fe–Fe interaction originates from 4s–3d, which corresponds to the decomposed ICOHP of –1.908 eV per pair.

The Fe₃P compound with the I₄ structure symmetry is stable at ambient conditions, and the first phase transition is observed at 64 GPa.²³ This phase is assigned as the P4/mnc or Pnma structure. Notably, we discovered an orthorhombic structure of Fe₃P (space group Cmcm, 4 formula units per cell) above 58 GPa. This structure is much more stable than the P4/mnc or Pnma structure in the pressure range of 0–400 GPa (Fig. S3b†). The lattice parameters of the Cmcm-structure Fe₃P are \( a = 5.446 \, \text{Å}, \quad b = 6.569 \, \text{Å}, \quad c = 3.255 \, \text{Å} \). This structure contains two inequivalent Fe atoms occupying the 4c (0.500, 0.900, 0.750) and 8g (0.784, 0.117, 0.750) positions, and one equivalent P occupying the 4c (0.000, 0.170, 0.250) site. For each P atom, there are 10

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**Fig. 2** Crystal structures of the Fe₄P and Fe₃P compounds at 200 GPa. (a) Fe₄P in the C2/c structure. (b) Coordination environment of the P atom in Fe₄P. (c) Coordination environment of the Fe atoms in Fe₄P. (d) Fe₃P in the Cmcm structure. (e) Coordination environment of the P atom in Fe₃P. (f) Coordination environment of the Fe atom in Fe₃P. Here, the large shallow blue or dark blue and small orange spheres represent Fe and P atoms, respectively.
nearest neighbours of Fe atoms, forming a P–Fe hendecahedron (Fig. 2e). The adjacent polyhedrons have common edges. The two inequivalent Fe atoms are labelled with two different colors (Fig. 2f). The nearest Fe–Fe distance in the Cmcm phase is 2.10 Å, which is comparable to 2.14 Å in Fe₄P. Unlike Fe₄P with C₂/c symmetry, the Cmcm-structured Fe₃P structure is magnetic. An experimental study investigated the phase transitions of Fe₃P, in which the P₄/mnc or Pnma structure is stable above 64 GPa. However, our predicted Cmcm phase is more stable than the P₄/mnc or Pnma phase at 64 GPa. The specific ΔH values of the three Fe₃P phases can be found in Table S2.† To confirm the existence of the Cmcm phase, the simulated XRD pattern of the Cmcm phase is shown in Fig. 3 along with the experimental patterns. Several possible structures of Fe₃P are also included in Fig. 3. Clearly, the main peak of Cmcm-structured Fe₃P is close to the experimental peaks or that of the other phases. Thus, Cmcm-structured Fe₃P might disappear under these experimental conditions.

3.3. Electronic properties of stable phases

The electronic band structures and projected density of states (PDOS) of the Fe₄P and Fe₃P compounds are shown in Fig. 4. Similar to the electronic properties of the other Fe–P species,¹⁸ there are several bands crossing the Fermi level, which indicate that they are metallic. As can be seen in Fig. 4, there is a common character of a pronounced Fe 3d component near the Fermi level. Interestingly, the Cmcm-structured Fe₃P is a ferromagnetic phase with the magnetic moment of 3.31μₑ per cell at 200 GPa by our calculations. Its electronic band structures and spin-polarized DOS based on spin up and spin down are described by the black and green lines, respectively.

4. Conclusion

To find the stable phases of Fe–P binary compounds, we explored the crystal structures and stabilities of the Fe–P system with FeₓPᵧ (x = 1–4 and y = 1, or x = 1 and y = 2) compositions under high pressure using first principles calculations combined with swarm structural searches. It is found that the FeP, Fe₂P and Fe₄P compounds are stable in the whole pressure range of 0–400 GPa, and Fe₃P becomes unstable above 214 GPa. Several phases of iron-rich compounds can be synthesized at elevated conditions, which may be the possible phases in the Earth’s cores. The FeP and Fe₂P compounds are metallic, with a pronounced Fe 3d component across the Fermi level. Our study provides the theoretical basis for the further investigation of Fe–P binary compounds. It is also important to explore the constituents of the Earth’s core.

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