We studied thermo-optical property degradation of indium tin oxide (ITO)-coated aluminized polyimide thin films under exposure to vacuum ultraviolet radiation and low-energy (3 and 5 keV) protons during ground tests using the Complex Irradiation Facility at the DLR site in Bremen. Changes in solar absorption and thermal emission coefficients caused by the irradiation were analyzed. We report a significant increase in solar absorptance of the samples irradiated by protons. We also attempted to identify any defects on the surface of the samples. The study was motivated by a unique opportunity that is provided by the Complex Irradiation Facility to study the degradation effects induced by exposure to protons with an energy below 10 keV and short-wavelength light below 115 nm.

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1. INTRODUCTION

MATERIALS certified for use in space are characterized by their exceptional properties (for instance lightweight, resistance to ionizing radiation, multifunctional capabilities, self-healing capabilities and outstanding thermal stability) so that they can survive in an environment that combines, among others, ionizing radiation, extreme temperatures, micrometeorites, and deep vacuum. Many space applications require a coating that is applied to the surface of the material to protect it or change its properties. Both the materials and their coatings used for spacecraft applications must be easy to apply and should have low outgassing and stability in space environment.

However, despite the unique characteristics, space is a harsh environment for the materials used on spacecraft, especially on their exterior surfaces. Most of these materials exhibit some degree of degradation due to different external factors, including ultraviolet and particle damage. One of the crucial aspects in spacecraft design is the thermal control system whose function is to keep the temperature of spacecraft systems within their operating ranges. The absolute temperature of a given area of the spacecraft in interplanetary space far away from Earth’s atmosphere is directly determined by the ratio between solar absorptance $a_s$ and thermal emittance $e_t$, whereby $a_s$ is defined as the fraction of incident solar radiant flux that is absorbed by a surface and $e_t$ as the ratio of the radiant intensity of the surface to that which would be emitted from a black body radiator. Thus, any change in thermo-optical properties of the surface materials affect the operation of the thermal control system and—directly or indirectly—the space mission.

Thermal control of spacecraft is mainly achieved by isolating it from the harsh environment by means of the multi-layer insulation (MLI) that prevents the heat from the Sun entering the spacecraft and restricts the loss of internally generated heat from the spacecraft. Typically, MLI consists of many (from several to few tens) layers of usually lightweight low-emittance polyimide films that are aluminized on one or both sides, separated by low-conductance spacers that reduce the heat transfer between the aluminized polyimide that acts as a radiation shield. Passive thermal control is also achieved by the application of various coatings with known optical properties, for example, high emittance (to improve the heat radiation characteristics) or low emittance (to minimize the radiative coupling).

An example of materials widely used in space applications is Kapton, a polyimide film developed by DuPont™, that is stable across a wide range of temperatures and used, among others, in flexible electronics and thermal blankets on spacecraft, satellites, and space instruments (see Reference 4 and references therein). When required, materials for spacecraft applications must be coated. One of the examples of a special purpose coating is indium tin oxide (ITO) that is often
applied to the front surface of the second surface mirror (the second surface mirror consists of an optically smooth metallic surface and a protective layer so that the incident rays do not hit the high reflectance layer directly but pass through the protective layer before and after reflection). The ITO coating is used to drain static electricity that is induced on the outer surface whose charges can build potentials of 20,000 to 30,000 volts. The conductive coating is transparent, thus it does not significantly affect the emittance or solar absorptance of the second surface mirror.

Motivated by a very limited number of studies on low-energy proton radiation effects (see References 5, 6 and references therein) and vacuum ultraviolet (VUV) radiation in space materials, we used the Complex Irradiation Facility (CIF) that—to our knowledge—offers a unique system with sources of very low-energy protons (down to below 10 keV) and short-wavelength light (down to 40 nm) to study the influence of irradiation with the VUV light source and the bombardment of the probe with protons (for more details about the facility see Section II–A) on thermo-optical properties of the ITO-coated aluminized polyimide foils. The main goal of these studies was extrapolation of real processes following material exposure to space conditions. In this case, we extrapolated from degradation effects caused by irradiation of the probes in the CIF. It was achieved by measurements of changes in thermo-optical properties, $\varepsilon_s$ and $\varepsilon_t$, of the irradiated material after the exposure when compared to its state before the experiment. We also attempted to identify any defects induced by the irradiation.

### Table I. The Samples Parameters Before and After Exposure to VUV Light or Protons

| Sample  | $\varepsilon_t$ | $\varepsilon_s$ |
|---------|----------------|----------------|
| ITO#1   | 0.64           | 0.44           |
| ITO#1   | 0.74*          | 0.38           |

| Sample  | $t_{lab}$ [h] | $t_{space}$ [d] | Fluence [p$^+$/cm$^2$] | $\Delta t_s$ |
|---------|----------------|-----------------|------------------------|--------------|
| ITO#2   | 13.0           | 20              | 9.10 x 10$^{16}$       | 0.38         |
| ITO#3   | 19.5           | 30              | 2.18 x 10$^{17}$       | 0.49         |
| ITO#4   | 32.5           | 50              | 2.93 x 10$^{17}$       | 0.47         |

| Sample  | $t_{lab}$ [h] | $t_{space}$ [d] | Fluence [p$^+$/cm$^2$] | $\Delta t_s$ |
|---------|----------------|-----------------|------------------------|--------------|
| ITO#14  | 11.0           | 407.0           | 9.10 x 10$^{16}$       | 0.45         |
| ITO#13  | 26.8           | 981.0           | 2.18 x 10$^{17}$       | 0.49         |
| ITO#9   | 47.2           | 1393            | 2.93 x 10$^{17}$       | 0.47         |
| ITO#12  | 93.1           | 3312            | 7.79 x 10$^{17}$       | 0.50         |
| ITO#10  | 221            | 7929            | 1.72 x 10$^{18}$       | 0.52         |

| Sample  | $t_{lab}$ [h] | $t_{space}$ [d] | Fluence [p$^+$/cm$^2$] | $\Delta t_s$ |
|---------|----------------|-----------------|------------------------|--------------|
| ITO#8   | 5.96           | 394.0           | 9.16 x 10$^{16}$       | 0.48         |
| ITO#11  | 17.5           | 944.0           | 2.22 x 10$^{17}$       | 0.55         |
| ITO#7   | 22.0           | 1267            | 3.14 x 10$^{17}$       | 0.53         |
| ITO#6   | 62.8           | 3370            | 7.46 x 10$^{17}$       | 0.55         |
| ITO#5   | 137            | 7445            | 1.79 x 10$^{18}$       | 0.56         |

$t_{lab}$—the irradiation time in the laboratory.
$t_{space}$—the corresponding irradiation time in space.
$\Delta t_s < 1.5$ pct (after irradiation) for all samples.

In these studies, the Sheldahl brand material of ITO-coated aluminized polyimide films was investigated (see Figure 1). The structure of the films is as follows: a layer of an ITO coating was applied to the front side and a reflecting aluminum coating to the back side of a polyimide substrate (for more details visit the Red Book webpage.*). These films are often used as an outer layer in MLI blankets.

The polyimide layer is 1.0 mil (25.4 µm) thick and the nominal aluminum coating thickness is 1000 Å (0.1 µm). The ITO coating has a standard target surface resistance of 5000 $\Omega$/sq. According to the manufacturer, it corresponds to a thickness of less than 100 Å (0.01 µm); thus, it should have no measurable impact.

II. TEST DESCRIPTION

In these studies, the Sheldahl brand material of ITO-coated aluminized polyimide films was investigated (see Figure 1). The structure of the films is as follows: a layer of an ITO coating was applied to the front side and a reflecting aluminum coating to the back side of a polyimide substrate (for more details visit the Red Book webpage.*) These films are often used as an outer layer in MLI blankets.

Fig. 1—A picture of a foil that was used in our experiments and its specification.

*https://www.sheldahl.com/technology/thermal-control-materials/the-red-book.
on emittance and increases the solar absorptance by about 5 pct. The product was perforated prior to metalizing, which is a standard procedure to enhance its functionality. The samples were exposed to low-energy protons and VUV radiation in three sets of experiments using the CIF at a temperature of 300 K.

A. Complex Irradiation Facility

1. General description

The Complex Irradiation Facility at the DLR Institute of Space Systems in Bremen (Germany), presented in Figure 2, is an instrument designed to study material surface property changes after exposure to corpuscular and electromagnetic radiation. Such studies are mandatory e.g., for a proper thermal design of satellite components exposed directly to radiation in a space environment. The facility is an Ultra High Vacuum system, equipped with two linear particle accelerators and three light sources. Both electrons and protons can be generated in an energy range from 2 to 100 keV. The maximum current of the particles can reach 100 \( \mu \text{A} \) over a maximum area of 60 \( \times \) 60 mm. All radiation sources can expose test specimens simultaneously, which presents a unique opportunity to study synergistic degradation effects on the materials under study. More importantly, each of the CIF sources can be used with an acceleration factor higher than one, which allows us to study how radiation with an intensity much higher than the equivalent of the Sun intensity influences degradation mechanisms.

2. The particles irradiation system

The corpuscular radiation of the CIF is achieved by use of two independent linear electron and proton accelerators. Both can irradiate specimens separately or simultaneously. The particle current can be set from 1 nA to 100 \( \mu \text{A} \). Only one energy level (between 2 and 100 keV) of the particles can be chosen; however, each particle source can irradiate the sample with different energy. Electrons are produced by heating up a LaB\(_6\) cathode. By changing the heater current and the Wehnelt voltage, one can control the electron current. Protons are produced by ionization of Hydrogen. \( \text{H}_2 \) gas is pumped into a glass bulb where the ionization is...
achieved by radio frequency excitation. Protons, H\textsuperscript{2+}, and H\textsuperscript{3+} ions are produced, which makes a mass selection necessary. It is accomplished by using a magnetic field to bend the particles beam. A careful selection of the field’s strength enables to beam only the protons into the irradiation chamber, while the other ions strike the beamline wall.

The cross section of both beams has a circular shape and its diameter depends on the incident energy of the particles. The maximum diameter is approx. 4 cm, while the minimum can be achieved by use of a retractable aperture plate with a drilled hole of approx. 1 mm in diameter, located in front of the irradiation chamber. The sample can be exposed to a maximum effective irradiation area of 6 × 6 cm. It is achieved by use of a beam sweep mechanism, which is accomplished by two pairs of plates, each with an opposite potential. Working together, they sweep the beam across the whole sample area.

3. The light sources
The CIF has three light sources: a Xenon lamp, a Deuterium lamp, and an Argon-VUV-source. All of the light sources can work simultaneously and cover a wide wavelength range from 40 to 2150 nm. The Xenon lamp is calibrated on demand before each test, while the Deuterium lamp and the Argon source were calibrated at the Physikalisch-Technische Bundesanstalt (PTB) in Berlin. The Argon-VUV source that we used during our tests produces light by exciting an Argon (98.5 pct), Krypton (1 pct), and Helium (0.5 pct) gas mixture with a well-defined beam of 1 keV electrons. The electrons pass through the gas flow located in the center of a vacuum chamber and between two nozzles. One pushes the gas with a defined flow rate, the second pulls the gas out of the chamber. The rest of the gas condensates on a cold baffle which surrounds the two nozzles. The gas atoms excited by the electrons drop their energy to the ground state and emit photons. Since the VUV source is continuously supplied with gas, which is then used to produce the light, the light intensity of the source remains relatively stable over time. The calibration campaign, the light spectrum for three different gas flows, together with identified emission lines, and the analysis of the source stability can be found in Reference [8]. The source produces VUV light with an intensity and wavelength range exceeding those produced by commonly used Deuterium lamps. These lamps have a wavenumber cut-off at around 115 nm which is related to internally used MgF\textsubscript{2} or LiF\textsubscript{2} windows. It means that, wavenumber cut-off at around 115 nm which is related to the commonly used Deuterium lamps. These lamps have a wavelength range exceeding those produced by xenon lamps. When this pressure is reached, the turbo molecular pump is activated to lower the vacuum level to 10\textsuperscript{-8} mbar. The further step is achieved with an ion getter pump.

4. The sample holder
Our standard sample holder is a frame consisting of two parts. The specimen is placed between the two and held by means of four screws. The maximum effective irradiation area is 60 × 60 mm. The frame is made up of stainless steel.

5. Vacuum control
Vacuum level in the empty irradiation chamber reaches a low 10\textsuperscript{-9} mbar range. The vacuum is achieved using a differential pumping system. After venting, the chamber is evacuated by the use of a scroll pump to a level of 10\textsuperscript{-3} mbar. When this pressure is reached, the turbo molecular pump is activated to lower the vacuum level to 10\textsuperscript{-8} mbar. The further step is achieved with an ion getter pump.

6. Temperature control
Temperature control is performed by means of a thermostat. Temperature is measured with two PT100 sensors. The first is located at the copper block, which is heated from behind by a set of halogen lamps and the second is inserted into the sample holder from the side while sliding it in to the sample station. The first is controlled by software, which—based on current temperature—turns on/off the lamps.

7. Experiment procedure
The sample material is first cut into a size of 65 × 75 mm and inserted between two frame plates and screwed together. Then it is inserted into the so-called load lock chamber of the CIF. It is mounted on the end of a magnetic arm. The lock is then closed and sealed, and the chamber is evacuated. Once the pressure in the chamber reaches the level of vacuum pressure in the irradiation chamber (10\textsuperscript{-8} mbar), a valve separating the two can be opened. The sample holder with the specimen is then transferred to the sample station where it can be exposed to radiation sources of the CIF.

B. VUV Exposure
Three samples of the ITO-coated aluminized polyimide thin films (ITO#2, ITO#3 and ITO#4) were exposed to VUV light using the Argon VUV source. A gas flow of 1200 sccm was chosen to reach the highest acceleration factor of 36.86 at 40 nm. The experiments are equivalent to 20, 30, and 50 days, respectively (see Table I), in interplanetary space conditions (i.e., the region between celestial bodies that is far enough from Earth and other planets for the influence of the interplanetary plasma to be the dominant factor in materials degradation) at a distance of 1 AU from the Sun at 40 nm.\textsuperscript{[8]}

C. Protons Exposure
Ten samples of ITO-coated aluminized polyimide thin films were exposed to different proton fluence values (see Table I) using the linear proton accelerator with proton energies of 3 keV (ITO#9, ITO#10, ITO#12, ITO#13 and ITO#14) and 5 keV (ITO#5, ITO#6, ITO#7, ITO#8, ITO#11). The average fluence was 6.33 × 10\textsuperscript{17} p\textsuperscript{-2}cm\textsuperscript{-2} for an energy of 5 keV and 6.20 × 10\textsuperscript{17} p\textsuperscript{-2}cm\textsuperscript{-2} for an energy of 3 keV. The exposure times, lasting from hours to days, were chosen to obtain five equivalent fluence values for both energies.
III. RESULTS AND DISCUSSION

A. Overview of Sample Degradation

Each sample was carefully investigated with the naked eye before and after exposure. The samples exposed to VUV radiation had their whole surface irradiated. We report no visible degradation of the samples. A limited area of each sample was irradiated by low-energy protons. We report a visible color change caused by the proton bombardments (see Figure 3).

We also examined the samples with the field-emission and atomic force microscopes (FEM and AFM) to identify any defects induced by the irradiation. In both cases, there was no apparent change to the surface structure that could be directly linked to degradation effect caused by the VUV light or the proton bombardment.

B. Spectral Reflectance Change

Figures 4 and 5 show changes in spectral reflectance of the samples exposed to protons of an energy of 3 keV or 5 keV respectively. Low-energy protons clearly cause a general decrease of the spectral reflectance. The change is especially significant in the visible and infrared regions of the electromagnetic spectrum. For both energies, the increase in proton fluence up to above $5 \times 10^{19}$ protons cm$^{-2}$ (ITO#14 and ITO#8 respectively) results in a strong reduction of the reflectance. Then, with further rise in proton fluence, at first, the reflectance shows a temporary reverse of the reduction trend to then maintain a less intense but steady decrease.

C. Thermo-Optical Properties

Table I presents $x_s$ for all samples. The change in the thermal emission coefficient after irradiation was less than 1.5 pct. Thus, only the value for a non-irradiated material is shown.

1. VUV irradiation

A first noticeable change in $x_s$ was observed after an equivalent of 50 days in space at distance of a 1 AU from the Sun. These results suggest a VUV radiation hardness of this material at very short wavelengths for short-term exposures.

2. Proton irradiation

Low-energy protons caused a significant increase in the solar absorption coefficient when compared to the non-irradiated material. Figure 6 presents $x_s$ change with increasing protons fluence for both energies of 3 keV and 5 keV. The increase in the solar absorption coefficient varies from 18 pct for ITO#14 to 36 pct for ITO#5 at 3 keV and from 28 pct for ITO#8 to 49 pct for ITO#5 at 5 keV when compared to a non-irradiated sample.

Additionally, there is a nearly perfect scaling between the solar absorption coefficients for equivalent fluence values at 3 keV and 5 keV as follows:

\[
x_s = 1 - \frac{\int_{2500 \text{nm}}^{500 \text{nm}} R(\lambda) S(\lambda) d\lambda}{\int_{2500 \text{nm}}^{500 \text{nm}} S(\lambda) d\lambda}, \tag{1}
\]

and

\[
\varepsilon_r = \frac{\int_{20 \mu m}^{50 \mu m} A(\lambda) E(\lambda) d\lambda}{\int_{3 \mu m}^{5 \mu m} E(\lambda) d\lambda}, \tag{2}
\]

where $R(\lambda)$ is the spectral reflectance after 100 pct reference correction, $S(\lambda)$ is the spectral solar irradiance, $A(\lambda) = 1 - R(\lambda)$, and $E(\lambda)$ is the black body emittance spectrum at 300 K.

We used a VERTEX 80v vacuum FTIR spectrometer from Bruker, equipped with integrating spheres with a sample holder on the sidewall. The optical measurements were performed ex situ. Reflectivity was compared to the known mirror standards. The uncertainties were calculated as the standard deviation from series of repeated spectral reflectance measurements in a given range.
Fig. 3—ITO-coated aluminized polyimide thin films, samples after exposure to proton energies of 3 keV (ITO#13 and ITO#10) and 5 keV (ITO#5 and ITO#11), respectively. The darker areas were irradiated and show a visible change in color when compared to the non-irradiated part of the sample. For more details of the setup and parameters of the experiments, see text and Table I.

Fig. 4—Spectral reflectance of the investigated samples of aluminized Kapton film with an ITO coating after 3 keV protons irradiation in comparison to the non-irradiated sample. See Table I for the values of the fluence and the corresponding irradiation time in space.

Fig. 5—Spectral reflectance of the investigated samples of aluminized Kapton film with an ITO coating after 5 keV protons irradiation in comparison to the non-irradiated sample. See Table I for the values of the fluence and the corresponding irradiation time in space.
Degradation in solar absorption, as expected from the reflectance spectra, generally increases when the fluence value rises. We also observe a temporary decrease of $x_{s}$ (ITO#9 and ITO#7 respectively). The increase is steeper for the first two fluence values than for the two last ones.

$$x_{5\text{keV}} = 1.1 \times x_{3\text{keV}}. \quad [3]$$

IV. DISCUSSION

A. Ultraviolet

MLI was reported by Sharma and Sridhara[1] to degrade after UV, sequential UV, and charged particle irradiation exposure. A significant increase of $x_{s}$ and only a slight rise in $\varepsilon_{s}$ were observed. However, UV sources available in most facilities used for material tests do not allow irradiation with wavelengths shorter than approx. 115 nm (an intensity cut-off of a common Deuterium lamp). Despite the fact that VUV radiation (10 to 200 nm) constitutes only about 0.1 pct of the near ultraviolet (NUV) irradiation intensity, it can still significantly contribute towards material degradation. The energy of a single photon in VUV range is high enough to cause photoionization and photodissociation effects that cannot be generated by NUV photons. Thus, one would expect a degradation of materials exposed to VUV radiation, especially for longer periods of time. Our results show no degradation of materials exposed to VUV radiation, especially for longer periods of time. The results of exposure of ITO-coated aluminized polyimide thin films to VUV radiation in our facility suggest a VUV radiation hardness of this material at very short wavelengths for short-term exposures. According to the presented results, low-energy protons cause a significant rise in the solar absorption coefficient of ITO-coated aluminized polyimide thin films.

V. CONCLUSION

Further studies of degradation of the material caused by the VUV radiation would require much longer exposures of the films to get more reliable information. As discussed by Li et al.[5] forming new chromophores and auxochromes could cause the visible color change of the samples induced by protons and the increase in absorbance in the visible region. However, we report the temporary healing-like effect between samples ITO#13 and ITO#9 for a proton energy of 3 keV, and ITO#8 and ITO#11 for a proton energy of 5 keV (see Figure 6), despite increasing the fluence. Thus, to explain the observed effect in a quantitative and qualitative manner, further research, both experimental and theoretical, is required. We intend to perform further irradiation of aluminized Kapton thin films with low-energy protons. Additionally, since Ding et al.[14] reported an increase in solar absorptance of ITO-coated aluminized Kapton under electron irradiation, an investigation of synergic influence of both protons and electrons on aluminized Kapton thin films would contribute to our understanding of the changes in internal structure and surface properties of the material when exposed to space conditions. Other aspects worth addressing are a possible influence of the temperature of the sample and the acceleration factor.
The main motivation for using the Complex Irradiation Facility is to extend the existing results of ground experiments with a unique proton energy range (down to below 10 keV) and short-wavelength light (down to 40 nm).

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