HIGH RESOLUTION OBSERVATIONS OF CHROMOSPHERIC JETS IN SUNSPOT UMBRA

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ABSTRACT

Recent observations of a sunspot’s umbra have suggested that it may be finely structured on a subarcsecond scale representing a mix of hot and cool plasma elements. In this study, we report the first detailed observations of umbral spikes, which are cool jet-like structures seen in the chromosphere of an umbra. The spikes are cone-shaped features with a typical height of 0.5−1.0 Mm and a width of about 0.1 Mm. Their lifetime ranges from 2 to 3 minutes and they tend to re-appear at the same location. The spikes are not associated with photospheric umbral dots and they instead tend to occur above the darkest parts of the umbra where magnetic fields are strongest. The spikes exhibit up and down oscillatory motions and their spectral evolution suggests that they might be driven by upward propagating shocks generated by photospheric oscillations. It is worth noting that triggering of the running penumbral waves seems to occur during the interval when the spikes reach their maximum height.

Key words: Sun: chromosphere – Sun: photosphere
Online-only material: animations, color figures

1. INTRODUCTION

Recent advances in observations and modeling created a picture of a sunspot as being a very dynamic and complex magnetic structure (e.g., Shimizu et al. 2012; Scharmer et al. 2011; Rempel & Schlichenmaier 2011; Rempel 2011). A closer look at the dark sunspot umbra using photospheric spectral lines revealed the detailed structure of the bright and nearly circular small intensity patches called umbral dots (UDs; Kilcik et al. 2012; Louis et al. 2012, and references therein). The magnetic fields in UDs are weaker than in their darker surroundings and UDs show plasma up-flows of a few hundred m s$^{-1}$ (Wiehr & Degenhardt 1993; Socas-Navarro et al. 2004; Rimmele 2004, 2008; Watanabe et al. 2009, 2012). According to realistic three-dimensional (3D) simulations (e.g., Schüssler & Vögler 2006; Bharti et al. 2011), UDs result from magneto-convection in the sunspot umbra and represent narrow convective up-flow plumes with adjacent down-flows, which become almost field-free near the surface layer.

The umbra appears to be much more dynamic when observed using chromospheric spectral lines. Based on recent high resolution observations, it has been suggested that the sunspot’s umbra may be finely structured, and consists of hot and cool plasma elements intermixed at subarcsecond scales (Socas-Navarro et al. 2000, 2009; Rouppe van der Voort & de la Cruz Rodríguez 2013). The well-known three-minute sunspot oscillations (e.g., Bogdan 2000; Maurya et al. 2013, and references therein) often lead to the appearance of bright umbral flashes (UFs), which are emissions in the cores of chromospheric lines caused by hot shocked plasma (Bard & Carlsson 2010) and which have been reported to display a filamentary structure (Socas-Navarro et al. 2009). Later studies also found evidence for a two-component structure of the umbra (Centeno et al. 2005; de la Cruz Rodríguez et al. 2013; Rouppe van der Voort & de la Cruz Rodríguez 2013).

Bharti et al. (2013) recently reported on transient jet-like structures seen in the Ca II H images of the sunspot umbrae, which they called umbral microjets. The microjets appear to be aligned with the umbral field, and no one-to-one correspondence between the microjets and the UDs was found. Rouppe van der Voort & de la Cruz Rodríguez (2013) described small-scale, periodic, jet-like features in the chromosphere above sunspots, which result from long-period waves leaking into the chromosphere along inclined sunspot fields.

In this paper, we present the first detailed observations of the fine-scale chromospheric phenomena in a sunspot’s umbra using Hα imaging spectroscopy data obtained with the New Solar Telescope (NST; Goode et al. 2010) that operates in Big Bear Solar Observatory. The data allowed us to fully resolve ubiquitous dynamic umbral Hα jet-like features and measure their general properties.

2. DATA

On 2013 June 14, the NST was pointed at the leading spot of active region NOAA 11768 located at S11W46 (Figure 1). Simultaneous observations were performed for the titanium oxide (TiO) band using the Broadband Filter Imager and at the Hα spectral line using the Visible Imaging Spectrometer (VIS). The photospheric images of the sunspot were acquired every 25 s using a 1 nm passband TiO filter centered at 705.7 nm with a pixel scale of 0'0375. This absorption line (the head of the TiO γ-system) is only formed at low temperatures below 4000 K, i.e., inside sunspots (see Figure 10 Berdyugina et al. 2003).

The VIS combines a 0.5 nm interference filter with a Fabry–Pérot etalon to produce a bandpass of 0.007 nm over a 70″ field of view (FOV) centered at the Hα spectral line. Imaging of the chromosphere was performed at 11 positions along the spectral line with a 0.02 nm step along the spectrum and a pixel size of 0'029. The difference in the acquisition time at two sequential line positions (e.g., +0.02 nm and −0.02 nm) was about 2 s. The 11 point line scan was recorded every 25 s along the following sequence: −0.1, +0.1, −0.08, +0.08, −0.06, +0.06, −0.04, +0.04, −0.02, +0.02, 0.0 nm. At each line position, we acquired a burst of 25 images with exposure times ranging from 7 ms (at −0.1 nm) to 25 ms (at the line center). These bursts were used for speckle reconstruction, so that each line position was speckled separately. To estimate an error in the line profile, we used the 50 profiles shown in the last panel of Figure 5. Each profile was constructed from a series of 11
Figure 1. Main spot of NOAA AR 11768 as seen in the chromospheric Hα line (top and lower right panels) as well as the photospheric TiO 705 nm line (lower left). The top row shows off-band images of the sunspot taken at −0.02 nm (left, blue) and +0.02 nm (right, red) off the line center. All of the images were unsharp masked. The D.C. arrow in the lower left image points toward the disk center. The two line segments indicate the location of xτ cuts. The large circles indicate several bright photospheric umbral dots with fine structures inside them. The small circles mark the base of several spikes (see text for more details). The long tick marks separate 1 Mm intervals. The images were not corrected for a projection effect.

Animations and a color version of this figure are available in the online journal.

speckle-reconstructed off-band Hα images (see above), so that the errors introduced by local misalignment and residual seeing are taken into account. These profiles were measured inside of a stable, slowly evolving super-penumbra canopy, so that we accept that all variations in the line profiles are due to varying seeing as well as errors introduced by the speckle reconstruction code. We calculated the relative standard deviation using 50 data points at each line position and found that it ranges from 2% to 9% with an average of 5%.

All images were acquired with the aid of an adaptive optics (AO) system, which incorporates a 357 actuator deformable mirror, a Shack–Hartmann wavefront sensor with 308 sub-apertures, and a digital signal processor system. The Kiepenheuer-Institut für Sonnenphysik’s software package for speckle interferometry of AO corrected solar data (KISIP; Wöger & von der Lühe 2007) was applied to all of the acquired images as a post AO reconstruction algorithm, allowing us to achieve the diffraction limit of the telescope over a large FOV.

3. RESULTS

In Figure 1, we show contrast-enhanced (unsharp masked) Hα−0.02 nm (top left), Hα+0.02 nm (top right), Hα (lower right), and TiO (lower left) images of the sunspot. The dynamics of the umbral area is best seen in the Hα−0.02 nm animation (S1.mp4), where the diffuse, dark, and dynamic intensity fronts are seen to rapidly travel across the umbra, seemingly reflecting back and forth from the umbra-penumbra boundary triggering penumbral waves. The animation images were not processed with the unsharp masking method. The chromospheric “surface” of the umbra seems to be vertically oscillating, giving an impression of standing surface waves in an enclosed body of fluid. The bright diffuse patches seen within the umbra perimeter
around the three central mid-size circles are UFs. At a spatial resolution of \( \sim 0.1'' \), the UFs do not show any intrinsic fine structure. Instead, the bright patches appear to be interrupted by vertical dark spikes projected onto the surrounding UFs background.

According to Figure 1, these umbral spikes can be detected everywhere within the umbra with nearly the same number density, although it seems that they preferably occur in darker parts of the umbra. Using the large circles, we indicate in the TiO image several bright photospheric UDs that have dark lanes and/or dots inside of them (see Rimmele 2008). Comparison with the chromospheric images shows that in four cases out of five, an umbral spike can be found in the vicinity of a UD. It is possible that the surrounding features simply overlap with the UDs, since we did not find strong evidence in those cases for the spikes to originate above UDs. We also plotted small circles at the base of several spikes in the off-band images and it appears that all of them map back to dark areas in the photospheric umbra.

In Figure 2, the position of the base of the umbral spikes (yellow dots) is plotted over an image of the umbra as observed at H\( \alpha -0.1 \) nm. There were 101 spikes identified in total in this image and the position of their base was determined manually using the H\( \alpha -0.02 \) nm image shown in Figure 1. This plot and Figure 2 further show that the spikes tend to avoid bright UDs and are either co-spatial with dark umbral patches or are found in the vicinity of a UD. This tendency probably explains the presence of spikes near large UDs. In Figure 3, we plot the brightness distribution (solid line) determined for the areas enclosed by the black contour plotted at an intensity level of \( 1.5 \times 10^3 \) DN. The distribution peaks at the \( 1.0 \times 10^3 \) DN intensity level marked by the white contour line. The dotted line shows the base intensity distribution measured under the base of the umbral spikes. In this plot, the base intensity distribution was multiplied by a factor of 400 to make these two curves comparable. The distributions clearly show a preference for spikes to originate in the dark umbral patches.

The spikes have a cone-like appearance with a wide base (typically \( \sim 0.1 \) Mm) and a pointed tip. Their typical length is between 0.5 and 1.0 Mm. Using the off-band images in Figure 1 and the line scan data (see S2.mp4 animation), we conclude that (1) umbral spikes simultaneously show up in the blue and red wings of the H\( \alpha \) spectral line, (2) their number density rapidly decreases when observed further away from the line center, and (3) they show oscillatory up and down motions. When observing near the H\( \alpha \) line center, the umbra is seen to be filled with chromospheric plasma forming various extended structures, so that the spikes are less prominent and only their dark tips can be identified when carefully comparing off-band and line center images.

Leenaarts et al. (2012) performed MHD simulations and 3D non-LTE radiative transfer computations to understand the
details of the Hα line formation. They reported that Hα opacity in the upper chromosphere is mainly sensitive to the mass density and only weakly sensitive to the temperature. Also, the intensity of the Hα line core is related to the formation height, and the intensity decreases as the formation height is shifted toward the upper chromosphere. In their simulations, the fibril-like Hα structures represent ridges of enhanced mass density, which displace the formation level of the Hα line-core to higher levels, so these structures appear darker. Based on these simulation results, we interpret umbral spikes as plasma structures with enhanced mass density.

In Figure 4, we plot Doppler shift maps obtained by subtracting red wing Hα+0.02 nm images from the corresponding blueshifted images taken 2 s prior. In general, the Doppler shift maps are finely structured but do not show the presence of any strong directed out- or down-flows associated with the spikes. Instead, they are seen as very low contrast Doppler shift features. The redshifted Doppler features seen in these maps at the limb side of the umbra and penumbra appear to be longer (1–2 Mm) and they are mostly seen in the projection on the sunspot’s penumbra. The sunspot was located at a longitude of 45 deg west, and therefore near the disk side umbra-penumbra boundary, where the field inclination is nearly 45 deg and the line of sight is approximately parallel to the field lines, while at the limb side of the boundary the line of sight is nearly perpendicular to the magnetic field lines. Thus, because of the projection effect, the spikes at the limb side are more easily observed as compared to their point like projection at the disk side of the umbra.

The four Doppler maps in Figure 4 illustrate that the location of the redshifted thorns within the umbra varies rapidly in time, which is consistent with the idea of their wave origin. The observed behavior suggests that the spikes may have a wave origin rather than being linear jet-like outflows from the sunspots umbra. As follows from the maps, the intensity of the redshifted fine structures evolve in time, however, no equally strong blueshifted signal appears instead as would be expected from simple oscillatory motions. The distribution patterns of these
Figure 5. Examples of spectral evolution as measured for two different umbral spikes (panels (a)–(b)). Panels (c) and (d) show spectral evolution immediately next to a spike and in UFs, correspondingly. Panel (e) represents spectral evolution in a super-penumbral filament outside the sunspot. Fifty line profiles acquired within a 20 minutes time span were used to produce these plots.

redshifted features varies rapidly in time, which is consistent with the idea of their wave origin. We therefore interpret these jets as the outflows in the sunspot dynamic fibrils (DF) described in Rouppe van der Voort & de la Cruz Rodríguez (2013).

In Figure 5, we show a $\lambda$ vs. $t$ plot representing the time evolution of the Hα line profiles measured at various locations inside the umbra. Panel (e) is shown for comparison and it represents a co-temporal spectral evolution measured in a super-penumbral filament outside the sunspot. The (a)–(b) profile stacks were measured inside the spikes and they show typical signatures of shock driven flows (Hansteen et al. 2006; Langangen et al. 2008), i.e., the initially blueshifted spectral line gradually becomes redshifted due to deceleration and the subsequent fall of the previously upward moving chromospheric plasma (see, e.g., panel (b) between $t = 13$ and $t = 17$ minutes). The diagonal spectral features seem to repeat approximately every 2–3 minutes, although the line intensity and the spectral shift vary over time and from one location to another. Panel (c) shows the spectral line evolution measured immediately next to the spikes. The shocks are present there too, although the plasma seems to be compressed to a lesser degree as evidenced by the much shallower line profiles. Panel (d) shows similar data but for pixels belonging to UFs. The diagonal shock features are largely absent in this case. Instead, periods of emission (e.g., at $t = 8–9$ minutes) and absorption (e.g., at $t = 6–5$ minutes) follow each other with a periodicity of about 2–3 minutes.

In Figures 6 and 7, we show time variations of the chromospheric intensity along the cuts indicated in Figure 1 by two line segments. The cuts were made along the axis of spikes and the width of the slit was 0.32 Mm (15 pixels=0′′43). The repeating dark half-parabola feature seen above 1 Mm is the

Figure 6. Top panel: time–distance plots showing the chromospheric dynamics as it is seen at the Hα−0.02 nm along the longer slit in Figure 1. The width of the slit was chosen to be 0.32 Mm. The white curve in both panels has a three-minute period and is plotted for illustration only. The white box outlines an area enlarged in Figure 7. The bottom panel shows similar dynamics, however, observed at the Hα−0.06 nm along the shorter slit.

Figure 7. Enlarged fragment of the $x$–$t$ plot shown in the top panel of Figure 6. Arrows indicate umbral spikes.
propagating front of penumbral waves that are excited at the umbra-penumbral boundary (\( x = 1 \text{ Mm} \)) and travel toward the outer bounds of the penumbra. The umbral spikes can be seen here as the dense and short features visible below \( x < 0.5 \text{ Mm} \) and between \( t = 50 \) and \( t = 58 \text{ min} \) (Figure 7) with a lifetime of about 3 minutes. The tip of the spikes appears to be rising up with velocities of about 4–5 \( \text{km s}^{-1} \), which is expected to reflect the local sound speed and is in the range of the reported velocities associated with oscillations and shock formation (e.g., Steiner et al. 1998; Maurya et al. 2013); the temporal resolution of the data did not allow us to make a more detailed conclusions about their dynamics. After \( t = 60 \text{ min} \), they briefly appear nearly every 3 min approximately at the moments of the minimum in the white curve (i.e., at \( t = 66, 69, 72, ..., 84, \) and 87 minutes). In general, the data set gives an impression that the spikes tend to re-occur at the same locations at fairly regular time intervals; however, their parameters vary, possibly reflecting the strength of local oscillations. The lower plot in Figure 6 shows similar features and behavior except that it was made using H\( \alpha - 0.06 \text{ nm} \) data. The penumbral waves are still present, although their front is less dense and narrower. The same is true about the umbral spikes, which are only weakly seen in the plot (i.e., at \( t = 65, 67, 73, 76, \) and 79 minutes).

It is worth noting that each occurrence of the spikes in the center of the penumbra seems to be correlated in time with the onset of a new penumbral wave: the wave begins when spikes are at the peak of their development and is best displayed in the lower panel of Figure 6 between \( t = 63 \) minutes and \( t = 81 \) minutes. The spikes are visible under the white curve and the middle of their lifetime occurs when the white curve is in the maximum. At the same moment of time, a dark front, representing penumbral waves, begins to propagate above the white curve. This correlation may indicate that the spikes and the running penumbral waves are triggered and powered by the same processes that appear to be coherent on scales comparable to the size of a sunspot, such as sunspot oscillations.

4. CONCLUSIONS

We summarize our findings as follows. High resolution NST observations of the solar chromosphere using the H\( \alpha \) spectral line revealed the existence of dynamic spike-like chromospheric structures inside sunspot’s umbra, which we called umbral spikes. The spikes are on average about 0.1 Mm wide, their height does not exceed 1 Mm. They are mainly vertical with a slight tendency to fan out closer to the periphery of the umbra. Since umbral fields show a similar inclination distribution, we suggest that the umbral spikes are aligned with the umbral magnetic fields. The spikes show a nearly uniform distribution over the umbra with a tendency to be more concentrated in its darkest parts occupied by the strongest fields in the umbra, while UDs are considered to be field-free magneto-convection features. Thus, our findings seem to indicate that the umbral spikes may be co-spatial with strong magnetic flux concentrations rather than with the weaker magnetic fields above UDs.

These first detailed observations of umbral spikes presented here suggest that they are a wave phenomena and result from sunspot oscillations. Bharti et al. (2013) observed transient jet-like structures in Ca \( \text{II} \) images of sunspot umbrae, which they called umbral microjets. The authors speculated that the microjets, which are shorter than 1 Mm and not wider than \( 0.3 \), may be either upflow jets driven by the pressure gradient above the photospheric UDs or they may be caused by reconnection of hypothetical opposite polarity fields that might exist around large UDs. Although the H\( \alpha \) spikes and the microjets are of a comparable size, the lifetimes of spikes appear to be longer (2–3 minutes) than those of microjets (50 s). While a possible relationship between them has yet to be established, it seems that the two phenomena are different and the scenarios presented in Bharti et al. (2013) may not be able to explain the spikes, since the latter show a preference to be more frequent in the darkest cores of the umbra, dominated by strong fields and devoid of large and bright UDs.

In a very recent paper, Rouppe van der Voort & de la Cruz Rodríguez (2013) described DF observed in the chromosphere above a sunspot. Their \( \alpha \) plot generated from a series of H\( \alpha \) line center images (see Figure 3 in Rouppe van der Voort & de la Cruz Rodríguez 2013) shows the presence of parabolic intensity features, which were interpreted by these authors as being very short jet-like features precisely above the umbra. At the same time, individual umbral spikes were not resolved in their data and only diffuse, low contrast dark specks can be distinguished inside the umbra in the corresponding H\( \alpha \) images. The umbral spikes and the DFs may be the same wave phenomenon, while their differences in appearance are caused by the fine scaled structures in magnetic fields, viewing angle, and the resolution of the data.

Heggland et al. (2011) studied wave propagation in different magnetic configurations. They found that the peaks of power of the 3 minute oscillations and the high amplitudes of the vertical velocities (5–8 \( \text{km s}^{-1} \)) are located above strong photospheric flux concentrations or, in other words, inside vertical flux tubes. They also found that rising and falling jets, which form as a result of the oscillations, have their axes aligned with the magnetic axes of field concentrations. The general appearance of the simulated jets and their oscillatory motions are strikingly similar to those of the spikes (see Figure 25 in Heggland et al. (2011) and Figure 1 in this paper), although the simulated jets appear to be, on average, more extended (0.5–6.0 Mm), longer living (2–5 minutes), and show higher velocities (10–40 \( \text{km s}^{-1} \)). Thus one possible interpretation of the observed phenomena is the penetration of photospheric oscillations into the chromosphere along thin and vertical magnetic flux tubes.

Using spectropolarimetric observations, Socas-Navarro et al. (2000) suggested the existence of an unresolved active component with upward directed velocities. More precisely, the anamalous polarization profiles could only be explained by emission from an unresolved mixture of an upward propagating shock and cool, slowly downflowing surroundings. Later, Centeno et al. (2005) specified that the active component is present throughout the entire oscillation cycle. They also inferred that the shock waves propagate into the umbra inside channels of subarcsecond width, which could be the flux tubes discussed above. Finally, Socas-Navarro et al. (2009) presented high resolution \( \text{Hinode} \) data on UFs, concluding that UFs show fine filamented structure.

We argue that the \( \alpha \)’t wide umbral spikes may be the unresolved active component of the sunspot’s umbra discussed above. They represent finely structured shocked plasma showing up- and down-flows of the order of 5–7 \( \text{km s}^{-1} \). The data seem to suggest that the spikes are associated with the interiors of strong flux tubes, thus confirming the idea about the existence of narrow channels conducting photospheric oscillations into the chromosphere (Socas-Navarro et al. 2000; Centeno et al. 2005).
Finally, Rouppe van der Voort & de la Cruz Rodríguez (2013) suggested that the fine structure of UFs is related to the sunspot DF. The new data presented here show that the dark filaments reported inside the UF area are the vertical umbral spikes projected onto the otherwise uniform and unstructured UFs.

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