Research Article
External Magnetic Field and Air Mass Effects on Carrier’s Effective Lifetime of a Bifacial Solar Cell under Transient State

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Abstract: This paper presents a theoretical study of the external magnetic field and the air mass effects on the carrier density, the transient open circuit voltage decay and the effective minority carrier’s lifetime in a polycrystalline silicon solar cell. From a theoretical approach based on the columnar model of the grains and the quasi-neutral base, the three-dimensional diffusion equation is established and the boundaries conditions are defined in order to use Green’s functions to solve this equation. New analytical expressions of carrier’s density and transient voltage are also found. The magnetic field and the air mass impacts on transient electrons density and transient open circuit voltage are then analyzed. The effective minority carrier lifetime is extracted from the curve versus time of transient open circuit voltage decay for different values of magnetic field and air mass.

Keywords: Air mass, bifacial, effective lifetime, magnetic field, solar cell, transient state

INTRODUCTION

In order to reduce dependence on fossil fuels, renewable energy sources have been studied intensively for many years (Ellabban et al., 2014). As solar energy, one major source of renewable energies, depends on the fluctuating solar radiation and solar cell properties.

However, some external factors as magnetic field (Betser et al., 1995; Erel, 2002; Madougou et al., 2007) external electric field (Coors et al., 1998; Pelanchon et al., 1992; Dieng et al., 2009), intense light (Pelanchon et al., 1992; Agroui et al., 2003) and internal factors as grain size (Samb et al., 2009), grain boundary recombination velocity (Ba et al., 2003), air mass (Rida et al., 2016; Shnishil et al., 2011) electric field in the bulk due to carrier concentration (Pelanchon et al., 1992) can influence the solar cell quality.

In this study, we present a three-dimensional study of a bifacial solar cell under external variable magnetic field and air mass. From a theoretical approach, the calculations were carried out in the case of a polycrystalline device with columnar model of grain orientation under front side illumination and the quasi-neutral base. The three dimensional diffusion equation is established and the boundaries conditions are defined in order to use Green’s functions to solve this equation.

METHODOLOGY

In this approach, new analytical expressions of excess minority carriers density and transient open circuit voltage were established for a front side illumination.

The effects of magnetic field and air mass on excess minority carriers generation, transient voltage decay and effective lifetime profile are then analyzed.

Considering the weak thickness of the emitter and zone of space load, we neglected their contributions to the photocurrent, so the quasi-totality of the current is provided therefore by the base (Charles et al., 2000; Toure et al., 2012).

The grains are under parallelepipedic shape (2a; 2b; H) and the joints of grains are perpendicular to the junction.

The surfaces between two adjacent grains and perpendicular to the junction are characterized by the same carrier recombination process evaluated by a grain boundaries recombination velocity $S_{gx} = S_{gy} = S_g$.
The electric field of crystal lattice is negligible (Sam et al., 2012).

Only the contribution of the base in the process of generation is considered.

The magnetic field is oriented according to the direction \( \overrightarrow{OY} \), \( \overrightarrow{B} = B\overrightarrow{OY} \). It is therefore perpendicular to the depth of light penetration in order to observe the effect of the strength of Lorentz on the carriers.

The illumination is uniform. Then we have a generation rate depending only on the depth in the base \( z \).

The solar cell is front side illuminated under the external magnetic field with the tree type of air mass is illustrated on Fig. 1 (Sam et al., 2016).

In the base of the cell, the excess minority carriers are electrons and their density satisfies to the equation below:

\[
\frac{\partial \delta n}{\partial t} = \frac{1}{e} \nabla j_n + g_n - r_n
\]

(1)

where,

\[ \delta(\vec{r}, t) = \delta(x, y, z, t) \]

and,

\[ \theta = 1 + \left( \frac{\mu B}{D} \right)^2 \]

(4)

\( D' \) is the diffusion coefficient of excess minority carriers in the base of the bifacial cell in presence of magnetic field (Dieng et al., 2011; Flohr and Helbig, 1989); it can be expressed as:

\[ D' = \frac{D}{1 + \left( \frac{\mu B}{D} \right)^2} \]

(5)

\( D, \mu, B \) are respectively the coefficient of electrons diffusion, the mobility of electrons and the value of magnetic field applied.

The presence of the magnetic field in our model leads to new values of carriers diffusion length \( L' \) and carriers diffusion coefficient \( D' \) which depends on magnetic field. \( \theta \) is a coefficient which depends on magnetic field intensity.

The carriers’ generation rate under multispectral light at the depth \( z \) in the base can be written by the following expression:

\[
g(z,t) = \begin{cases} 
  n \sum_{n=1}^{3} a_n \exp(-b_n z) \text{ si } 0 \leq t \leq T_e \\
  0 \text{ si } t > T_e 
\end{cases} 
\]

(6)

In this expression of \( g(z,t) \), \( n = \frac{I_{sun}}{I_{eco}} \) indicates the illumination level (sun number) and \( n = 1 \) for our study.
The coefficients $a_m$ and $b_m$ are the modeling coefficients of AM1.5, AM1 and AM0 of air mass (Mohammad, 1987; Furlan and Amon, 1985).

Equation (1) is solved with the following boundary conditions:

- At the junction $z = 0$:
  \[
  \frac{\partial \delta(r,t)}{\partial z} \bigg|_{z=0} = \frac{S_f}{D^*} \delta(r(z=0),t) \quad (7)
  \]

- At the rear side $z = H$:
  \[
  \frac{\partial \delta(r,t)}{\partial z} \bigg|_{z=H} = \frac{S_b}{D^*} \delta(r(z=H),t) \quad (8)
  \]

- At surfaces limited by $x = \pm a$ and $y = \pm b$
  \[
  \frac{\partial \delta(r,t)}{\partial x} \bigg|_{x=\pm a} = \pm \frac{S_g}{D^*} \delta(r(x=\pm a),t) \quad (9)
  \]

  \[
  \frac{\partial \delta(r,t)}{\partial y} \bigg|_{y=\pm a} = \pm \frac{S_g}{D^*} \delta(r(y=\pm b),t) \quad (10)
  \]

$S_f$, $S_b$ and $S_g$ are the recombination velocity of minority carriers respectively at surfaces $z = 0$, $z = H$ and $x = \pm a$ (or $y = \pm b$).

$a$, $b$ and $H$ are the grain sizes as indicated on Fig. 1.

A solution of Eq. (1) is given by Eq. (11) according to the author of the reference (Sam et al., 2016):

\[
G = \sum_{i=1}^{\infty} \sum_{j=1}^{\infty} \sum_{k=1}^{\infty} M_i M_j M_k \exp\left(-\beta(t-t')\right) \quad (11)
\]

where,

\[
M_i = A^2_{k_i} \cos(k_i x') \cos(k_i x) \quad (12)
\]

\[
M_j = A^2_{l_j} \cos(l_j y') \cos(l_j y) \quad (13)
\]

\[
M_k = A^2_{\mu_k} \cos(\mu_k z' + \varphi_k) \cos(\mu_k z + \varphi_k) \quad (14)
\]

$\beta'$ and $l_j^*$ are expressed by:

\[
\beta' = D' \left(k^2_{i} + l^2_{j} + \mu^2_{k} + \frac{1}{L^2}\right) \quad (15)
\]

\[
l_j^* = \frac{l_{j}}{\sqrt{\theta}} \quad (16)
\]

The quantities $A_i$, $A_j$ and $A_{\mu_k}$ are obtained by normalizing $M_i$, $M_j$ and $M_k$. The parameters $k_i$, $l_j$ and $\mu_k$ are the eigen values obtained from the boundaries conditions. $\varphi_k$ is the initial phase and obtained by solving the equation:

\[
\tan(\mu_k H + \varphi_k) = \frac{S_b}{\mu_k D^*} \quad (17)
\]

The transient voltage decay is defined by:

\[
V(t) = V_i F_v(k_i, l_j^*, \mu_k) r \exp(-\beta(t-T_e)) \quad (18)
\]

where

\[
r = \exp\left(\frac{\Delta V}{V_r}\right) - 1 \quad (19)
\]

and

\[
\Delta V = V_o - V_F \quad (20)
\]

The quantities $\Delta_o(0,0)$ and $\Delta(0,0)$ are defined by:

\[
\Delta_o(0,0) = \int_{-a}^{b} \int_{-b}^{b} Z_{111}(x,y,0) dxdy \quad (22)
\]

\[
\Delta(0,0) = \int_{-a}^{b} \int_{-b}^{b} d_{111}(x,y,0) dxdy \quad (23)
\]

$Z_{111}(x,y,0)$ and $d_{111}(x,y,0)$ are respectively the spatial component of $\delta(x,y,z,t)$ and the minority carriers density during the phase of illumination:

\[
l_j^* = \frac{l_{j}}{\sqrt{\theta}} \quad (24)
\]

According to expression (18), we notice two types of $V(t)$ decays:

- If:
\[ Fv(k; l^*; \mu_t) r \exp(-\beta^*(t - Te)) \ll 1 \]  
(25)

The time dependent tension is rewritten in the following way Eq. (26):

\[
V(t) = V_r \left( -\beta^*(t - Te) + \ln\left(1 + Fv(k; l^*; \mu_t) r\right) \right)
\]  
(26)

This is a linear function of the time with a negative slope \(-V_r \beta^* \):

- If:

\[ Fv(k; l^*; \mu_t) r \exp(-\beta^*(t - Te)) \ll 1 \]  
(27)

\( V(t) \) can be rewritten in the form:

\[
V(t) = V_r Fv(k; l^*; \mu_t) r \exp(-\beta (t - Te))
\]  
(28)

\( V(t) \) is a time dependent decay exponential function.

**RESULTS AND DISCUSSION**

We present here the simulation results obtained from the previous modeling equations by using the Origin pro and Mathcad software.

**Effects of magnetic field and air mass on carrier density:** Figure 2 and 3 present respectively curves of carriers density versus air mass and versus magnetic field. This figure shows a quasi-exponential decay versus time of the electrons density.

This decay of the Fig. 2 can be explained by the reducing of the generation rate during the time to the different air mass. So when the air mass increases, the direction of normal irradiance decreases reducing carriers generation.

A quantitative difference is observed on the curves corresponding to these three values of air mass. These differences are more sensitive on the amplitudes of the curves, i.e., during the illumination phase. So, hole-electron pair generation rate and air mass number varies in reverse.

Figure 3 presents the variation of carriers’ density versus magnetic field and time. These curves show that for magnetic field increase, the maximum of the excess minority carriers density increases also and are shifted left to the junction. This increasing of the maximum of the excess minority carriers density translate an increase of the carrier concentration.

The increase of carriers density cross to the junction with the magnetic field is more important; so this displacement of the maximum of the excess minority carriers’ density shows that there is a decrease of carriers flow through the junction. We can also explain this by the reducing of recombination at the grain boundaries because of the magnetic field.

**Effects magnetic field and air mass on transient voltage:** We use the expression of transient voltage to show the influence of air mass and magnetic field on transient voltage by the following figures.

Figure 4 presents the curve of transient voltage decay when the air mass number air increases but we observe the same variation for air mass 0 and 1. This behavior is explained by the substantially identical values of air mass 0 and air mass 1. The decrease of transient voltage to air mass 0 to air mass 1.5 is explained by an important photogeneration of excess minority carriers under AM0 than AM1.5.

The increase of air mass creates the decrease of normal irradiance by reducing illumination then the number of excess minority carriers. In other words, much charge carriers are stored near the junction under air mass 0 because much carriers are photo-generated. The profile of Fig. 5 presents the transient voltage increase when the magnetic field increases.

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Fig. 2: Carriers density curves versus time and air mass
Fig. 3: Carriers density curves versus time magnetic field

Fig. 4: Transient voltage decay curve versus time for various air mass

Fig. 5: Transient voltage decay curves versus magnetic field
The transient voltage which means an accumulation of charges across the junction presents a maximum for very low $S_f$ values (open circuit). The open circuit voltage increases with increasing magnetic field. This is a consequence of charges accumulation in the base and the diminution of junction recombination with increasing magnetic field. This phenomenon is accompanied by an increase of recombination in the bulk and to grain boundaries (Dieng et al., 2011; Zoungrana et al., 2007).

The behavior of the transient voltage under magnetic field characterizes an increase of resistive losses in the base (due to material structure and electrical grids). This increase of resistive losses leads to a decrease in the output voltage.

**Effects of magnetic field and air mass on effective lifetime:** The Eq. (26) gives the effective lifetime and permit us to obtain the values of effective lifetime by making linear fitting with the software of origin lab of the linear expression of transient voltage decay:

$$ m = \frac{V_T}{\tau_{eff}} $$

where,

$m$ : The slope of the linear expression of transient open circuit voltage

$\tau_{eff}$ : The effective lifetime

$V_T = 26mV$ : The thermic voltage

We use this form of transient open circuit voltage because the linear parts of the curves don’t present capacitance effects. The results of linear fitting are saving on the Table 1.

Table 1 shows the variation of effective lifetime versus magnetic field and air mass. The study of carriers’ density and transient voltage show that the air mass AM0 has the maximum of carriers which permit to obtain the transient voltage that we make the linear fitting to get values of effective lifetime.

The increase of magnetic field explains the decreasing of diffusion length and diffusion coefficient. This reducing of diffusion length and diffusion coefficient is responsible of this decreasing because the minority of carriers’ recombination is fast. So the time of recombination becomes short when the magnetic field increases. All these situations are responsible of the decreasing of effective lifetime.

Another analysis of this Table 1 presents an inchange of effective lifetime versus air mass number increase because the effective lifetime is independent on the generated rate or the photogeneration globally.

We notice that the effects of air mass are less important than the magnetic field. The same results of effective lifetime with the magnetic field are found by Sam et al. (2016).

**CONCLUSION**

In this study, a three dimensional approach of electrons diffusion in the p region of a polycrystalline silicon solar cell is presented. The solar cell is front side illuminated by a pulsed light under external magnetic field versus air mass values. The resolution of the time dependent and the three dimensional diffusion equation leads to new expressions of the transient voltage and the reduced amplitude of transient voltage.

This study shows the effects of air mass and magnetic field on carriers density, transient voltage decay and carriers effective lifetime. In summary, the limiting effects of magnetic field on the quality of solar cell are pronounced. The air mass does not affect the solar cell carriers effective lifetime.

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