The absolute motion of the peculiar cluster NGC 6791.*

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Abstract. We present improved values of the three components of the absolute space velocity of the open cluster NGC 6791. One HST ACS/WFC field with two-epoch observations provides astrometric measurements of objects in a field containing the cluster center. Identification of 60 background galaxies with sharp nuclei allows us to determine an absolute reference point, and measure the absolute proper motion of the cluster. We find \((\mu_\alpha \cos \delta, \mu_\delta)_{J2000.0} = (-0.57 \pm 0.13, -2.45 \pm 0.12) \text{ mas yr}^{-1}\), and adopt \(V_{\text{rad}} = -47.1 \pm 0.7 \text{ km s}^{-1}\) from the average of the published values. Assuming a Galactic potential, we calculate the Galactic orbit of the cluster for various assumed distances, and briefly discuss the implications on the nature and the origin of this peculiar cluster.

Key words. astrometry — open clusters: NGC 6791 — dynamics

1. Introduction

NGC 6791 is a unique object in our Galaxy. Usually classified as an open cluster, it has a number of peculiarities which leave its origin and its nature quite enigmatic. It is more massive (at least 4000 \(m_\odot\)), more metal-rich ([Fe/H] \(\sim +0.4\), Carraro et al. 2006, Gratton et al. 2006), and older (\(\approx 9\) Gyr, King et al. 2005) than most known open clusters. In contrast with other open clusters, which lie close to the Galactic plane, its distance of \(\sim 4000\) pc (King et al. 2005) and Galactic latitude of 11° put NGC 6791 \(\sim 1\) kpc above the plane. The cluster is also anomalous with respect to the radial abundance gradient and the age-metallicity relation of the Galactic disk. (See discussion in Carraro et al. 2006.)

All of these peculiarities make NGC 6791 both an interesting and a challenging object, and they stimulated us to undertake a study of it with deep HST imaging. In such HST programs we have, whenever possible, included a second epoch of observation, in order to use proper motions to separate cluster stars from the field. The second epoch also allows us, here, to study the motion of the cluster itself.

We have already published two papers based on the first-epoch observations alone. The first of them reported the discovery (Bedin et al. 2005) of an anomalous white dwarf cooling sequence. The second paper (King et al. 2005) was devoted to the main sequence, and included a preliminary mass function, which we found to be rather flat.

Along with our proper motion, the availability in the literature of radial velocities for a number of cluster members allows us to determine all three components of the absolute motion of NGC 6791 and to infer some properties of its orbit, which will shed some light on the possible origin of this object. A detailed study of the cluster main sequence (down to the hydrogen-burning limit) and of the white dwarf cooling sequence will be presented in forthcoming papers.

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2. Observations, data reduction, and proper motions

All of the observations used for the proper-motion measurement come from our HST programs GO-9815 and GO-10471 (PI King), which were separated by ~2 years. For precise astrometric measurements and a more accurate assessment of the errors, we had taken particular care to dither our images properly, with both whole-pixel and fractional-pixel offsets (following the general recipes given in Anderson & King 2000). Table I describes the ACS/WFC observations used in this work. The first of the two programs also included images in the F606W band, but we used only F814W images for our proper motions, in order to avoid any possible filter-dependent systematic errors.

We measured positions and fluxes for every star in every F814W _FLT_ exposure, using library effective PSFs and the software programs documented in Anderson & King (2006). We then generated a master list of all the stars, and collated all the observations of each star. As in Bedin et al. (2003), we used the best distortion corrections available (Anderson 2002, 2005) to correct the raw positions that we had measured from the _FLT_ exposures.

We carefully constructed the reference frame, as follows. We measured a simple centroid position for each bright star in the first _DRZ_ image, and found by least squares a linear transformation between those positions and the positions of the same stars in the corresponding _FLT_ image. We then rotated the frame of the _FLT_ image so that the y axis points exactly north, and rescaled it to agree with the pixel size of the _DRZ_ image, which is 50 mas/pixel. This is our reference frame. From the header of the _DRZ_ image we got the R.A. and dec of the point in our reference frame that corresponds to the position given in the header; this allows us to combine the internal accuracy of our PSF-measured positions with the absolute orientation and scale information from calibrated pipeline products.

As reference stars we identified cluster members in the color-magnitude diagram, and used only those stars for the transformation from each exposure into the reference frame. (See Bedin et al. 2003 for details.) We thus ensured that the proper motions are measured relative to the bulk cluster motion. Carraro et al. (2006) give an estimate of 2.2 ± 0.4 km s^{-1} for the internal dispersion of radial velocities; for a distance of 4 kpc, this corresponds to a proper-motion dispersion of ~0.1 mas yr^{-1} (~0.005 ACS/WFC pixel in two years). This means that the relative motions of the cluster stars should all be zero to within the measurement errors. We iteratively removed from the member list some stars that had field-star-type motions, even though their colors placed them near the fiducial cluster sequence. Field-star proper motions will be discussed very briefly at the end of this section.

Finally, in order to minimize the influence of any uncorrected distortion on transformations into the reference frame, we used for each object a local transformation based on the nearest ~50 well-measured cluster stars. With all these precautions, we found that for stars with >500 DN in their brightest pixel we could measure positions in a single image with an error <0.05 pixel in each coordinate. We note that the relatively high background in our images (~65 DN) makes the astrometric effect of charge-transfer efficiency negligible, and also that the red-halo effect that disturbs ACS/HRC images is negligible in WFC images taken with the F814W filter.

A visual inspection of the images reveals many background galaxies. Since 60 of these show a point-like nucleus, we used our ePSF-fitting procedure to measure positions for them also. We transformed the position of each galaxy into the reference frame using the same kind of transformations that we used for the stars. With _N_ observations for each galaxy at each epoch, we calculated the random error for each galaxy as \(\sigma/\sqrt{N-1}\), where _\sigma_ comes from the agreement among the independent measurements. We also added in quadrature the error in the transformation, based on the residuals of the stars used to compute the transformation. As expected, the errors in multiple measurements of galaxy positions are several times as large as typical errors in star positions, and depend strongly on the galaxy morphology. Nevertheless, we found that the result depended almost completely on the ~15 galaxies whose errors are smallest.

Figure II shows the vector point diagram for the objects measured. Dots show the proper motions of the stars, and filled circles those of the galaxies. To avoid confusion, we show error bars for only the best 15 galaxies, each of which has an error <1 mas yr^{-1} in each coordinate. The galaxies that have larger errors agree with the mean but contribute almost no weight to it. Cluster members form a tight clump below the middle of Fig. II within the 0.75-mas yr^{-1} circle that we have adopted to define cluster membership. (This radius was chosen as a compromise between losing cluster members and including field stars that have a motion close to that of the cluster stars.) The separation between field stars and cluster stars is well defined.

The zero point of the figure was placed at the weighted mean motion of the galaxies, and is marked with dotted lines and with the error bars of the zero point. This zero-point determination is the dominant source of uncertainty in the absolute motion of NGC 6791.

With this zero point we find for NGC 6791 an absolute proper motion, in the J2000.0 system, of

\[
(\mu_\alpha \cos \delta, \mu_\delta) = (-0.57, -2.45) \pm (0.13, 0.12) \text{ mas yr}^{-1}
\]

Table 1. Data set used in this work.

| Epoch(date) | EXPTIME | FILT | GO  |
|-------------|---------|------|-----|
| I (17July2003) | 2×1142s+4×1185s | F814W | 9815 |
| II (13July2005) | 2×1206s+4×1264s | F814W | 10471 |
These values represent a considerable change from the
\((\mu_\alpha \cos \delta, \mu_\delta) = (-0.10, -0.70) \pm (0.90, 0.80) \text{ mas yr}^{-1}\)
used by Carraro et al. (2006). In Galactic coordinates
(Fig. 2) our motions correspond to \((\mu_\ell \cos b, \mu_b) =
(-2.45, -0.56) \pm (0.13, 0.13) \text{ mas yr}^{-1}\). At the compressed
scale of Fig. 2 many more field stars can be seen. The dis-
tribution of their proper motions is clearly elongated; this
is the consequence of streaming effects due to differential
rotation, for stars at different distances along the line of
sight.

3. Radial velocity
For the third component of the motion of NGC 6791
(along the line of sight), we used the results in two re-
cently published papers. In the first of these Gratton et al.
(2006) obtained high-resolution spectra of four red-clump
stars; combining their measured radial velocities gives a
mean value of \(V_{\text{rad}} = -47.2 \pm 1.5 \text{ km s}^{-1}\). In the second
paper, Carraro et al. (2006) derived a mean radial velocity
\(V_{\text{rad}} = -47.1 \pm 0.8 \text{ km s}^{-1}\) from the spectra of 15 proba-
ble cluster members. We adopted as the radial velocity of
NGC 6791 the weighted mean of the two measurements,
\(V_{\text{rad}} = -47.1 \pm 0.7 \text{ km s}^{-1}\).

Table 2. Input conditions for orbit calculation. Distances
are in kpc, and velocities in km s\(^{-1}\).

| d    | U    | V    | W (= Z) | Π | Θ     |
|------|------|------|---------|----|--------|
| 3.6  | 34 ± 4| -51 ± 2| -11 ± 2 | 40 ± 4 | 168 ± 3 |
| 4.0  | 38 ± 4| -52 ± 2| -12 ± 2 | 43 ± 4 | 167 ± 3 |
| 4.4  | 42 ± 5| -54 ± 2| -13 ± 2 | 47 ± 4 | 165 ± 3 |

4. Calculation of the orbit
Our absolute proper motion for NGC 6791, along with
the radial velocity and an assumed distance of 4 kpc ±
10\% (King et al. 2005), allows us to derive its three veloc-
ity components and calculate its Galactic orbit. The orbit
should allow us to study the dynamical history of the clus-
ter, and to assess the possible impact of the motion on its
internal dynamics, its mass function, and its origin.

The integration of an orbit requires adopting a model
of the potential of the Milky Way. We chose that of Allen
& Santillan (1991), which assumes a Ga]actocentric dis-
tance and rotation velocity for the Sun of \(R_0 = 8.5 \text{ kpc}
and \(\Theta_0 = 220 \text{ km s}^{-1}\), and takes densities in bulge, disk,
and halo components whose combined gravitational force
fits a rotation curve that agrees with observation. Besides
being time-independent, their potential is axisymmetric,
fully analytic, and mathematically very simple. It has al-
ready been used to derive the Galactic orbits of open clus-
ters (Carraro & Chiosi 1994, Carraro et al. 2006) and disk
and halo globular clusters (Odenkirchen & Brosche 1992,
Milone et al. 2006). The potential is time-independent—
clearly a crude approximation, because a significant varia-
tion of the Galactic potential is expected over the lifetime
of this cluster. Nevertheless, it is reasonable to believe that

Table 3. Orbit parameters, for the three different dis-
cances. Units: \(d \text{[kpc]}\), \(L_z \text{[kpc km s}^{-1}]\), \(E_{\text{tot}} \times 10^5 \text{[km}^2\text{s}^{-2}]\),
\(P \text{[Myr]}\), \(R_0 \text{[kpc]}\), \(R_p \text{[kpc]}\), \(z_{\text{max}} \text{[kpc]}\), \(e \text{[pure number]}\).

| d    | \(L_z\) | \(E_{\text{tot}}\) | \(P\) | \(R_0\) | \(R_p\) | \(z_{\text{max}}\) | \(e\) |
|------|--------|-----------------|------|--------|--------|----------------|------|
| 3.6  | 1101   | -11085          | 120  | 9.50   | 3.32   | 0.86           | 0.48 |
| 4.0  | 1060   | -11174          | 130  | 9.83   | 3.09   | 0.98           | 0.52 |
| 4.4  | 994    | -11232          | 141  | 10.12  | 2.81   | 1.00           | 0.56 |
the real Galactic potential has not changed much in the last few Gyr, so that the parameters that we derive for the present-day orbit of NGC 6791, such as the apo- and perigalactic distances, can be considered to be reasonable estimates.

The initial conditions are given in Table 2 for three different heliocentric distances. The integration routine is a modified second-order Bulirsch-Stoer integrator (Press et al. 1992). The orbits were integrated back in time for 1 Gyr, and are shown in Fig. 3 both in the $xy$ plane (left panels), and in the meridional plane (right panels). The orbital parameters are given in Table 3, where successive columns give the assumed heliocentric distance of the cluster ($d$), the $z$-component of the angular momentum ($L_z$), the total energy ($E_{\text{tot}}$), the orbital period ($P$), the apocenter ($R_a$) and pericenter ($R_p$) of the orbit, the maximum vertical distance the cluster reaches ($z_{\text{max}}$) and the eccentricity ($e$), defined as $(R_a - R_p)/(R_a + R_p)$.

Our newly derived orbital parameters are based on a more accurate space velocity, and are thus an improvement over previous orbits. In particular, this is the first time that proper motions referred to extragalactic objects have been used in studying the orbit of NGC 6791.

![Orbits calculated back in time for 1 Gyr, for assumed cluster distances of 3.6, 4.0, and 4.4 kpc.](image)

The ±10% range in distance does not change the shape of the orbit significantly, nor do the orbital parameters change greatly. The orbit is of boxy type. The eccentricity is significantly higher than is typical for an open cluster. As the assumed heliocentric distance increases (from top to bottom panel) the cluster tends to have a longer period, to reach greater heights above the Galactic plane, and to show a larger epicyclic amplitude, dipping closer to the Galactic Center. The cluster never moves very far from the Sun toward the anticenter, while in the other direction it reaches rather small Galactocentric distances.

Over one radial period it crosses the Galactic plane three times. (This is quite clear in the bottom panel on the right of Fig. 3 for other assumed distances it is less obvious, but equally true.)

Note that the new tangential motion derived in this paper leads to some sizable differences from previous results. The most important change is in the apogalactic distance $R_\alpha$, which in the present paper is strikingly smaller than previous values. This weakens significantly the likelihood of an extragalactic origin for NGC 6791, as proposed by Carraro et al. (2006).

In its life in relatively dense regions of the Milky Way, NGC 6791 has had a difficult time dynamically. In each orbital period of $\sim 130$ Myr it has endured a rapid Galactocentric passage at $R \sim 3$ kpc, a disk crossing at $R \sim 9$ kpc, and two more rapid crossings through the denser part of the disk at $R \sim 5$ kpc—all four of them producing tidal shocks. The survival of the cluster till the present era is probably due only to its high density and large mass. The mass has been decreasing with time, however. Internal equipartition keeps the lower-mass stars preferentially in the outer parts, and the tidal buffeting has detached much of low-mass population, leading to the flat mass function noted by King et al. (2005).

Another application of knowledge of the orbit of NGC 6791 is to the question of the origin of a super-metal-rich cluster that is $\sim 8$ kpc from the Galactic center. Grenon (1999) has suggested that the stars in the solar neighborhood that have comparably high metallicity could have originated in the Galactic bulge and then been perturbed by the central bar into orbits that bring them out to here. Our orbit suggests that a similar dynamical history might apply to NGC 6791.

5. Summary and Conclusions

By identifying and measuring galaxies with point-like centers, we have been able to measure an absolute proper motion for NGC 6791. With the known radial velocity and the distance, we have computed the Galactic orbit of the cluster. Uncertainties in the orbit are due mainly to the inaccuracy of the distance, but partly to the unknown gravitational influence that the central bar of the Galaxy may have had on the orbit of the cluster. It is quite plausible that the high metallicity of NGC 6791 is associated with an origin in the inner region of the Galaxy.

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