Conformable derivative applied to experimental Newton’s law of cooling

J. Rosales-García, J. A. Andrade-Lucio, and O. Shulika
División de Ingenierías Campus Irapuato-Salamanca, Universidad de Guanajuato,
Carretera Salamanca-Valle de Santiago, km. 3.5+1.8 Comunidad de Palo Blanco, Salamanca, Guanajuato, México.
e-mail: rosales@ugto.mx

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It has been proved that the integer order differential equation does not represent the real behaviour of nature for Newton’s law of cooling. Then, we solve Newton’s cooling law using the conformable derivative; as a result, we obtain the Kohlrausch stretched exponential function.

2. Conformable Newton law of cooling

Let \( f : [0, \infty) \rightarrow \mathbb{R} \) be a function. Then the \( \gamma \)-th order conformable derivative of \( f \) is defined by [12] as

\[
T_\gamma(f)(t) = \frac{d^\gamma f(t)}{dt^\gamma} = f_\gamma(t) = \lim_{\epsilon \to 0} \frac{f(t + \epsilon\gamma^\gamma) - f(t)}{\epsilon},
\]

for all \( t > 0 \) and \( 0 < \gamma \leq 1 \). This expression is not the only possible [22,23].

When \( \gamma = 1 \) from (1), we obtain

\[
\frac{d}{dt}f(t) = \lim_{\epsilon \to 0} \frac{f(t + \epsilon) - f(t)}{\epsilon}.
\]

The CD is easily computed, but it is not conformable at \( \gamma = 0 \). The differentiation properties of this conformable derivative are analogous to the properties of the ordinary derivative, and the relation

\[
T_\gamma(f)(t) = \frac{d^\gamma f}{dt^\gamma} = t^{1-\gamma} f'(t),
\]

is fulfilled [12]. On the other hand, Newton’s law of cooling says that under certain circumstances the rate of change of temperature of a body is proportional to the difference between the temperature of the body and that of the surrounding medium.
environment [1], as long as the temperature difference is not very large. The above is described by the equation
\[ \frac{dT}{dt} = -k(T - T_m), \quad T(0) = T_0, \] (3)
where \( T(0) = T_0 \) is the initial temperature of the body, \( T_m \) is the temperature of the medium considered constant, and \( k \) is the convection coefficient. The coefficient \( k \) is measured in inverse unity of time, \( s^{-1} \). The Eq. (3) has an analytical solution in the form
\[ T(t) = T_m + (T_0 - T_m)e^{-kt}. \] (4)

In the previous work [21], to maintain the homogeneity of the fractional differential equation for Newton’s law of cooling, the ordinary derivative was changed by a fractional one, in the following way,
\[ \frac{d}{dt} = k^{1-\gamma} \frac{d^\gamma}{dt^\gamma}, \quad 0 < \gamma \leq 1, \] (5)
substituting in (3), we have the corresponding fractional differential equation of order \( \gamma \),
\[ d^\gamma T = -k^\gamma (T - T_m), \quad T(0) = T_0, \quad 0 < \gamma \leq 1. \] (6)

Solutions of this fractional equation using other fractional derivatives have been reported in [24].

However, with these fractional derivatives, it is not easy to calculate the convective coefficient \( k \). Due to this, in [21] was applied the conformable derivative [12]. For this, we take into account the expression (2), we have
\[ \frac{d}{dt} = k^{1-\gamma} \frac{d^\gamma}{dt^\gamma} = k^{1-\gamma} \frac{d}{dt^\gamma} \] (7)
Substituting this time-fractional conformable transform in (6), we obtain an ordinary differential equation
\[ \frac{dT}{dt} = -k^\gamma t^{\gamma-1}(T - T_m), \quad T(0) = T_0, \quad 0 < \gamma \leq 1. \] (8)
This equation has the particular solution [21],
\[ T(t; \gamma) = T_m + (T_0 - T_m)e^{-(k^\gamma/\gamma)^t}, \quad 0 < \gamma \leq 1. \] (9)

So, in the case of CD the Newton’s law of cooling has the stretched exponential function (or Kohlrausch function) of order \( 0 < \gamma \leq 1 \) as a solution. This function has many applications in science: it fits many relaxation processes in disordered and molecular systems [25-27], this is due to the free parameter \( \gamma \) in the exponential function. For a given time \( \tau \) and temperature \( T_1 \) we can compute, in a closed form, the conformable convective coefficient
\[ k(\gamma) = \left[ \frac{\gamma}{\tau^{\gamma}} \ln \left( \frac{T_0 - T_m}{T_1 - T_m} \right) \right]^{1/\gamma}, \quad 0 < \gamma \leq 1, \] (10)
unlike the results obtained in [24] using the Riemann-Liouville and Caputo fractional derivatives, where does not possible to obtain.

### 3. Data acquisition

Cooling experiments were performed using 300 ml of hot water in an isolated ceramic cup. The temperature was measured using a data acquisition system based on two DS18B20 temperature digital sensors attached to the Arduino board (AB), model UNO-R3. One of the sensors was kept fully submerged in the center of hot water volume. The other one was placed outside to monitor the ambient temperature.

The waterproof sensor DS18B20 provides data in digital form, with an accuracy of ±0.5°C and a maximum resolution of 0.25 × 10⁻² °C without the need for any analog-to-digital conversion or calibration. The sensor comes with three wires: the red one is connected to 5 V, the black one is connected to ground, and the yellow one is connected to a digital input pin of the AB (pin 10 in our case). For these sensors, a pull-up resistor \( R = 4.7 \, \text{k}\Omega \) is needed between the digital input pin and the 5 V source. The connection diagram for the AB (created with Fritzing) is shown as an inset in Fig. 1.

The communication between the sensors and the AB is implemented in three steps:

1. Send the initialization command \texttt{reset()}.  
2. Call to all the \texttt{OneWire} devices to identify themselves in the AB; in our case, we have two temperature sensors.  
3. Read from or write to a \texttt{OneWire} device.

The Arduino-script we used has two parts, the setup subroutine and the loop subroutine using the \texttt{OneWire} library into the Arduino Integrated Development Environment (IDE).

![Figure 1](image)

**Figure 1.** Experimental data fitted by integer order model Eq. (4), with the best value of \( k = 0.02438 \) 1/min.
4. Data Processing and Results

The temperatures are registered with an average time step $\Delta t = 28.45$ ms for almost 3.5 hours. The measured average ambient temperature is $T_m = 20.9147^\circ C$. Then these data have been fitted by both integer (4) and conformable (9) models. The $\chi^2$-criterion has been applied when fitting experimental data. Results in Fig. 2(a) shows a drastic divergence of integer model from experimental data. However, the conformable model provides an ideal fitting, as we can see from the quantitative analysis of residuals in Fig. 2(b). The conformal model provides results within experimental accuracy in the whole range of measurements, and thus, they are indistinguishable from the measured data.

5. Conclusions

The main aim of this work was to compare the theoretical results obtained in [21] with the real data obtained experimentally. We conclude that even in a very simple case, as is Newton’s law of cooling, the integer order differential equation cannot correctly predict the behavior of the system. In this work, it was shown that the solution (9) of the conformable differential equation (8) fits very well with the experimental data when $\gamma = 0.77269$ and $k = 0.018765 \text{ l/min}$. Unlike other fractional formalisms [24], the conformable derivative has an analytical solution very similar to the ordinary case, and fits very well with the experimental data; this is due to the free parameter $0 < \gamma \leq 1$. This is interesting since the conformable derivative contains the same differentiation rules as the ordinary derivative, which does not contain the fractional derivatives known until today.

We hope that the way of analyzing the fractional differential equations using the conformable fractional transform (7) will be helpful in solving fractional equations that represent more complex systems.

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