THE SPACE OF SOLSOLITONS IN LOW DIMENSIONS

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Abstract. Up to now, the only known examples of homogeneous nontrivial Ricci soliton metrics are the so called solsolitons, i.e. certain left invariant metrics on simple connected solvable Lie groups. In this paper, we describe the moduli space of solsolitons of dimension \( \leq 6 \) up to isomorphism and scaling. We start with the already known classification of nilsolitons and, following the characterization given by Lauret in [10], we describe the subspace of solsolitons associated to a given nilsoliton, as the quotient of a Grassmannian by a finite group.

1. Introduction

A complete Riemannian metric \( g \) on a differentiable manifold \( M \) is said to be a Ricci soliton if its Ricci tensor satisfies
\[
\text{Rc}(g) = cg + LX g \quad \text{for some } c \in \mathbb{R}, \ X \in \chi(M) \text{ complete,}
\]
where \( L_X \) denotes the Lie derivative in the direction of the vector field \( X \). Ricci soliton metrics came up in the study of the Ricci flow since they are the fixed points of the flow up to isometry and scaling. In the homogeneous case, all known nontrivial examples (i.e. not the product of an Einstein homogeneous manifold with a euclidean space) are isometric to a left-invariant metric \( g \) on a simply connected solvable Lie group \( G \). Moreover, by identifying \( g \) with an inner product on the Lie algebra \( g \) of \( G \), one has
\[
\text{Ric}(g) = ci + D \quad \text{for some } c \in \mathbb{R}, \ D \in \text{Der}(g),
\]
where \( \text{Ric}(g) \) denotes the Ricci operator of \( g \). These metrics have been deeply studied, specially in the nilpotent case, where they are called nilsolitons (see for example [9], [12], [14], [15], [16]). In this case, Lauret proved many nice properties: they are the nilpotent parts of Einstein solvmanifolds, they are unique up to isometry and scaling and they are the critical points of certain natural variational problem (see the survey [11]).

If \( G \) is solvable, metrics for which (2) holds are called solsolitons and have been recently characterized by Lauret in [10]. It is proved there that any solsoliton is the extension of a nilsoliton by an abelian subalgebra of symmetric derivations of its metric Lie algebra. They are also unique in the sense that, among all left invariant metrics on a given simply connected solvable Lie group, there is at most one solsoliton up to isometry and scaling (see [10]).

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Table 1. Nilsolitons of dimension \( \leq 3 \)

| nilsoliton metric | eigenvalue type | rank |
|-------------------|----------------|------|
| (0)               | (1; 1)         | 1    |
| (0, 0)            | (1; 2)         | 2    |
| (0, 0, 0)         | (1; 3)         | 3    |
| (0, 0, 12)        | (1 < 2; 2, 1)  | 2    |

The aim of this paper is to approach the classification of solsolitons in low dimensions. We first show in Section 3 that, given an \( n \)-dimensional nilsoliton \( N \), the space of solsolitons which are extensions of \( N \) of dimension \( r + n \) is parameterized, up to isometry and scaling, by

\[
\text{Gr}_r(a)/W.
\]

Here \( a \) is a maximal abelian subalgebra of symmetric derivations of the Lie algebra of \( N \), \( \text{Gr}_r(a) \) denotes the Grassmannian of \( r \)-dimensional subspaces of \( a \) and \( W \) is a finite group. Thus for each \( 0 \leq r \leq \text{rank}(N) := \dim(a) \), each nilsoliton \( N \) provides a family of \( (r+n) \)-dimensional solsolitons depending on \( r(\text{rank}(N) - r) \) parameters. The number \( \text{rank}(N) \), called the rank of \( N \), is therefore the crucial datum to know on a nilsoliton if one is interested in classifying solsolitons.

In Section 4, we compute the rank of all nilsolitons of dimension \( \leq 6 \) (see Tables 1-5). In most cases, we give in addition an explicit description of the coset \( \text{Gr}_r(a)/W \) as a real semialgebraic set (see [7] for a similar approach in the case of nilsolitons attached to graphs). All this allows us to present quite a detailed picture of the moduli space \( \text{Sol}(m) \) of \( m \)-dimensional solsolitons for any \( m \leq 6 \) (see Section 5). Finally, in Section 6, we make some remarks on negatively curved (sectional and Ricci) solsolitons.

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2. Preliminaries

Let \( S \) be a solvmanifold, that is, a simply connected solvable Lie group endowed with a left invariant Riemannian metric. We will identify \( S \) with its metric Lie algebra \( (s, \langle \cdot, \cdot \rangle) \), where \( s \) is the Lie algebra of \( S \) and \( \langle \cdot, \cdot \rangle \) denotes the inner product on \( s \) determined by the metric. If \( n \) is the nilradical of \( s \) (i.e. the maximal nilpotent ideal), consider the orthogonal decomposition

\[
s = a \oplus n.
\]

\( S \) is called standard if \( [a, a] = 0 \).
Table 2. 4-dimensional nilsolitons

| N° | nilsoliton metric $\eta_i$ | eigenvalue type | rank |
|----|-----------------------------|-----------------|------|
| 1  | (0, 0, 12, 13)              | (1 < 2 < 3 < 4; 1,..., 1) | 2    |
| 2  | (0, 0, 0, 12)               | (2 < 3 < 4; 2, 1, 1)    | 3    |
| 3  | (0, 0, 0, 0)                | (1; 4)             | 4    |

**Definition 2.1.** A left-invariant metric $g$ on $S$ is said to be a solsoliton if

$$\text{(3)} \quad \text{Ric}(g) = cI + D,$$

where $\text{Ric}(g)$ denotes the Ricci operator of $g$.

Hence solsolitons are natural generalizations of Einstein metrics (i.e., when $D = 0$). Einstein solvmanifolds have been studied by Heber who gave in [6] many structural results. In particular, if $S$ is an Einstein solvmanifold, he showed that there exists an element $H \in \mathfrak{a}$ such that the eigenvalues of $\text{ad} \, H |_{\mathfrak{n}}$ are positive integers without common divisors, say $k_1 < \cdots < k_r$ with multiplicities $d_1, \ldots, d_r$. The eigenvalue type of $S$ is defined as the tuple

$$(k; d) = (k_1 < \cdots < k_r; d_1, \ldots, d_r).$$

We have that $\text{ad} \, H |_{\mathfrak{n}}$ is a multiple of $D$ in equation (3).

When the group $S$ is nilpotent, metrics for which (3) holds are called nilsolitons. They have been deeply studied by Lauret (see [11]) who has proved, among others properties, that:

- a given nilpotent Lie group $N$ can admit at most one nilsoliton up to isometry and scaling among all its left-invariant metrics,
- nilsolitons are precisely the metric nilradicals of Einstein solvmanifolds,
- any Ricci soliton left invariant metric is a nilsoliton.

By this uniqueness result we can associate to a nilsoliton the eigenvalue type of its uniquely defined rank-one Einstein solvable extension.

We give in what follows a summary of Lauret’s results we are using in our classification on solsolitons (see [10]). Note that we have slightly changed the notation in [10] to be consistent with the one that is useful in this context. For any metric nilpotent Lie algebra $\langle \mathfrak{n}, \langle \cdot, \cdot \rangle \rangle$, let $\text{sym}(\mathfrak{n})$ denote the space of symmetric transformations of $\mathfrak{n}$ with respect to $\langle \cdot, \cdot \rangle$.

It is proved in [10] that any solsoliton can be constructed, up to isometry and scaling, from a nilsoliton $\langle \mathfrak{n}, \langle \cdot, \cdot \rangle \rangle$ and an abelian subalgebra of symmetric derivations of $\mathfrak{n}$ in the following way (see [10] Corollary 4.10).

**Proposition 2.2.** [10] Proposition 4.3] Let $\langle \mathfrak{n}, \langle \cdot, \cdot \rangle \rangle_1$ be a nilsoliton with Ricci operator $\text{Ric}_1 = cI + D_1$, $c < 0$, $D_1 \in \text{Der}(\mathfrak{n})$, and consider a any abelian Lie algebra of symmetric derivations of $\langle \mathfrak{n}, \langle \cdot, \cdot \rangle \rangle_1$. Then the solvmanifold $S$ with Lie
Proposition 2.3. [9, Proposition 5.3] Let \((\mathfrak{n}, \langle \cdot, \cdot \rangle)\) be a nilsoliton and let \(S\) and \(S'\) be two solsolitons constructed as in Proposition 2.2 for abelian subalgebras \(\mathfrak{a}, \mathfrak{a}' \subset \text{Der}(\mathfrak{n}) \cap \text{sym}(\mathfrak{n})\), respectively. Then \(S\) is isometric to \(S'\) if and only if there exists \(h \in \text{Aut}(\mathfrak{n}) \cap O(\langle \cdot, \cdot \rangle)\) such that \(\mathfrak{a}' = h \mathfrak{a} h^{-1}\).

Let \(\mu : \mathbb{R}^n \times \mathbb{R}^n \to \mathbb{R}^n\) be a skew-symmetric bilinear form. If \(\mu\) satisfies the Jacobi identity then it defines a Lie algebra \((\mathbb{R}^n, \mu)\). We will denote by \(\mathcal{N} = \{\mu : \mathbb{R}^n \times \mathbb{R}^n \to \mathbb{R}^n : \mu\text{ skew-symmetric, bilinear, nilpotent and satisfies Jacobi}\}\). \(\mathcal{N}\) is often called the variety of nilpotent Lie algebras (of dimension \(n\)). If we also fix the canonical basis on \(\mathbb{R}^n\), \(\{X_1, \ldots, X_n\}\), and the inner product \(\langle \cdot, \cdot \rangle\) that makes this basis orthonormal then each \(\mu \in \mathcal{N}\) define a metric nilpotent Lie algebra \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\). We will say then that \(\mu \in \mathcal{N}\) is a nilsoliton if \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\) is so.

Nilsolitons metrics are classified up to dimension 6 (see [9] and [20]). They are listed in Tables 1 through 5. In these tables, an element \(\mu \in \mathcal{N}\) is represented as a

| \(N^o\) | nilsoliton metric \(\lambda_i\) | eigenvalue type | rank |
|---|---|---|---|
| 1 | \((0,0,3,12,4,13,3,14)\) | \((2 < 9 < 11 < 13 < 15;1,...,1)\) | 2 |
| 2 | \((0,0,3^{1/2},12,3^{1/2},13,2^{1/2},14 + 2^{1/2},123)\) | \((1 < 2 < 3 < 4 < 5;1,...,1)\) | 1 |
| 3 | \((0,0,0,2^{1/2},12,23 + 2^{1/2},14)\) | \((3 < 4 < 6 < 7 < 10;1,...1)\) | 2 |
| 4 | \((0,0,0,0,12 + 34)\) | \((1 < 2;4,1)\) | 3 |
| 5 | \((0,0,2,12,3^{1/2},13,3^{1/2},23)\) | \((1 < 2 < 3;2,1,2)\) | 2 |
| 6 | \((0,0,0,12,13)\) | \((2 < 3 < 5;1,2,2)\) | 3 |
| 7 | \((0,0,0,0,12)\) | \((2 < 3 < 4;2,2,1)\) | 4 |
| 8 | \((0,0,0,12,14)\) | \((1 < 2 < 3 < 4;1,1,2,1)\) | 3 |
| 9 | \((0,0,0,0,0)\) | \((1;5)\) | 5 |

Table 3. 5-dimensional nilsolitons
| $N^*$ | Nilsoliton metric : $\mu_i$ | eigenvalue type | rank |
|------|------------------|-----------------|-----|
| 1    | $\begin{pmatrix} 0, 0, (13)^\frac{1}{2} 12, 4 13, (12)^\frac{1}{2} 14 + 2 23, \\
                     \quad (12)^\frac{1}{2} 34 + (13)^\frac{1}{2} 52 \end{pmatrix}$ | $(1 < 2 < 3 < 4 < 5 < 7; 1, ..., 1)$ | 1 |
| 2    | $\begin{pmatrix} 0, 0, 12, (\frac{1}{4})^\frac{1}{2} 13, 14, 34 + 52 \end{pmatrix}$ | $(1 < 3 < 4 < 5 < 6 < 9; 1, ..., 1)$ | 2 |
| 3    | $\begin{pmatrix} 0, 0, 2 12, 6^\frac{1}{2} 13, 6^\frac{1}{2} 14, 2 15 \end{pmatrix}$ | $(1 < 9 < 10 < 11 < 12 < 13; 1, ..., 1)$ | 2 |
| 4    | $\begin{pmatrix} 0, 0, (22)^\frac{1}{2} 12, 6 13, (22)^\frac{1}{2} 14 + (30)^\frac{1}{2} 23, \\
                     \quad 5 24 + (30)^\frac{1}{2} 15 \end{pmatrix}$ | $(1 < 2 < 3 < 4 < 5 < 6; 1, ..., 1)$ | 1 |
| 5    | $\begin{pmatrix} 0, 0, 7^\frac{1}{2} 12, (\frac{1}{\sqrt{3}})^\frac{1}{2} 13, 3 14, \\
                     \quad (\frac{1}{\sqrt{3}})^\frac{1}{2} 23 + 2 15 \end{pmatrix}$ | $(1 < 3 < 4 < 5 < 6 < 7; 1, ..., 1)$ | 1 |
| 6    | $\begin{pmatrix} 0, 0, 12, 13, 23, 14 \end{pmatrix}$ | $(1 < 2 < 3 < 4 < 5; 1, 1, 1, 1, 2)$ | 2 |
| 7    | $\begin{pmatrix} 0, 0, 2 12, 5^\frac{1}{2} 13, 5^\frac{1}{2} 23, 2 14 + 2 25 \end{pmatrix}$ | $(1 < 2 < 3 < 4; 2, 1, 2, 1)$ | 2 |
| 8    | $\begin{pmatrix} 0, 0, 2 12, 5^\frac{1}{2} 13, 5^\frac{1}{2} 23, 2 14 + 2 25 \end{pmatrix}$ | $(1 < 2 < 3 < 4; 2, 1, 2, 1)$ | 1 |
| 9    | $\begin{pmatrix} 0, 0, (\frac{1}{2})^\frac{1}{2} 12, 14 - 23, (\frac{1}{2})^\frac{1}{2} 15 + 34 \end{pmatrix}$ | $(6 < 11 < 12 < 17 < 23 < 29; 1, ..., 1)$ | 2 |
| 10   | $\begin{pmatrix} 0, 0, 12, (\frac{1}{2})^\frac{1}{2} 14, 15 + 23 \end{pmatrix}$ | $(4 < 9 < 12 < 13 < 17 < 21; 1, ..., 1)$ | 2 |
| 11   | $\begin{pmatrix} 0, 0, (\frac{1}{2})^\frac{1}{2} 12, (\frac{2}{\sqrt{3}})^\frac{1}{2} 12, (\frac{2}{\sqrt{3}})^\frac{1}{2} 14 \\
                     \quad - (\frac{1}{\sqrt{3}})^\frac{1}{2} 13, (\frac{2}{\sqrt{3}})^\frac{1}{2} 15 + (\frac{1}{\sqrt{3}})^\frac{1}{2} 24 \end{pmatrix}$ | $(1 < 2 < 3 < 4 < 5; 1, 1, 2, 1, 1)$ | 1 |
| 12   | $\begin{pmatrix} 0, 0, 3^\frac{1}{2} 12, 3^\frac{1}{2} 14, 2^\frac{1}{2} 15 + 2^\frac{1}{2} 24 \end{pmatrix}$ | $(3 < 6 < 9 < 11 < 12 < 15; 1, ..., 1)$ | 2 |
| 13   | $\begin{pmatrix} 0, 0, 3^\frac{1}{2} 12, 2 14, 3^\frac{1}{2} 15 \end{pmatrix}$ | $(2 < 9 < 11 < 12 < 13 < 15; 1, ..., 1)$ | 3 |

Table 4. 4 and 5-step 6-dimensional nilsolitons

vector of $n$ coordinates in such a way that

$$ k^{th}-\text{coordinate} = cij + dri \iff \mu(X_i, X_j) = cX_k, \mu(X_r, X_l) = dX_k. $$

Hence, for example, $(0, 0, 0, 2^\frac{1}{2} 12, 23 + 2^\frac{1}{2} 14)$ represents the 5-dimensional Lie algebra with Lie bracket given by

$$ \mu(X_1, X_2) = 2^\frac{1}{2} X_4, \quad \mu(X_2, X_3) = X_5, \quad \mu(X_1, X_4) = 2^\frac{1}{2} X_5. $$

We will denote by $d(a_1, \ldots, a_n)$ the diagonal matrix with entries $a_1, \ldots, a_n$, and more in general

$$ d(M_1, \ldots, M_r) := \begin{bmatrix} M_1 & \cdots & M_r \end{bmatrix}, $$

for matrices $M_i \in \mathfrak{gl}(n_i)$ such that $n_1 + \cdots + n_r = n$.

3. General facts on the classification of solsolitons

Let $(\mathbb{R}^n, \mu)$ be a nilpotent Lie algebra. We fix the canonical inner product $\langle \cdot, \cdot \rangle$ on $\mathbb{R}^n$ and denote by $\mathfrak{so}(n)$, $\mathfrak{sym}(n)$ and $\mathfrak{O}(n)$ the subspaces of skew-symmetric,
symmetric linear maps on $\mathbb{R}^n$ and the orthogonal group relative to $\langle , \rangle$, respectively. Let $G_\mu$ denote the closed subgroup of automorphisms given by

$$G_\mu := \{ g \in \text{Aut}(\mu) : g^t \in \text{Aut}(\mu) \}. $$

Since $\text{Aut}(\mu)$ is algebraic, we have that $G_\mu$ is a real reductive group (see [19] Section 2.1) with Lie algebra

$$\mathfrak{g}_\mu = \{ A \in \text{Der}(\mu) : A^t \in \text{Der}(\mu) \},$$

and Cartan decompositions

$$G_\mu = K_\mu \exp(p_\mu), \quad \mathfrak{g}_\mu = \mathfrak{k}_\mu \oplus \mathfrak{p}_\mu,$$
where
\[ K_\mu := G_\mu \cap O(n) = \text{Aut}(\mu) \cap O(n), \quad \mathfrak{p}_\mu := \mathfrak{g}_\mu \cap \mathfrak{so}(n) = \text{Der}(\mu) \cap \mathfrak{so}(n), \]
\[ p_\mu := g_\mu \cap \text{sym}(n) = \text{Der}(\mu) \cap \text{sym}(n). \]

It is proved in [10, Lemma 4.7] that \( g_\mu \) is precisely the space of derivations of \((\mathbb{R}^n, \mu)\) which are normal (i.e. \([D, D'] = 0\)).

Let us assume from now on that \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\) is a nilsoliton. According to Proposition 2.3, the solsolitons associated with the nilsoliton \( \mu \) are parameterized by abelian subspaces of \( p_\mu \) up to conjugancy by \( K_\mu \).

Let \( a_\mu \) be a maximal abelian subalgebra of \( p_\mu \). Since \( G_\mu \) is real reductive, it is well known that any abelian subspace of \( p_\mu \) is \( K_\mu \)-conjugate to a subspace of \( a_\mu \) (see [19, 2.1.9]). Moreover, any two subspaces of \( a_\mu \) are \( K_\mu \)-conjugate if and only if they are conjugate by an element of the Weyl group \( W_\mu \) of \( G_\mu \) (see [5, Proposition 2.2, Ch.VII]), which is given by
\[ W_\mu := N_{K_\mu}(a_\mu)/Z_{K_\mu}(a_\mu), \]
where
\[ N_{K_\mu}(a_\mu) := \{ g \in K_\mu : \text{Ad}(g)a_\mu \subset a_\mu \}, \]
\[ Z_{K_\mu}(a_\mu) := \{ g \in K_\mu : \text{Ad}(g)A = A, \quad \forall A \in a_\mu \}. \]

We therefore obtain that the set of \((r + n)\)-dimensional solsolitons associated with the nilsoliton \( \mu \) is parameterized, up to isometry and scaling, by the coset
\[ (5) \quad \text{Gr}_r(a_\mu)/W_\mu, \]
where \( \text{Gr}_r(a_\mu) \) is the Grassmanian of \( r \)-dimensional subspaces of \( a_\mu \). We note that \( 0 \leq r \leq \text{rank}(\mu) \), where \( \text{rank}(\mu) := \dim a_\mu \) will be called the rank of the nilsoliton \( \mu \). As \( W_\mu \) is a finite group, the quotient \( \text{Gr}_r(a_\mu)/W_\mu \) depends on \( r(\text{rank}(\mu) - r) \) parameters.

The aim of this paper is to compute the rank for any nilsoliton of dimension \( \leq 6 \), and also to give in most cases an explicit description of the quotient in (5).

We finish this section by dealing with the cases that admit a simultaneous study. In the next section we will work out a case by case analysis of the remaining nilsolitons.

### 3.1. Abelian nilsolitons

A first case that can be studied in general is the abelian case, that is, \( n = (\mathbb{R}^n, \mu) \) with \( \mu = 0 \). We have that \( \text{Der}(0) = \mathfrak{gl}(n) \), and therefore a maximal abelian subspace of symmetric derivations can be chosen to be the set of diagonal matrices:
\[ p_0 = \text{sym}(n) \quad a_0 = \{ d(a_1, \ldots, a_n) : a_i \in \mathbb{R} \}, \quad \text{rank}(0) = n. \]

Since \( K_0 = O(n) \), the action by conjugation on diagonal matrices contains all the permutations of the entries. Hence, for \( 0 \leq r \leq n \), the space of \((r + n)\)-dimensional solsolitons is given by
\[ \text{Gr}_r(\mathbb{R}^n)/S_n, \]
where \( S_n \) is the permutation group of \( n \) elements. Note that the \((n + 1)\)-dimensional hyperbolic space \( \mathbb{R}H^{n+1} \) is an element of \( \text{Gr}_r(\mathbb{R}^n)/S_n = \mathbb{P}\mathbb{R}^n/S_n \) and is characterized as the only point that corresponds to an Einstein metric.
3.2. Maximal and minimal dimension solsolitons. If \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\) is a nilsoliton such that \(\text{rank}(\mu) = k\), then

\[
\text{Gr}_0(\mathfrak{a}_\mu) = \{\ast\}, \quad \text{Gr}_k(\mathfrak{a}_\mu) = \{\ast\}.
\]

In other words, there is only one solsoliton \(s = a \oplus \mathfrak{n}_\mu\) of dimension \((k+n)\) associated to the nilsoliton \(\mu\), which is always Einstein (see Proposition 2.2), and only one of dimension \(n\), the nilsoliton itself, which is never Einstein unless \(\mu = 0\).

3.3. Rank-one nilsolitons. Another case we can study separately is when \(\mathfrak{a}_\mu\) is one dimensional. Then

\[
\text{Gr}_1(\mathfrak{a}_\mu) = \{\ast\},
\]

that is, there is only one solsoliton associated to \(\mu\) which is always Einstein. For dimension \(\leq 6\) they are all listed in Table 6.

3.4. Multiplicity-free eigenvalue type nilsoliton. It is well known that skew-symmetric and symmetric derivations, as well as orthogonal automorphisms of \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\), all commute with Ric\(\mu\) (see [6, Lemma 2.2]). In this way, if \(\mu\) is an nilsoliton, say with Ric\(\mu = c \text{Id} + \mathcal{D}_\mu\), all the above mentioned operators also commute with \(\mathcal{D}_\mu\). Then if we fix \(\beta = \{X_1, \ldots, X_n\}\) a basis of eigenvectors of \(\mathcal{D}_\mu\) we get the following observation.

**Lemma 3.1.** Let \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\) be a nilsoliton with eigenvalue type \((k_1 < \cdots < k_r; n_1, \ldots, n_r)\). Then the elements in \(\mathfrak{p}_\mu\) and in \(\mathfrak{K}_\mu\) are diagonal block matrices of dimension \(n_1, \ldots, n_r\), as in [4].

**Definition 3.2.** An eigenvalue type of the form \((k_1 < \cdots < k_r; 1, \ldots, 1)\) will be called multiplicity-free.

If \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\) is a nilsoliton with multiplicity-free eigenvalue type, then the elements in \(\mathfrak{p}_\mu\) are diagonal matrices and \(\mathfrak{K}_\mu \subseteq \mathfrak{d}(\pm 1, \ldots, \pm 1)\). This implies that \(\mathfrak{K}_\mu\) acts trivially by conjugation on \(\mathfrak{a}_\mu\) and consequently, the solsolitons of dimension \(r + n\) are just parameterized by

\[
\text{Gr}_r(\mathfrak{a}_\mu), \quad 0 \leq r \leq \text{rank}(\mu).
\]

The multiplicity-free eigenvalue type nilsolitons of rank \(\geq 2\) are listed in Table 7 where the relevant information is given.

| metric | \(\mathfrak{a}_\mu\) | dim. |
|--------|-----------------|-----|
| \(\lambda_2\) | \(\mathbb{R}d(1, 2, 3, 4, 5)\) | 5   |
| \(\mu_1\) | \(\mathbb{R}d(1, 2, 3, 4, 5, 7)\) | 6   |
| \(\mu_4\) | \(\mathbb{R}d(1, 2, 3, 4, 5, 6)\) | 6   |
| \(\mu_5\) | \(\mathbb{R}d(1, 3, 4, 5, 6, 7)\) | 6   |
| \(\mu_8\) | \(\mathbb{R}d(1, 1, 2, 3, 4)\) | 6   |
| \(\mu_{11}\) | \(\mathbb{R}d(1, 2, 3, 4, 5)\) | 6   |
| \(\mu_{18}\) | \(\mathbb{R}d(1, 1, 2, 3, 3)\) | 6   |

Table 6. Rank-one nilsolitons of dimension \(\leq 6\).
In this section, we study the nilsolitons of dimension \(\leq 6\) that are not included in Tables 6 or 7. For each one of these nilsolitons \(\mu\), we give a set that parameterizes the space of \((n + 1)\)-dimensional nilsolitons with nilradical \(\mu\) (up to isometry and scaling) by explicitly describing \(\mathbb{P}a_\mu/W_\mu\) (see 5).

Note that the nilsolitons we are studying have rank at most 5 and therefore it follows from what was explained in Sections 3.2, 3.3, and 3.5 that this is the only case we need to consider for those of rank one, two or three. If the rank is greater than three, to have a complete parametrization of the set of nilsolitons associated to \(\mu\), it remains to consider \(\text{Gr}_k(a_\mu)/W_\mu\), but this is a problem that exceeds the scope of this paper. We hope to deal with it in a future work.

We will denote by \(n_i\), the Lie algebra \((\mathbb{R}^n, \mu, \langle \cdot, \cdot \rangle)\), where \(\langle \cdot, \cdot \rangle\) is the canonical inner product we have fixed, and by \(\sigma_{ij}\) the element of the symmetric group \(S_n\) that permutes \(X_i\) with \(X_j\) and fixes all other elements of the basis.

### Table 7. Multiplicity-free eigenvalue type nilsolitons of rank \(\geq 2\).

| metric | \(a_\mu\) | Einstein condition | rank | dim. |
|--------|------------|-------------------|------|------|
| \(\eta_1\) | \((a, b, a + b, 2a + b)\) | \(2a = b\) | 2 | 4 |
| \(\lambda_1\) | \((a, b, a + b, 2a + b, 3a + b)\) | \(9a = b\) | 2 | 5 |
| \(\lambda_3\) | \((a, b, 2a, a + b, 2a + b)\) | \(4a = b\) | 2 | 5 |
| \(\mu_2\) | \((a, b, a + b, 2a + b, 3a + b, 3a + 2b)\) | \(3a = b\) | 2 | 6 |
| \(\mu_3\) | \((a, b, a + b, 2a + b, 3a + b, 4a + b)\) | \(9a = b\) | 2 | 6 |
| \(\mu_9\) | \((a, b, 2a, a + b, 2a + b, 3a + b)\) | \(11a = 6b\) | 2 | 6 |
| \(\mu_{10}\) | \((a, b, 3a, a + b, 2a + b, 3a + b)\) | \(9a = 4b\) | 2 | 6 |
| \(\mu_{12}\) | \((a, 2a, b, 3a, a + b, 2a + b)\) | \(11a = 3b\) | 2 | 6 |
| \(\mu_{13}\) | \((a, b, c, a + b, 2a + b, 3a + b)\) | \(a/2 = b/9 = c/12\) | 3 | 6 |
| \(\mu_{14}\) | \((a, b, \frac{a+b}{2}, a + b, \frac{a+b}{2}, 2a + b)\) | \(2a = b\) | 2 | 6 |
| \(\mu_{19}\) | \((a, b, 2b, a + b, 2a + b, a + 2b)\) | \(6a = 5b\) | 2 | 6 |
| \(\mu_{21}\) | \((a, b, 2a, a + b, 3a, 2a + b)\) | \(5a = 3b\) | 2 | 6 |
| \(\mu_{22}\) | \((a, b, c, a + b, a + c, a + 2b)\) | \(a/6 = b/5 = c/9\) | 3 | 6 |
| \(\mu_{23}\) | \((a, b, c, a + b, a + c, 2a + b)\) | \(a/2 = b/5 = c/6\) | 3 | 6 |
4.1. Nilsolitons of dimension 3. This case has been considered in [14] but we will include it again, for completeness.

- \( h_3 \), the 3-dimensional Heisenberg Lie algebra, given by \((0, 0, 12)\) (see Table 1).

Therefore

\[
p_{h_3} = \left\{ \begin{bmatrix} a & c \\ c & b \\ a+b \end{bmatrix} : a, b, c \in \mathbb{R} \right\}, \quad K_{h_3} = \{ d(H, \det H) : H \in O(2) \},
\]

\[
a_{h_3} = \{ d(a, b, a+b) : a, b \in \mathbb{R} \}.
\]

\( W_{h_3} \) is the group generated by \( \{ \text{Id}, \sigma_{12} \} \).

The corresponding set \( \mathbb{P}a_{h_3}/W_{h_3} \) of 4-dimensional solsolitons is parameterized by:

\[
\mathcal{F}_1^4 = \left\{ (a, b) \in S^3 : |a| \leq b \right\}.
\]

where, in general, \( S^n \) is the \( n \)-dimensional sphere in \( \mathbb{R}^{n+1} \).

Einstein condition: \( a = b \).

This correspondence assigns to each \((a, b) \in \mathcal{F}_1^4\) the solsoliton \( s = \mathbb{R}\left[ \begin{bmatrix} a & b \\ b & a \end{bmatrix} \oplus h_3 \right], \)

where Lie bracket and inner product are respectively given by

\[
\begin{align*}
[X_1, X_2] &= X_3, \quad [X_0, X_1] = a X_1, \quad [X_0, X_2] = b X_2, \quad [X_0, X_3] = (a + b) X_3, \\
\langle X_i, X_j \rangle &= \delta_{ij}, \quad \text{for } 0 \leq i, j \leq 3, \quad \langle X_0, X_0 \rangle = \frac{4}{3} (a^2 + b^2 + ab).
\end{align*}
\]

4.2. Nilsolitons of dimension 4. (see Table 2)

- \( \eta_2 \simeq h_3 \oplus \mathbb{R}X_3 \) Derivations of \( \eta_2 \) have been studied in [8].

\[
p_{\eta_2} = \{ d(A, c, \text{tr} A) : A \in \text{sym}(2), c \in \mathbb{R} \}, \quad K_{\eta_2} = \{ d(H, \pm 1, \det H) : H \in O(2) \},
\]

\[
a_{\eta_2} = \{ d(a, b, c, a+b) : a, b, c \in \mathbb{R} \}.
\]

\( W_{\eta_2} \) is the group generated by \( \{ \text{Id}, \sigma_{12} \} \).

The corresponding 5-dimensional solsolitons are parameterized by:

\[
\mathcal{F}_2^5 = \left\{ (a, b, c) \in S^2 : \begin{array}{l}
c > 0 \\
c = 0 \\
|a| \leq b \end{array}, \quad a \leq b, \text{ or } \right\}.
\]

Einstein condition: \( a = b = 2/3c \).

The above correspondence is given by

\[
(a, b, c) \mapsto \mathbb{R}\left[ \begin{bmatrix} a & b \\ b & a \end{bmatrix} \oplus \eta_2 \right],
\]

where \( \eta_2 = (\mathbb{R}^4, \eta_2) \).

The Lie bracket and the inner product of \( s \) are explicitly given by:

\[
\begin{align*}
[X_1, X_2] &= X_4, \quad [X_0, X_1] = a X_1, \quad [X_0, X_2] = b X_2, \quad [X_0, X_3] = c X_3, \quad [X_0, X_4] = (a+b) X_4, \\
\langle X_i, X_j \rangle &= \delta_{ij}, \quad \text{for } 0 \leq i, j \leq 4, \quad \langle X_0, X_0 \rangle = \frac{4}{3} (a^2 + b^2 + ab + c^2).
\end{align*}
\]
4.3. Nilosolitons of dimension 5. (see Table 3)

- \( \lambda_4 \) (5-dimensional Heisenberg Lie algebra).

\[
p_{\lambda_4} = \left\{ \begin{pmatrix} a & a & b \\ B & c & \beta \\ a+b-c & a+b-c & \alpha \end{pmatrix} : a, b, c, \alpha, \beta \in \mathbb{R}, \quad B = \begin{bmatrix} u & u \\ v & -u \end{bmatrix} \right\},
\]

\[
a_{\lambda_4} = \{ \mathbf{d}(a, b, c, a + b - c, a + b) : a, b, c \in \mathbb{R} \},
\]

\[
K_{\lambda_4} = \left\{ \begin{pmatrix} [T(A)]_1 \\ [S(A)]_{-1} \end{pmatrix} : A \in \mathbb{U}(2) \right\},
\]

where if for any \( z = a + ib \in \mathbb{C} \) we denote by \( M(z) = \begin{pmatrix} a & b \\ -b & a \end{pmatrix} \) and \( \tilde{M}(z) = \begin{pmatrix} b & -a \\ a & b \end{pmatrix} \), then for any \( A = \begin{pmatrix} z_1 & z_2 \\ z_3 & z_4 \end{pmatrix} \in \mathbb{U}(2) \), \( T(A) \) and \( S(A) \) are the matrices in \( \mathfrak{gl}(4) \) given by

\[
T(A) = \begin{pmatrix} M(z_1) M(z_2) \\ M(z_3) M(z_4) \end{pmatrix}, \quad S(A) = \begin{pmatrix} \tilde{M}(z_1) \tilde{M}(z_2) \\ \tilde{M}(z_3) \tilde{M}(z_4) \end{pmatrix}.
\]

It is easy to see that by considering \( A_1 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \) and \( A_2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \) then the element in \( K_{\lambda_4} \) corresponding to \( T(A_1) \) permutes \( a, b \) with \( c, d \) and the one corresponding to \( T(A_2) \) permutes \( a \) with \( b \) and fixed \( c, d \). Therefore \( W_{\mu_4} \) is generated by \( \{ \text{Id}, \sigma_{12}, \sigma_{34}, \sigma_{13}\sigma_{24} \} \).

It can be seen that any derivation \( \mathbf{d}(a, b, c, a + b - c, a + b) \) is related under the action of \( W_{\mu_4} \) to a multiple of an element in

\[
\{ \mathbf{d}(a, b, c, a + b - c, a + b) : a > 0, \quad 2c \geq a + b \geq 0, \quad a \geq c \geq b \}.
\]

Therefore, the set \( P_{a_{\lambda_4}}/W_{\lambda_4} \) of corresponding 6-dimensional solsolitons is parameterized by:

\[
\mathfrak{F}_2 = \{ (a, b, c) \in S^2 : a > 0, \quad 2c \geq a + b \geq 0, \quad a \geq c \geq b \},
\]

where the correspondence is given by \( (a, b, c) \rightarrow \mathbf{d}(a, b, c, a + b - c, a + b) \).

Einstein condition: \( a = b = c \).

Remark 4.1. As \( \text{rank}(\lambda_4) = 3 \), the set \( \mathfrak{F}_2 \) also parameterizes \( \text{Gr}_2(a_{\lambda_4})/W_{\lambda_4} \), the space of 7-dimensional solsolitons with nilradical \( \lambda_4 \).

- \( \lambda_5 \) (free 2-step nilpotent Lie algebra with 3 generators).

\[
p_{\lambda_5} = \{ \mathbf{d}(\begin{pmatrix} a & 0 \\ a & 0 \end{pmatrix}, a + b, \begin{pmatrix} 2a+b & a \\ a+2b \end{pmatrix}) : a, b, a \in \mathbb{R} \}, \quad \mathfrak{t}_{\lambda_5} = \{ 0 \}.
\]

So we can choose

\[
a_{\lambda_5} = \{ \mathbf{d}(a, b, a + b, 2a + b, a + 2b) : a, b \in \mathbb{R} \}.
\]

It is easy to see that by acting with \( A = \mathbf{d}(\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, -1, \begin{pmatrix} 0 & -1 \\ 0 & 1 \end{pmatrix}) \in K_{\lambda_5} \), we can permute \( a \) with \( b \). The corresponding 6-dimensional solsolitons are parameterized by \( \mathfrak{F}_1 \) (see (2)).

Einstein condition: \( a = b \).

- \( \lambda_6 \).

\[
p_{\lambda_6} = \{ \mathbf{d}(a, \begin{pmatrix} b & \beta \\ \beta & c \end{pmatrix}, \begin{pmatrix} a+b & \beta \\ \beta & a+c \end{pmatrix}) : a, b, c, \beta \in \mathbb{R} \},
\]

\[
a_{\lambda_6} = \{ \mathbf{d}(a, b, c, a + b, a + c) : a, b, c \in \mathbb{R} \}, \quad K_{\lambda_6} = \{ \mathbf{d}(\varepsilon, H, \varepsilon H) : H \in \mathbb{O}(2), \varepsilon = \pm 1 \}.
\]

The corresponding 6-dimensional solsolitons are parameterized by: \( (c, b, a) \in \mathfrak{F}_1 \) (see (2)).

Einstein condition: \( \frac{3}{2} a = b = c \).
\[ \lambda_7 \mapsto \eta_3 \oplus \mathbb{R}X_3 \oplus \mathbb{R}X_4 \].

\[ p_{\lambda_7} = \{ d([a \ b \ c \ d], a + b, 2a, b, c, b, 2a + b) : a, b, c, d, a \in \mathbb{R} \} , \]

\[ K_{\lambda_7} = \{ d(H_1, H_2, d, \det H_1), H_1, H_2 \in O(2) \} , a_{\lambda_7} = \{ d(a, b, c, d, a + b) : a, b, c, d \in \mathbb{R} \} . \]

The corresponding 6-dimensional solitons are parameterized by:

\[
S_3^3 = \{(a, b, c, d) \in S^3 : a > 0, a \geq b, c \geq d, \text{ or } \ a = b = 0, c \geq |d| \} .
\]

Einstein condition: \( a = b, c = d, 3a = 2c \).

\[ \lambda_8 \mapsto \eta_1 \oplus \mathbb{R}X_3 \].

\[ a_{\lambda_8} = p_{\lambda_8} = \{ d(a, b, c, a + b, 2a + b) : a, b, c \in \mathbb{R} \} , \quad K_{\lambda_8} \subseteq d(\pm 1, \ldots, \pm 1) . \]

The corresponding 6-dimensional solitons are parameterized by \( \mathbb{R}^3 \). Recall that \( \mathbb{R}^3 \) can be parameterized as

\[ \mathbb{R}^3 = \{(x_1, x_2, x_3) \in S^2 : x_3 > 0, \text{ or } x_3 = 0, (x_1, x_2) \in \mathbb{R}^2 \} , \]

where \( S^n \) is the \( n \)-dimensional sphere in \( \mathbb{R}^{n+1} \) and \( \mathbb{R}^2 \) is seen as in (10) below. Einstein condition: \( 2a = b = 2 \frac{c}{3} \).

4.4. Nilsolitons of dimension 6. (see Tables 5 and 2)

\[ \mu_6 \]

\[ a_{\mu_6} = p_{\mu_6} = \{ d(a, b, a + b, 2a + b, a + 2b, 3a + b) : a, b \in \mathbb{R} \} , \quad K_{\mu_6} \subseteq d(\pm 1, \ldots, \pm 1) . \]

The space of corresponding 7-dimensional solitons is parameterized by \( \mathbb{R}^2 \), that can be seen as

\[
\mathbb{R}^2 = \{(x, y) \in S^1 : y \geq 0, x > -1 \} .
\]

Einstein condition: \( 2a = b \).

Remark 4.2. As we shall see, although \( \mu_7 \) and \( \mu_8 \) are very similar, \( p_{\mu_7} \) is completely different from \( p_{\mu_8} \) (see Table 5). In particular, \( p_{\mu_7} \) has non-diagonal elements (with respect to our fixed basis) but \( p_{\mu_8} \) does not. The same holds for \( \mu_{15} \) and \( \mu_{16} \).

\[ \mu_{17} \]

\[ a_{\mu_{17}} = p_{\mu_{17}} = \{ d([a \ b \ c] \ b, 2a, [3a \ b] \ 3a, 4a) : a, b \in \mathbb{R} \} , \quad K_{\mu_{17}} = \{ d(A, \varepsilon, A, A) : A = \begin{bmatrix} 0 & \pm 1 \\ \pm 1 & 0 \end{bmatrix} \} . \]

The corresponding set of 7-dimensional solitons is parameterized by \( \mathbb{R}^2 \). Einstein condition: \( b = 0 \).

\[ \mu_{15} \]

For simplicity we will change the basis to \( \beta = \{ X_1, X_3, X_2, X_4, X_5, X_6 \} \). We then get that

\[ a_{\mu_{15}} = p_{\mu_{15}} = \{ d([a \ c] \ b, [a + b \ c \ a + b] \ 2a + b) : a, b, c \in \mathbb{R} \} . \]

It is not hard to see that the action of \( K_{\mu_{15}} \) can at most permute \( c \) with \( -c \) (that is \( a + c \) with \( a - c \) and this can be done by acting with \( A = d([0 \ 1], 1, [0 \ -1], -1) \in K_{\mu_{15}} \).
The corresponding set of 7-dimensional solsollitons is parameterized by:

\[ \mathfrak{S}_3^2 = \left\{ (a, b, c) \in S^2 : \begin{array}{l} a > 0, \ c \geq 0, \ \text{or} \\ a = 0, \ b, c \geq 0 \end{array} \right\} . \]

Einstein condition: \( a = b, c = 0. \)

- **\( \mu_{16} \).** As in the previous case, we will change the basis to \( \beta \) to get nicer matrices. In this way we get that

\[ \mathfrak{p}_{\mu_{16}} = \{ d(a, a, b, a + b, 2a + b) : a, b \in \mathbb{R} \}, \]

\[ \mathfrak{K}_{\mu_{16}} = \{ d(H_1, \pm 1, H_2, \pm 1) : H_i \in O(2) \} . \]

The corresponding set of 7-dimensional solsollitons is parameterized by \( \mathbb{P} \mathbb{R}^2 \).

Einstein condition: \( a = b, \)

- **\( \mu_{17} \) (\( \simeq \lambda_5 \oplus \mathbb{R} X_3 \)).**

\[ \mathfrak{p}_{\mu_{17}} = \{ d\left( \begin{bmatrix} a & d \\ d & b \end{bmatrix} , c, a + b, 2a + b \right) : a, b, c, d \in \mathbb{R} \}, \]

\[ \mathfrak{a}_{\mu_{17}} = \{ d(a, b, c, a + b, 2a + b, a + 2b) : a, b, c \in \mathbb{R} \} , \]

\[ \mathfrak{K}_{\mu_{17}} = \{ d(A, \pm 1, \varepsilon, \varepsilon A) : A \in O(2), \varepsilon = \det A \} . \]

The corresponding set of 7-dimensional solsollitons is parameterized by \( \mathfrak{F}_1^2 \) (see (7)). Einstein condition \( a = b \).

- **\( \mu_{20} \).**

\[ \mathfrak{p}_{\mu_{20}} = \{ d(a, a, b + a, a + b, 2a + b, a + 2b) : a, b \in \mathbb{R} \} . \]

A straightforward calculation shows that one can permute \( a \) with \( b \) by acting with

\[ A = d(\begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} , \frac{1}{\sqrt{2}} \frac{i}{\sqrt{2}} , \frac{-1}{\sqrt{2}} \frac{1}{\sqrt{2}} , \) \( \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} ) \in \mathfrak{K}_{\mu_{20}}. \]

The corresponding set of 7-dimensional solsollitons is parameterized by \( \mathfrak{F}_1^1 \) (see (6)). Einstein condition: \( a = b \).

- **\( \mu_{24} \) (free 3-step nilpotent Lie algebra with 2 generators).**

\[ \mathfrak{p}_{\mu_{24}} = \{ d\left( \begin{bmatrix} a & d \\ c & e \end{bmatrix} , \begin{bmatrix} a + d & e \\ e & a + f \end{bmatrix} : a, b, c, d, e, f \in \mathbb{R} \} , \]

\[ \mathfrak{a}_{\mu_{24}} = \{ D(a, b, c) = d(a, b, c + a + b, a + c, b + c) : a, b, c \in \mathbb{R} \} . \]

It is known that \( \mathfrak{K}_{\mu_{24}} \simeq O(3) \) and from this we can see that the corresponding set of 7-dimensional solsollitons is parameterized by \( \mathbb{P} \mathbb{R}^3 / S_3 \). This set can be explicitly realized as

\[ \mathfrak{F}_1^2 = \left\{ (a, b, c) \in S^2 : \begin{array}{l} b > 0, \ a \geq b \geq c \text{, or} \\ b = 0, \ c \leq a \end{array} \right\} . \]

(see also \( \mathfrak{F}_1^3 \)). It will be Einstein if and only if \( a = b = c \).

- **\( \mu_{25} \).**

\[ \mathfrak{p}_{\mu_{25}} = \{ d(a, b, \begin{bmatrix} c \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} , a + b, 2a + b) : a, b, c, \beta \in \mathbb{R} \} , \]

\[ \mathfrak{a}_{\mu_{25}} = \{ d(a, b, c, 2a + b - c, a + b, 2a + b) : a, b, c \in \mathbb{R} \} , \]

\[ \mathfrak{K}_{\mu_{25}} = \{ d(\nu, \varepsilon, H, \varepsilon \nu, \varepsilon), H \in O(2), \varepsilon = \det H, \nu = \pm 1 \} . \]

The corresponding set of 7-dimensional solsollitons is parameterized by \( \mathbb{P} \mathbb{R}^3 \).

Einstein condition: \( a/5 = b/8 = c/9 \)
\( \bullet \mu_{26} \ (\cong \eta_1 \oplus (\mathbb{R}X_3 \oplus \mathbb{R}X_4)) \).

\[ p_{\mu_{26}} = \{ d(a, b, C, d, e) : a, b, d, e \in \mathbb{R}, C \in \text{sym}(2) \}, \]

\[ a_{\mu_{26}} = \{ d(a, b, c, d, a + b, 2a + b) : a, b, c, d \in \mathbb{R} \}, \]

\[ K_{\mu_{26}} = \{ d(\varepsilon, \nu, H, c\varepsilon\nu), H \in O(2), \varepsilon = \pm 1, \nu = \pm 1 \}. \]

\( W_{\mu_{26}} \): group generated by \( \{ 1d, \sigma_{34} \} \).

The corresponding set of 7-dimensional solitons is parameterized by (see (7)):

\[ \mathfrak{g}^3_2 = \left\{ (a, b, c, d) \in S^3 : \begin{array}{l}
    a > 0, c \leq d, \text{or} \\
    a = 0, (c, d, b) \in \mathfrak{g}^2_1
\end{array} \right\}. \]

Einstein condition: \( 6a = 3b = 2c = 2d \).

\( \bullet \mu_{27} \ (\cong \lambda_3 \oplus \mathbb{R}X_3) \).

\[ a = p_{\mu}(\mu_{27}) = \{ d(a, b, c, 2b, a + b, a + 2b) : a, b, c \in \mathbb{R} \}, K(\mu_{27}) \subset d(\pm 1, \ldots, \pm 1). \]

The corresponding set of 7-dimensional solitons is parameterized by \( \mathbb{PR}^3 \).

Einstein condition: \( 4a = 3b = 2c \).

\( \bullet \mu_{28} \).

\[ p_{\mu_{28}} = \left\{ d\left( \begin{bmatrix}
    a & B \\
    b & a
\end{bmatrix}, a + b, a + b : B = \begin{bmatrix}
    a & \beta \\
    \beta & a
\end{bmatrix}, a, b, \alpha, \beta \in \mathbb{R} \right\}, \]

\[ a_{\mu_{28}} = \{ d(a, b, b, a + b, a + b) : \alpha, \beta \in \mathbb{R} \} \]

The group \( \text{Aut}(\mu_{28}) \) has been calculated in [18], although it is easy to see that we can interchange \( a \) with \( b \) (the only possible permutation) by conjugating by

\[ A = \begin{bmatrix}
    1 & 1 & 1 \\
    1 & 1 & 1 \\
    1 & 1 & 1
\end{bmatrix} \in K_{\mu_{28}}. \]

The corresponding set of 7-dimensional solitons is parameterized by \( \mathfrak{g}_1 \) (see (8)).

Einstein condition: \( a = b \).

\( \bullet \mu_{29} \).

\[ p_{\mu_{29}} = \left\{ d\left( \begin{bmatrix}
    a & B \\
    0 & c
\end{bmatrix}, a + b, b + c : a, b, c, \alpha \in \mathbb{R} \right\}, \]

\[ a_{\mu_{29}} = \{ d(a, b, c, a + b, a + b, b + c) : a, b, c \in \mathbb{R} \}, \]

\[ K_{\mu_{29}} = \{ d\left( \begin{bmatrix}
    a & d \\
    c & b
\end{bmatrix}, d, c, \varepsilon, 1 : H = \begin{bmatrix}
    a & b \\
    c & d
\end{bmatrix} \in O(2), \varepsilon = \det H \right\}. \]

The corresponding set of 7-dimensional solitons is parameterized by \( \mathbb{PR}^3 \).

Einstein condition: \( a = b, 4a = 3c \).

\( \bullet \mu_{30} \ (\cong h_3 \oplus h_3) \).

\[ p_{\mu_{30}} = \left\{ d\left( \begin{bmatrix}
    a & B \\
    0 & c
\end{bmatrix}, a + b, c + d : a, b, c, \alpha, \beta \in \mathbb{R} \right\}, \]

\[ a_{\mu_{30}} = \{ d(a, b, c, d, a + b, c + d) : a, b, c, d \in \mathbb{R} \}, \]

\[ K_{\mu_{30}}^0 = \{ d(H_1, H_2, \varepsilon_1, \varepsilon_2) : H_i \in O(2), \varepsilon_i = \det H_i, i = 1, 2 \}. \]

In this case there is another orthogonal automorphism that act non trivially on subspaces of \( a \) given by \( H = d\left( \begin{bmatrix}
    0 & H' \\
    H' & 0
\end{bmatrix}, H' \right) \) where \( H' = \begin{bmatrix}
    0 & 1 \\
    1 & 0
\end{bmatrix} \), and therefore we get that \( W_{\mu_{30}} \) is the group generated by \( \{ 1d, \sigma_{12}, \sigma_{34}, \sigma_{13}\sigma_{24} \} \). It is easy to see that \( \mathbb{R}^4/W_{\mu_{30}} \) can be parameterized by

\[ \mathcal{D} = \left\{ (a, b, c, d) \in \mathbb{R}^4 : a \geq b, c \geq d, \begin{array}{l}
    a + b > c + d \text{ or} \\
    a + b = c + d \text{ and } a \geq c
\end{array} \right\}. \]
Hence, the corresponding set of 7-dimensional solsolitons is parameterized by
\[ \mathfrak{F}^3_3 = \{(a, b, c, d) \in S^3 \cap \mathcal{D} : a > 0, a + b \geq 0, \text{ if } d, c + d < 0 \text{ then } a > |d|\} . \]

Einstein condition: \(a = b = c = d\).

- \( \mu_{31} (\simeq \mathfrak{h}_6 \oplus \mathbb{R}X_4) \).
  \[ p_{\mu_{31}} = \{d(a, [b, \alpha]_c, d, \alpha + b, a + c] : a, b, c, d, \alpha \in \mathbb{R}\} , \]
  \[ a_{\mu_{31}} = \{d(a, b, c, d, a + b, a + c) : a, b, c, d, \in \mathbb{R}\} , \]
  \[ K_{\mu_{31}} = \{d(\varepsilon, H, \pm 1, \varepsilon H), H \in O(2), \varepsilon = \pm 1\} . \]

\( W_{\mu_{31}} \) : generated by \{Id, \sigma_{23}\}.

The corresponding set of 7-dimensional solsolitons is parameterized by: \((a, b, c, d) \in \mathfrak{F}^3_3\) (see (12)).

Einstein condition: \(4b = 4c = 6a = 3d\).

- \( \mu_{32} (\simeq \mathfrak{h}_6 \oplus \mathbb{R}X_5) \) (see \( \lambda_4 \)).
  \[ p_{\mu_{32}} = \{d(a, b, c, d, a + b) : a, b, c, d, \alpha, \beta \in \mathbb{R}, B = [\alpha, \beta]_{a+b-c} , \]
  \[ a_{\mu_{32}} = \{d(a, b, c, a + b-c, d, a+b) : a, b, c, d, \in \mathbb{R}\} . \]

Concerning \( K_{\mu_{32}} \) we can use the information we have from \( \lambda_4 \) to obtain that the corresponding set of 7-dimensional solsolitons is parameterized by:
\[ \mathfrak{F}^3_4 = \left\{(a, b, c, d) \in S^3 : d > 0, a \geq c \geq b, 2c \geq a + b, \text{ or }\right\} . \]

Einstein condition: \(a = b = c = \frac{3}{4}d\).

- \( \mu_{33} (\simeq \mathfrak{h}_3 \oplus (\mathbb{R}X_3 + \mathbb{R}X_4 + \mathbb{R}X_5)) \).
  \[ p_{\mu_{33}} = \{d(A, B, \text{tr} A) : A \in \text{sym}(2), B \in \text{sym}(3)\} , \]
  \[ a_{\mu_{33}} = \{d(a, b, c, d, e, a + b) : a, b, c, d, e, \in \mathbb{R}\} , \]
  \[ K_{\mu_{33}} = \{d(H_1, H_2, \varepsilon) : H_1 \in O(2), H_2 \in O(3), \varepsilon = \det H_1\} . \]

The corresponding set of 7-dimensional solsolitons is parameterized by (see (4)):
\[ \mathfrak{F}^4_1 = \left\{(a, b, c, d, e) \in S^4 : e > 0, a \geq b, c \leq d \leq e, \text{ or }\right\} . \]

Einstein condition: \(3a = 3b = 2c = 2d = 2e\).

5. The Space of \( m \)-dimensional SolSolitons

We now describe the space of solsolitons for each dimension \( \leq 6 \), by using the information we have obtained in the previous sections. Let \( \mathfrak{n}_\mu = (\mathbb{R}^n, \mu) \) denote a metric nilpotent Lie algebra as in Sections 3 and 4 (recall that we have fixed a basis and an inner product on \( \mathbb{R}^n \)). Hence we may denote as
\[ \mathcal{N} \text{il}(n) = \{\mathfrak{n}_\mu = (\mathbb{R}^n, \mu) : \mu \text{ is a nilsoliton}\} \]
the moduli space of nilsolitons of dimension \( n \) up to isometry and scaling.
According to Proposition 2.2 and 3, we have that $\text{Sol}(m)$, the moduli space of $m$-dimensional solsolitons (up to isometry and scaling), is given by:

$$\text{Sol}(m) = \bigcup_{\mu \in \text{Nil}(m)} \{ a \oplus n_\mu : a \in \text{Gr}_{m-n}(a_\mu)/W_\mu \},$$

where $\text{Gr}_r(a_\mu)$ is the Grassmanian of $r$-dimensional subspaces of $a_\mu$ (a maximal abelian subspace of symmetric derivations of $n_\mu$), and $a \oplus n_\mu$ is the (metric) solvable Lie algebra constructed in Proposition 2.2. Recall that $m - n \leq k_\mu = \text{rank}(\mu) = \dim(a_\mu)$.

We note that each $\mu \in \text{Nil}(m)$ is a point in $\text{Sol}(m)$ and the same holds when $m = k_\mu + \dim n_\mu$, that is, there is a single point $a_\mu \oplus n_\mu$ in the space of solsolitons of dimension $k_\mu + \dim n_\mu$, which is always Einstein (see Section 3.2).

More explicitly, we can see this moduli space as

$$\text{Sol}(m) = \text{Nil}(m) \cup \bigcup_{\mu \in \text{Nil}(m-1)} \text{Gr}_1(a_\mu)/W_\mu \cup \bigcup_{\mu \in \text{Nil}(m-2)} \text{Gr}_2(a_\mu)/W_\mu$$

$$\cup \cdots \cup \bigcup_{\mu \in \text{Nil}(m-r)} \text{Gr}_r(a_\mu)/W_\mu \cup \cdots \cup \bigcup_{\mu \in \text{Nil}(m/2)} \text{Gr}_{m/2}(a_\mu)/W_\mu$$

for $m$ even. For $m$ odd, the same expression holds but without the last union.

We therefore get that:

$\text{Sol}(2) = \{ \mathbb{R}^2 \} \cup \{(a_\mathbb{R} \oplus \mathbb{R})\},$

$\text{Sol}(3) = \{ \mathfrak{h}_3, \mathbb{R}^3 \} \cup \mathfrak{sl}_2(\mathbb{R}),$

$\text{Sol}(4) = \{ \eta_1, \eta_2, \mathbb{R}^4 \} \cup \mathfrak{sl}_2(\mathbb{R}) \cup (\mathbb{P}^3/\mathbb{S}_3)(\mathbb{R}^3) \cup \{ a_{\mathbb{R}^2} \oplus \mathbb{R}^2 \},$

$\text{Sol}(5) = \{ \lambda_1, \ldots, \lambda_5, \mathbb{R}^5 \} \cup \mathbb{P}^2(\eta_1) \cup \mathfrak{sl}_2^3(\mathbb{R}) \cup (\mathbb{P}^3/\mathbb{S}_3)(\mathbb{R}^4) \cup \{ a_\mathbb{R}_3 \oplus \mathfrak{h}_3 \} \cup \mathfrak{so}_3(\mathbb{R}^3),$

$\text{Sol}(6) = \{ \mu_1, \ldots, \mu_{33}, \mathbb{R}^6 \} \cup \mathbb{P}^2(\lambda_1) \cup \mathbb{P}^2(\lambda_2) \cup \mathbb{P}^2(\lambda_3) \cup \mathfrak{so}_2(\lambda_4) \cup \mathfrak{so}_3(\lambda_5)$

$$\cup \mathfrak{so}_2(\lambda_6) \cup \mathfrak{so}_2^3(\lambda_7) \cup \mathbb{P}^3(\lambda_8) \cup (\mathbb{P}^3/\mathbb{S}_3)(\mathbb{R}^5) \cup \{ a_{\mathbb{R}^3} \oplus \eta_1 \} \cup \mathfrak{so}_3^1(\eta_2)$$

$$\cup (\mathbb{Gr}_2(\mathbb{R}^3)/\mathbb{S}_3)(\mathbb{R}^4) \cup \{ a_{\mathbb{R}^3} \oplus \mathbb{R}^3 \},$$

where to distinguish among them we have added the metric to the notation. In this way, for example, $\mathbb{P}^2(\lambda_1)$ denotes the set of $\mu$ that parameterize the solsolitons associated to $\lambda_1$ in the corresponding dimension. We note that $\mathfrak{so}_3^j(\mu)$ is a real semialgebraic set and $\dim \mathfrak{so}_3^j(\mu) = j$.

Moreover, with the results we have obtained in Sections 3 and 4, one can describe $\text{Sol}(7)$. It is given by the union of the following:

- one point for each element of $\text{Nil}(7)$;
- six points coming from the 6-dimensional nisilons from Table 4, eight copies of $\mathbb{P}^2$ and three copies of $\mathbb{P}^3$ coming from the 6-dimensional algebras in Table 7 and the one-dimensional extensions of the elements in $\text{Nil}(6)$, described in Section 4;
three points corresponding to $\lambda_1, \lambda_3$ and $\lambda_5$, $\mathbb{F}^2(\lambda_4)$, $\mathbb{F}^2(\lambda_6)$, $\mathbb{F}\mathbb{R}(\lambda_8)$,
(Gr$_2(\mathbb{R}^4)/W_{\lambda_4})/(\lambda_7)$ and (Gr$_2(\mathbb{R}^5)/S_6)(\mathbb{R}^5)$;
• one point corresponding to $n_2$ and (Gr$_3(\mathbb{R}^4)/S_4)(\mathbb{R}^4)$.

We note that in this union not everything is explicitly described. Indeed, on
the one hand we have $Nil(7)$ which, as far as we know, has not been completely
classified yet. On the other hand, for $r > 0$, besides from the abelian cases, there
is just one set which is not explicitly described in this paper: (Gr$_2(\mathbb{R}^4)/W_{\lambda_7})/(\lambda_7)$.

6. On negatively curved solsolitons.

To conclude we will make some remarks on the curvature of solsolitons. Recall
that for each nilsoliton $\mu \in Nil(n)$ and each $1 \leq r < k_\mu = \text{rank}(\mu)$, we have a
family $X_{\mu,r}$ of solsolitons obtained from $\mu$ as in Proposition 2.2. Moreover, this
family is parameterized, up to isometry and scaling, by

$$X_{\mu,r} = \text{Gr}_r(\mathfrak{a}_\mu)/W_{\mu}.$$  

We then have that $X_{\mu,r} \subset \text{Sol}(m)$ where $m = r + n$, and since $W_{\mu}$ is a finite
group, this quotient inherits many of the properties of Gr$_r(\mathfrak{a}_\mu)$. In particular, $X_{\mu,r}$
is always connected and compact. Recall that if rank($\mu$) = 1, there is only one
solsoliton associated to $\mu$ (up to isometry and scaling) and therefore $X_{\mu,1}$ is just a
single point, as it always happens for $X_{\mu,k_\mu}$.

**Remark 6.1.** Since a solsoliton $\mathfrak{s} = \mathfrak{a} \oplus \mathfrak{n}_s \in X_{\mu,r}$ is Einstein if and only if $D_1 \in \mathfrak{a}$
(see Proposition 2.2), then for $1 \leq r < k_\mu$ the set of non-Einstein solsolitons is open
and dense in $X_{\mu,r}$.

**Remark 6.2.** Let $\mathfrak{s}_0 \in X_{\mu,r}$ be an Einstein solvmanifold with sectional curvature
$K_{\mathfrak{s}_0} < 0$. By continuity, if $\mathfrak{s} \in X_{\mu,r}$ is sufficiently close to $\mathfrak{s}_0$ then $K_{\mathfrak{s}} < 0$ as well.
The classification of Einstein solvmanifolds with non-positive sectional curvature
of dimension $\leq 6$ is given in [13, Theorem 2] and Table 3. Let us look at $\lambda_6$, for example. We have that rank($\lambda_6$) = 3
and $X_{\lambda_6,1} \cong \mathbb{F}^2$ (see [7]), and therefore around $\mathfrak{s}_0 = \mathbb{R}D_1 \oplus \mathfrak{n}_{\lambda_6} \in X_{\lambda_6,r}$ we get a
2-dimensional family of 6-dimensional non-Einstein solsolitons with $K_{\mathfrak{s}} < 0$.

**Remark 6.3.** We can apply the same argument to the Ricci curvature. In fact, if
$\mathfrak{s}_0 \in X_{\mu,r}$ is Einstein then Ric$_{\mathfrak{s}_0} = c \text{Id} < 0$ and therefore Ric$_{\mathfrak{s}} < 0$ for any $\mathfrak{s} \in X_{\mu,r}$
sufficiently close to $\mathfrak{s}_0$.

On the other hand, if $\mathfrak{a} \in \mathfrak{a}_\mu$ is orthogonal to $D_1$ then $\mathfrak{s}_A = \mathbb{R}A \oplus \mathfrak{n}_\mu \in X_{\mu,1}$ does
not have negative Ricci curvature. Indeed, it is easy to see that $\langle D_1, A \rangle = 0$ implies
that Ric$_{\mathfrak{s}_A}$ has both signs. Note that therefore $\mathfrak{s}_A$ has not negative sectional curvature either. The same holds if $\mathfrak{a} \subset \mathfrak{a}_\mu$ is a
subspace orthogonal to $D_1$. For the existence of these orthogonal subspaces we
need rank($\mu$) > 1.

Summarizing, if $1 \leq r < \text{rank}(\mu)$ then there are always some solsolitons in $X_{\mu,r}$
with negative Ricci curvature and others with positive and negative directions for the
Ricci tensor.

**Remark 6.4.** For $\mu = 0$ and $r = 1$, the Einstein solsoliton corresponds to the real
hyperbolic space $\mathbb{R}H^{n+1}$, which is a symmetric space of constant negative sectional
curvature. According to the above observations, in $X_{0,1} = \mathbb{R}^n/S_n$, there are families depending on $n - 1$ parameters of non-Einstein solsolitons with $K_s < 0$ as closed as one wishes to $\mathbb{R}H^{n+1}$.

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