GAMMA RADIATION SHIELDING PROPERTIES OF NATURAL GLASS OBSIDIAN

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ABSTRACT

The present study was carried out to estimate gamma radiation shielding properties of natural glass obsidian. For this purpose, linear attenuation coefficient, mass attenuation coefficient, mean free path, half-value layer, tenth-value layer, effective atomic number and effective electron number values of obsidian samples in black and brown colors were experimentally measured for 661.66, 1172.23 and 1332.48 keV gamma ray energies obtained from $^{137}$Cs, and $^{60}$Co radioactive sources. Measurements were performed at narrow-beam transmission geometry using a NaI(Tl) scintillation detector. In addition, all these parameters were theoretically calculated by using WinXCOM program in the energy region of 0.015 to 15 MeV. A good agreement was observed between theoretical and experimental values. Furthermore, energy absorption and exposure buildup factors (EABF and EBF) of obsidian specimens were determined in the energy range of 0.015 to 15 MeV using G-P fitting method. Finally, it can be concluded that these naturally occurring volcanic glasses can be used for radiation shielding applications.

Keywords: Gamma Radiation, Mass attenuation coefficient, Obsidian, G-P fitting method, Buildup factors

1. INTRODUCTION

Over the past few years, growing attention in proper designs for radiation protection has raised up due to the increased use of radiation in various fields such as science, technology, industry, agriculture, medicine, radiation biophysics, radiation protection, and the like [1,2]. Many researchers have focused on the design of barrier materials that reduce or retain X and gamma rays to protect both humans and their environment from the harmful impacts of radiation. Mass attenuation coefficients ($\mu_m$), linear attenuation coefficients ($\mu$) mean free path (mfp), half value layer (HVL), tenth value layer (TVL), effective atomic number ($Z_{eff}$), effective electron number ($N_{eff}$), energy absorption and exposure buildup factors (EABF and EBF) are basic qualities that shows the radiation shielding effectiveness of any materials [3].

Extensive surveys have been performed on investigation of radiation shielding features of various materials such as steels, alloys, superconductors, various ores concrete and etc. [4-14]. Moreover, there are several theoretical and experimental studies have been carried out to estimate radiation attenuation behaviors of various glass system. For example, Singh et al. experimentally determined $\mu_m$, $Z_{eff}$ and $N_{eff}$ values of $x$CaO•(0.3-$x$)SrO•0.7B$_2$O$_3$ glass system in the gamma ray energy range 511–1332 keV and compared with theoretical values obtained from WinXCOM program [15, 16]. The radiation shielding parameters ($\mu_m$, MFP, $Z_{eff}$ and $N_{eff}$) of the borosilicate glasses were theoretically computed using WinXCOM program by Chanthima and Kaewkhao [17]. The $\mu_m$, $Z_{eff}$ and $N_{eff}$ of WO$_3$ based glass system were experimentally determined for the 0.356 and 0.662 MeV energies and theoretically calculated at 1 keV-100 GeV using XCOM program by Mostafa et al. [18]. Sayyed et al. theoretically investigated gamma ray attenuation behaviors of different germanate glasses in the energy range of 15 keV to 10 MeV by using WinXCOM program [19]. Ahmed et al. experimentally measured different gamma ray attenuation parameters ($\mu$, $\mu_m$ and HVL) of some sodium iron phosphate glasses at the energy range of 121.8 keV to 1407.9 keV [20].
Obsidian, also called volcanic glass, is a type of aluminosilicate glass that contains iron impurities and some tiny particles (particle-dispersed glasses). It is an amorphous volcanic glass which have glassy brightness and is found in nature in black, gray, brown, red and green colors [21]. The chemical compositions of the studied obsidian are taken from the literature [22] and given in Table 1. When the main oxide contents of the both obsidian glasses evaluated, it is seen that they are rich in SiO₂ and alkalis (Na₂O + K₂O> 8%). The Na₂O content of each obsidian sample is higher than the K₂O content. As can be seen in Table 1, there is no significant difference in the main oxide contents of two samples and they have similar chemical composition in general terms. So far, there is not enough information in the literature about the radiation protection properties of this natural glass.

| Sample            | SiO₂ | TiO₂ | Al₂O₃ | Fe₂O₃ | MnO | MgO | CaO | Na₂O | K₂O | P₂O₅ |
|-------------------|------|------|-------|-------|-----|-----|-----|------|-----|------|
| Black Obsidian (BK) | 75,80 | 0,10 | 13,00 | 1,67  | 0,10| 0,10| 4,77| 3,97 | < 0,1|      |
| Brown Obsidian (BN) | 75,65 | 0,10 | 12,90 | 1,85  | 0,10| 0,10| 4,45| 4,20 | < 0,1|      |

The objectives of this present study was to estimate the gamma ray attenuation parameters such as μ, μₘ, mfp, HVL, TVL, Zeff and Neff in black and brown color obsidian specimens which were collected from Sarıkamış located in Turkey. These parameters were measured for 661.66, 1173.24 and 1332.48 keV photons obtained from 137Cs and 60Co point sources and the theoretical values of these parameters were calculated using the WinXCOM computer program. Besides these parameters EABF and EBF values were calculated at the 0.015 to 15 MeV energy range. The obtained experimental values for obsidian samples compared with the results of WinXCOM calculations. In addition, the radiation shielding capabilities of investigated naturally occurring volcanic glasses were investigated.

## 2. MATERIALS AND METHODS

Obsidian specimens were taken from Sarıkamış district of Kars province located at the eastern part of Turkey. The obtained samples were cut in a diameter of 18 mm and 5mm thickness using glass cutting machine to perform mass attenuation coefficient measurements. The prepared obsidian samples were irradiated with 661.66 keV gamma ray energy emitted from the 137Cs point source, 1173.23 and 1332.48 keV gamma ray energies from the 60Co point source.

A 3inch × 3 inch NaI(Tl) detector based on a gamma-ray spectrometer was employed to evaluate the unattenuated (I₀) and attenuated (I) photon intensities (Figure 1). The theoretical μₘ values for the prepared obsidian specimens were determined using WinXCOM software package [16]. Each spectrum was counted for a period of 600 s and repeated 3 times for both samples to obtain accurate measurement results. Maestro is a multi-channel analyzer (MCA) emulation software package from Ortec Company. The spectrum analyzed using an MCA system and Maestro based PC (Personal Computer) [23]. The typical spectra obtained from 137Cs and 60Co radioactive sources (for 661.66 keV, 1173.23 keV and 1332.48 keV) without attenuation and attenuated by present brown and black obsidian are given in Figure 2 (a and b). The μₘ values of obsidian specimens at different energies were experimentally determined using Lambert-Beer law given by equation 1 [12].

\[ I = I_0 e^{-\frac{\mu}{\rho}t} \]  

where, \( \mu/\rho = \mu_m \) is the mass attenuation coefficient in cm²/g, \( \rho \) is the density in g/cm³, \( t \) is sample mass thickness (the mass per unit area) in g/cm², \( I_0 \) and \( I \) are the incident and attenuated photon intensity, respectively.
2.1. Calculations

The theoretical $\mu_m$ values for any compound or mixture containing various elements is given by mixture rule [24]:

$$\mu_m = \frac{\mu}{\rho} = \sum w_i (\mu_m)_i$$  \hspace{1cm} (2)

where $w_i$ and $(\mu_m)_i$ are the weight fraction and mass attenuation coefficient for individual element in mixture, respectively. mfp which is the inverse of $\mu$ denotes the average distance between two successive interactions of photons and is calculated by the following formula [5, 25]:

$$mfp = \frac{1}{\mu}$$  \hspace{1cm} (3)

where, $\mu$ indicates the linear attenuation coefficient of absorber. A widely used important radiation shielding parameter (reflecting effectiveness of radiation protection for any materials) called half value layer is described as the sample thickness which reduces photon intensity to 50% of incident intensity and it can be computed according to following expression [19, 26]:

$$HVL = \frac{0.693}{\mu}$$  \hspace{1cm} (4)

Tenth value layer (TVL) in units of cm can be obtained by using following expression [5]:

$$TVL = \frac{\ln(10)}{M}$$  \hspace{1cm} (5)
The total atomic cross-sections in unit of barns/atom for investigated samples were computed using \( \mu_m \) values through the following expression [4, 25]:

\[
\sigma_a = \frac{\mu_m}{N_A \sum_i w_i A_i}
\]

where \( N_A \) is the Avogadro constant, \( A_i \) and \( w_i \) are atomic and fractional weights of the \( i^{th} \) element respectively. The total electronic cross-section, \( (\sigma_e) \), can be calculated according to equation 7 [7, 11].

\[
\sigma_e = \frac{1}{N_A} \sum_i f_i A_i Z_i (\mu_m)_i = \frac{\sigma_a}{Z_{eff}}
\]

where \( f_i \) and \( Z_i \) are the fractional abundance and atomic number of the \( i^{th} \) element respectively. The effective atomic number \( (Z_{eff}) \) represents the weighted average atomic number of the material consisting of different elements and can be obtained by using the following expression [3-5, 25]:

\[
Z_{eff} = \frac{\sigma_a}{\sigma_e}
\]

The effective electron number or electron density \( (N_{eff}) \), described as the number of electrons which interacted with photon per unit mass can be determined using \( \mu_m \) and \( \sigma_e \) values with the help of the following practical formula [3, 5, 25].

\[
N_{eff} = \frac{\mu_m}{\sigma_e}
\]

To compute EABF and EBF of obsidian glasses under the study, firstly equivalent atomic number \( (Z_{eq}) \) for specimens were calculated in the energy range 0.015–15 MeV using the ratio of Compton partial mass attenuation coefficient to total mass attenuation coefficient \( (R; (\mu_m)_{Compton} / (\mu_m)_{Total}) \) according to the following formula [13, 27-29];

\[
Z_{eq} = \frac{Z_1 (\log R_2 - \log R) + Z_2 (\log R - \log R_1)}{\log R_2 - \log R_1}
\]

where \( R \) represents the ratio of \( (\mu_m)_{Compton} \) and \( (\mu_m)_{Total} \) for an obsidian specimen at a given energy. \( Z_1 \) and \( Z_2 \) denotes atomic numbers of the elements corresponding to the ratios \( R_1 \) and \( R_2 \), respectively. More details can be found in Reference [29]. The values of \( (\mu_m)_{Compton} \) and \( (\mu_m)_{Total} \) at the selected energies were obtained using WinXCOM program for the elements \( Z = 4 \)–30 and both obsidian specimens. Then, G-P fitting coefficients \( (a, b, c, d, \) and \( X_k) \) for the specimens were generated by applying a similar interpolation method as given by the equation 12 [8, 28, 29];

\[
P = \frac{P_1 (\log Z_2 - \log Z_{eq}) + P_2 (\log Z_{eq} - \log Z_1)}{\log Z_2 - \log Z_1}
\]

where \( P \) denotes G-P fitting parameters of studied obsidian specimens. \( P_1 \) and \( P_2 \) are the values of G-P fitting coefficients corresponding to the \( Z_1 \) and \( Z_2 \) atomic numbers at a specific energy, respectively. The G-P fitting coefficients for elements were obtained from the ANSI/ANS-6.4.3 database, which provides the G-P fitting parameters for 23 elements \( (Z=4–92) \), water, air and concrete in the energy region of 0.015-15 MeV up to 40 mfp [30]. Finally, the computed G-P fitting coefficients were used to compute the EABF and EBF’s of obsidian glasses through the below equations [13, 31-33].
\[ B(E,X) = 1 + \frac{b + 1}{K - 1} (K^{X} - 1) \quad \text{for} \ K \neq 1 \]  
\[ B(E,X) = 1 + (b - 1)x \quad \text{for} \ K = 1 \]  
\[ K(E,X) = cx^a + d \frac{\tanh \left( \frac{x}{X_k} - 2 \right) - \tanh(-2)}{1 - \tanh(-2)} \quad \text{for} \ x \leq 40 \text{ mfp} \]  

where \( E, x \) and \( K(E, X) \) are the photon energy, penetration depth in mfp and dose multiplicative factor, respectively. \( a, b, c, d \) and \( X_k \) are the G-P fitting parameters.

### 3. RESULTS AND DISCUSSIONS

The mass attenuation coefficient for X and gamma rays is one of the most important quantity in many fields such as radiation physics, radiation dosimetry, biology, medicine, agriculture, industry and etc. [7]. It gives the number of photons scattered or absorbed and depends on the photon energy and the density of the absorber material [11]. The experimental \( \mu_m \) values of the natural obsidian glasses were measured at the 661.66, 1173.24 and 1332.48 keV photon energies and are summarized in Table 2. The average error in the experimental measurements in Table 2 is less than 6.2% and this error is caused by calculation of the peak area in the spectra, the measurement of the mass thickness of the samples and systematic errors. The measured \( \mu_m \) values were also compared with the theoretical values obtained by WinXCOM for the same photon energies and displayed in Table 3. As shown in this table, our measured \( \mu_m \) values are in a good agreement with the theoretically calculated values for obsidian glasses. The obtained \( \mu_m \) values of obsidian glasses are compared in Table 3 with the reported \( \mu_m \) values of the different materials for example, BB BG (barium–bismuth–borosilicate glass with the composition of 50BaO–5 Bi₂O₃ 45 borosilicate) [34], SIPG (sodium iron phosphate glass with the composition of 15 mole % Na₂O, 15 mole % Fe₂O₃. 70 mole % P₂O₅.5 CdO) [20], and also normal concrete [35]. It observed that the \( \mu_m \) values of our investigated both obsidians are compatible with the \( \mu_m \) values of BBG, SIPG and ordinary concrete (OC). For example, the theoretical \( \mu_m \) values of BBBG, SPIG and OC were 0.0795, 0.041 and 0.0780 cm²/g at 661.66 keV gamma ray energy. The measured \( \mu_m \) values of BNO and BKO for same gamma ray energy are similar to the value of BBBG and OC and slightly higher than the value of SPIG. Therefore, these obsidian glasses could be used as radiation shield especially at low energies.

**Table 2.** The experimental and theoretical mass attenuation coefficients (cm²/g) for obsidian specimens, BBG SPIG and OC

| Energy (keV) | BNO Exp. | BNO Theo. | BKO Exp. | BKO Theo. | BBBG Exp. | BBBG Theo. | SPIG Exp. | SPIG Theo. | OC Exp. | OC Theo. |
|-------------|----------|-----------|----------|-----------|-----------|------------|-----------|-----------|--------|---------|
| 661.66      | 0.07960 ± 0.00103 | 0.076891 | 0.07675 ± 0.00100 | 0.076890 | 0.0775 | 0.041 | 0.0780 |
| 1173.23     | 0.05831 ± 0.00462 | 0.058479 | 0.05746 ± 0.00459 | 0.058480 | 0.0561 | 0.044 | 0.0593 |
| 1332.48     | 0.05212 ± 0.00498 | 0.054789 | 0.05094 ± 0.00474 | 0.054790 | 0.0523 | 0.034 | 0.0555 |

In addition, the variations of experimentally and theoretically mass and linear attenuation coefficients with incident photon energy (0,015-15 MeV) were plotted in Figure 3. (a and b) to show the energy dependence of these parameters.
As can be seen from the Figure 3, values of both theoretical and experimental \( \mu_m \) values for investigated samples decreases with increase in the photon energy. As a result of our studies, it was seen that the \( \mu_m \) and \( \mu \) values for obsidian specimens decreases rapidly with the increasing of the incident photon energy up to 0.1 MeV. In the moderate energy region (0.1 - 5.0 MeV), these values decrease slowly with the increase of the incident photon energy. At last, in the \( E > 5 \) MeV the \( \mu_m \) and \( \mu \) values have a very weak dependence of the incident photon energy. This trend can be explained according to different interaction mechanism occurring between photons and material for different energy regions [4].

In the low energy region photoelectric interaction mechanism is dominant process and the photoelectric cross section depends inversely on the energy (i.e. \( E^{-3.5} \)) and depends upon the atomic number as \( Z^{4-5} \). In intermediate energies Compton Scattering becomes the dominant mechanism and the Compton cross section changes directly with the atomic number \( Z \), and inversely proportional to the photons energy (i.e. \( E^{-1} \)). Then, in the high energy region pair production is dominant interaction and the cross-section for pair production depends upon atomic number as \( Z^2 \) and logarithmically on energy (log \( E \)) [27, 29, 36].

In order to evaluate photon attenuation performance of studied samples mfp, HVL and TVL parameters were determined at photon energies of 0.015–15 MeV using the \( \mu \) values. The variations of mfp, HVL, and TVL values according to incident photon energy for brown and black obsidians are given in Figure 4 (a and b). It is seen that both the experimental and theoretical values of mfp, HVL and TVL, parameters increase with increasing photon energy. Our obtained trends in the mfp, HVL and TVL parameters are found to be similar with the reported these parameters for some AISI coded stainless steel [5]. At the 661.66, 1172.23 and 1332.48 keV gamma ray energies, the experimental mfp values were measured to be 4.907, 6.699 and 7.495 cm for Brown obsidian, 5.296, 7.074 and 7.980 cm for black obsidian respectively. The experimental HVL values of were measured to be 3.401, 4.633 and 5.195 cm for Brown obsidian, 3.671, 4.904 and 5.532 cm for black obsidian at the same gamma ray energies respectively. The experimental TVL values of were also measured to be 11.299, 15.474 and 17.257 cm for Brown obsidian, 12.195, 16.289 and 18.375 cm for black obsidian at the same gamma ray energies respectively.
Figure 4. (a, b) mfp, HVL and TVL values of brown and black obsidian versus incident photon energy.

On the other hand, the comparisons of measured and computed mfp, HVL and TVL values of both specimens with the values ordinary concrete are given in Figure 5 (a, b, c). According to these results, the mfp, HVL and TVL values of both specimens for 661.66, 1172.23 are lower than the values OC. At the same time, the experimental mfp, HVL and TVL values of BNO is slightly lower than the values OC for 1332.48 keV. It is seen that mfp, HVL and TVL values of black obsidian slightly higher than the values of brown obsidian, therefore, brown obsidian may be used as a better alternative gamma ray protection material than black obsidian. Because these parameters are commonly used quantities in order to estimate radiation shielding property of any material, having low mfp, HVL and TVL values generally exhibit better shielding features.
The effective atomic number (Z_{eff}) and effective electron numbers (N_{eff}) for obsidian specimens were calculated at photon energies of 0.015 – 15 MeV using equations 8 and 9, respectively. The variations of these parameters (Z_{eff} and N_{eff}) with the incident photon energy are given in Figure 6 (a and b), respectively. It observed that the measured Z_{eff} and N_{eff} values are consistent with the theoretically calculated values for obsidian glasses at 661.66 and 1173.24 keV, but Z_{eff} and N_{eff} values of both specimens for 1332.48 keV slightly lower than the calculated values. This dissimilarity may be arising from the difference between the measured and calculated \mu_m values. Hence Z_{eff} and N_{eff} values of BNO and BKO were compared with the values of OC. It seen that while the Z_{eff} values of both specimens are higher than the OC, N_{eff} values of both specimens are lower than the OC. It is well known that the N_{eff} value for any sample is closely related to the Z_{eff}, therefore variation of N_{eff} with the incident photon energy has a similar tendency to Z_{eff}. Also, one can be seen from Figure 6 (a and b) that the Z_{eff} and N_{eff} values of investigated obsidian specimens decreased with the increasing photon energy at low energies and maximum values of Z_{eff} and N_{eff} are below 0.1 MeV. In the moderate energies (0.1 - 5 MeV) Z_{eff} and N_{eff} have minimum values and started to slowly increase above about 5 MeV. The present trend in the Z_{eff} and N_{eff} values for the obsidian specimens can be understood according to the above mentioned three interaction mechanism between photon and material.

![Figure 6](image-url). The variations of Z_{eff} and N_{eff} against incident photon energies.

The buildup factors for obsidian samples were determined by using G–P fitting technique for some selected penetration depths (1, 5, 10, 20, 30 and 40 mfp) and in the region of 0.015 to 15 MeV. Table 3 and 4 represent the computed equivalent atomic numbers (Z_{eq}) and G-P fitting coefficients (a, b, c, d and Xk) for EABF and EBF of black and brown obsidian, respectively.
### Table 3. $Z_{eq}$ and G-P fitting coefficients for Brown Obsidian.

| Energy (MeV) | $Z_{eq}$ | EABF | EBF | Xk |
|--------------|----------|------|-----|----|
|              |          |      |     |    |
| 0.015        | 0.219    | 1.035| 0.391| -0.144|
| 0.02         | 0.205    | 1.082| 0.399| -0.106|
| 0.03         | 0.197    | 1.265| 0.432| -0.106|
| 0.04         | 0.156    | 1.580| 0.527| -0.084|
| 0.05         | 0.138    | 2.084| 0.598| -0.078|
| 0.06         | 0.092    | 2.617| 0.734| -0.065|
| 0.08         | 0.014    | 3.609| 0.999| -0.023|
| 0.1          | -0.037   | 4.155| 1.229| 0.001|
| 0.15         | -0.095   | 4.112| 1.545| 0.031|
| 0.2          | -0.113   | 3.658| 1.663| 0.039|
| 0.3          | -0.119   | 3.060| 1.698| 0.042|
| 0.4          | -0.113   | 2.753| 1.647| 0.037|
| 0.5          | -0.107   | 2.552| 1.598| 0.037|
| 0.6          | -0.100   | 2.415| 1.546| 0.034|
| 0.8          | -0.086   | 2.237| 1.450| 0.030|
| 1            | -0.075   | 2.114| 1.380| 0.027|
| 1.5          | -0.053   | 1.938| 1.251| 0.020|
| 2            | -0.035   | 1.841| 1.158| 0.013|
| 3            | -0.010   | 1.703| 1.053| -0.001|
| 4            | 0.008    | 1.612| 0.985| -0.011|
| 5            | 0.021    | 1.544| 0.940| -0.018|
| 6            | 0.024    | 1.478| 0.929| -0.027|
| 8            | 0.033    | 1.386| 0.904| -0.023|
| 10           | 0.034    | 1.321| 0.899| -0.028|
| 15           | 0.042    | 1.226| 0.882| -0.037|

### Table 4. $Z_{eq}$ and G-P fitting coefficients for Black Obsidian.

| Energy (MeV) | $Z_{eq}$ | EABF | EBF | Xk |
|--------------|----------|------|-----|----|
|              |          |      |     |    |
| 0.015        | 0.217    | 1.035| 0.393| -0.141|
| 0.02         | 0.201    | 1.083| 0.404| -0.104|
| 0.03         | 0.196    | 1.265| 0.434| -0.105|
| 0.04         | 0.154    | 1.590| 0.532| -0.082|
| 0.05         | 0.137    | 2.104| 0.602| -0.077|
| 0.06         | 0.090    | 2.642| 0.742| -0.063|
| 0.08         | 0.012    | 3.637| 1.099| -0.022|
| 0.1          | -0.039   | 4.175| 1.239| 0.003|
| 0.15         | -0.097   | 4.109| 1.555| 0.032|
| 0.2          | -0.114   | 3.655| 1.669| 0.040|
| 0.3          | -0.120   | 3.072| 1.072| 0.042|
| 0.4          | -0.114   | 2.748| 1.653| 0.038|
| 0.5          | -0.108   | 2.551| 1.601| 0.037|
| 0.6          | -0.100   | 2.415| 1.546| 0.034|
| 0.8          | -0.087   | 2.236| 1.453| 0.031|
| 1            | -0.075   | 2.114| 1.380| 0.027|
| 1.5          | -0.053   | 1.938| 1.251| 0.020|
| 2            | -0.035   | 1.842| 1.157| 0.013|
| 3            | -0.010   | 1.704| 1.053| -0.001|
| 4            | 0.008    | 1.613| 0.985| -0.011|
| 5            | 0.021    | 1.543| 0.941| -0.017|
| 6            | 0.024    | 1.478| 0.929| -0.027|
| 8            | 0.033    | 1.387| 0.904| -0.023|
| 10           | 0.035    | 1.321| 0.898| -0.028|
| 15           | 0.042    | 1.226| 0.883| -0.036|
The variation of EABF and EBF for both obsidian samples with incident photon energy at some selected penetration depth is plotted in Figure 7. The EABF and EBF values of both obsidian samples have similar tendency. This is because of the both obsidian specimens have similar chemical components. From this Figure 7, it can be observed that the EABF and EBF values of obsidians are small for all penetration depths at low energies where photoelectric absorption is the dominating photon interaction process and increases with the increase in the incident photon energy. In the intermediate energy region where the Compton scattering dominates EABF and EBF gradually increase and have highest value at 0.2 MeV for the samples. Then EABF and EBF values decrease with increase in photon energy for all penetration depths in the high energy region.

![Graphs of EABF and EBF for BNO and BKO samples](image)

Figure 7. (a-d) EABF and EBF values of the BNO and BKO versus photon energy at different mfps

The variations of EABF and EBF with penetration depth of obsidian specimens for 0.015, 0.15, 1.5 and 15 MeV are given in Figures 8 (a-d) and 9 (a-d), respectively. It is clear from these figures that EABF and EBF values increase with the increase of mean free path values. It was observed that the EABF and EBF values of black obsidian at 0.015 and 0.15 MeV energies are slightly higher than brown obsidian whereas EABF and EBF values obtained for both obsidians at 1.5 and 15 MeV energies are approximately the same.
Figure 8. (a–d). The EABF values for obsidian glasses versus the penetration depths at 0.015, 0.15, 1.5, 15 MeV energies.

Figure 9. (a–d). EBF values for obsidian glasses up to 40 mfp at 0.015, 0.15, 1.5, 15 MeV.
In addition, our obtained EABF and EBF values are consistent with the results obtained by Küçük and Gezer. They investigated EABF and EBF of black obsidian samples taken from Artvin and Van for 16 types of penetration depth in the energy range from 0.015 to 15 MeV [37].

Table 5 compares EBF values of examined samples with the OC some selected incident photon energies at 1 5, 10, 20 and 40 mfp penetration depths. The EBF values of both samples are slightly lower than the OC at 0.15 and 1.5 MeV photon energies for all penetration depths. At 15 MeV photon energy, EBF values of BNO and BKO slightly higher to the values of OC for all penetration depths except 1mfp. At 15 MeV photon energy, EBF values of BNO, BKO and OC very close to the each other for all penetration depths. From these comparisons, it can be concluded that these naturally occurring volcanic glasses can be used for radiation shielding applications especially for low energies.

4. CONCLUSION

In this paper, the gamma radiation shielding performance of naturally occurring obsidian glasses are investigated which are taken from the Sarıkamış district of Kars province located at the eastern part of Turkey. The $\mu$, $\mu_m$, mfp, HVL, TVL, $Z_{\text{eff}}$ and $N_{\text{eff}}$ parameters of these obsidian glasses have been investigated experimentally for the 661.66, 1172.23 and 1332.48 keV gamma ray energies and theoretically computed using the WinXCOM software. The results were compared with the OC commonly used. It was observed that the values of $\mu$, $\mu_m$, $Z_{\text{eff}}$ and $N_{\text{eff}}$ of obsidian glasses are quite high at low energies and decreased with increasing gamma ray energy. Also, the $\mu$ mfp, HVL and TVL values of BNO are smaller than the values of BKO and OC at 661.66, 1172.23 and 1332.48 keV gamma energies. The EABF and EBF values of these obsidian specimens were computed using G-P fitting method in the energy range 0.015-15 MeV up to 40 mfp penetration depths and variation of EABF and EBF photon energy and penetration depth were presented graphically. It was observed that EBF values brown obsidian generally less than black obsidian and ordinary concrete. Among the studied samples brown obsidian may be used as a better alternative gamma ray protection material than black obsidian especially for the low energy region.
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