Greatly Enhancing Catalytic Activity of Graphene by Doping the Underlying Metal Substrate

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Graphene-based solid-state catalysis represents a new direction in applications of graphene and has attracted a lot of interests recently. However, the difficulty in fine control and large-scale production of previously proposed graphene catalysts greatly limits their industrial applications. Here we present a novel way to enhance the catalytic activity of graphene, which is highly efficient yet easy to fabricate and control. By first-principles calculations, we show that when the underlying metal substrate is doped with impurities, the catalytic activity of the supported graphene can be drastically enhanced. Graphene supported on a Fe/Ni(111) surface is chosen as a model catalyst, and the chemical reaction of CO oxidation is used to probe the catalytic activity of graphene. When the underlying Fe/Ni(111) substrate is impurity free, the graphene is catalytically inactive. When a Zn atom is doped into the substrate, the catalytic activity of the supported graphene is greatly enhanced, and the reaction barrier of the catalyzed CO oxidation is reduced to less than 0.5 eV. Intriguing reaction mechanism of catalyzed CO oxidation is revealed. These studies suggest a new class of graphene-based catalysts and pave the way for future applications of graphene in solid-state catalysis.
Results and Discussions

Ni (111) is regarded as an excellent metal substrate for graphene due to the very small lattice mismatch between them. Graphene supported on Ni (111) has been extensively studied both experimentally and theoretically\(^1\)–\(^4\). It has been pointed out that the van der Waals (vdW) interactions play an essential role in theoretically reproducing the experimentally observed stable adsorption configuration of graphene on Ni (111)\(^5\)–\(^7\). A very recent experiment demonstrated that one or two monolayers of Fe can be ‘inserted’ between graphene and Ni (111)\(^8\). The Fe layers were shown to adopt the same lattice structure of Ni (111) surface and the supported graphene remains flat (without the Moiré pattern) after the Fe insertion. Unlike the case of graphene on Ni, previous DFT calculations without vdW interactions yielded adsorption energies and configurations of graphene on Fe/Ni (111) that agree well with experiments\(^9\). Since the vdW type of adsorption is not desired for graphene catalysis, in current study, we choose the two monolayer of Fe on Ni (111) as the underlying substrate, and then study the catalytic activity of the supported graphene.

The atomic model of graphene on Fe/Ni (111) is shown in Fig. 1a where five layers of Ni are included in the supercell to make sure that the correct surface properties of Ni (111) can be obtained. In structure optimizations, three bottom layers of Ni are fixed to bulk structures. Our DFT calculations show that the Ni-Fe and graphene-Fe bond lengths are around 2.48 Å and 2.15 Å, respectively. The adsorption energy of graphene is estimated to be 0.42 eV per carbon pair. All of these results agree very well with previous studies\(^9\). In Fig. 1b,c, we plot the isosurface of charge redistribution caused by the adsorption of graphene. The isosurface value of 0.008 e/Å\(^3\). The accumulation (depletion) of electrons is denoted by yellow (green).

Figure 1. (a) Atomic model of graphene supported on Fe/Ni (111) (Ni in dark red, Fe in grey, C in black). Five layers of Ni are included in the model. The supercell includes \(4 \times 4\) of pristine unit cell of graphene and 15 Å of vacuum in the direction perpendicular to the graphene surface; (b,c) are the top and side views of the isosurface of charge redistribution caused by the adsorption of graphene. The isosurface value of 0.008 e/Å\(^3\). The accumulation (depletion) of electrons is denoted by yellow (green).
and has been widely used as dopant in various kinds of applications. One Zn atom is doped in the top Fe layer as shown in Fig. 2. After the Zn-doping, graphene remains essentially flat, while the local electronic properties of graphene significantly change. To demonstrate this, we superimpose in Fig. 2 the impurity-induced charge redistribution which is defined as the difference between electron densities with and without the impurity. The figure suggests that the Zn-C bond loses electrons compared with Fe-C bonds, and there are electrons accumulated on the impurity atom and the graphene surface right on top of the impurity, which weaken the bonding between the C atom on top of the impurity and the metal surface and in turn change the local chemical reactivity. As a result, after the Zn doping, O$_2$ does bind to graphene at places around the C atom that forms chemical bond with the impurity Zn atom.

We show in Fig. 3a the most stable adsorption configuration of O$_2$ on graphene that is supported on Zn doped Fe/Ni (111). The adsorption energy of O$_2$ in this case is calculated to be around 0.42 eV. The O-O bond of the O$_2$ molecule is significantly elongated to 1.48 Å (compared with 1.24 Å in gas phase). We also tested other inequivalent adsorption sites near the Zn impurity (Figure S1 in the supporting information), and found that the site shown in Fig. 3a is the only one that shows chemisorption of O$_2$. When far away from the impurity atom, there is no binding of O$_2$. Note that CO still does not bind to graphene after doping. We also considered cases of doping multiple Zn atoms in the supercell and found that when two Zn impurities are close to each other, only one O$_2$ is adsorbed on graphene (see Figure S2 in the supporting information). The adsorption energy O$_2$ is calculated from equation,

$$ E_b = E_{\text{separated}} - E_{\text{combined}} $$

where $E_{\text{separated}}$ is the energy of the supported graphene and an O$_2$ that is 5 Å away from graphene surface.

We would like to point out here that since the system under study is big, the numerical error in total energy calculations could be big also so that the calculation of the adsorption energy of O$_2$ may not be accurate. To confirm the stability of the O$_2$ adsorption, we conducted ab initio molecular dynamics (MD) simulation of the system with adsorbed O$_2$ (Fig. 3a) at room temperature 300 K. The variations of bond lengths (the difference between the current length of the bond and its equilibrium value) of O-O
(O₂) and two C-O bonds (between graphene of O₂) as functions of MD steps are shown in Fig. 4. The variation of O-O bond length at room temperature is always less than 0.1 Å, and the bond lengths of two C-O bonds oscillate between ±0.1 Å after 2000 MD steps, indicating that the adsorption of O₂ on graphene after the impurity doping is stable at room temperature. To understand the interaction between the supported graphene and the O₂ molecule, we plot the isosurface of charge redistribution caused by O₂ adsorption in Fig. 3a. The side views of the charge redistribution are shown in Fig. 3b,c. It can be seen that electrons transferred to 2π’ orbital of O₂, causing the O-O bond elongated. The adsorption induced change of O₂ electronic structures is further illustrated in Fig. 3d where the projected density of

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Figure 3. (a) O₂ (O in red) adsorbed on graphene supported on Zn doped Fe/Ni (111). The adsorption energy of O₂ is around 0.42 eV. The isosurface of charge redistribution caused by O₂ adsorption is superimposed. (b,c) are side views of charge redistribution. (d) projected density of states (PDOS) of O₂ molecule before (upper) and after (lower) adsorption. The PDOS of O₂ before adsorption is calculated by putting the molecule 5 Å away from the graphene in the same supercell.
states (PDOS) of \( \text{O}_2 \) molecule before and after adsorption is plotted. Before adsorption, the \( \text{O}_2 \) molecule is magnetic and the \( 2\pi^* \) orbital of spin-down channel is unoccupied. After adsorption on graphene, three are electrons transferred to the originally empty \( 2\pi^* \), pulling the orbital below the Fermi energy and eliminating the magnetic moment of the molecule. This picture is consistent with the Bader charge analysis that nearly one electron is transferred to \( \text{O}_2 \) after adsorption.

The significantly elongated O-O bond of the adsorbed \( \text{O}_2 \) implies that the molecule is activated for CO oxidation. Since CO does not bind to graphene, there are no CO poisoning problems and we only need to consider the Eley-Rideal (ER) mechanism of the reaction for which the CO molecule approaches the \( \text{O}_2 \) from vacuum. As aforementioned, the reaction barrier is determined by c-NEB method. In Fig. 5, we show the energy profile along the reaction paths for two steps of the CO oxidation in a whole cycle of the catalyzed reaction (the first step: \( \text{CO} + \text{O}_2 \rightarrow \text{CO}_2 + \text{O} \) and the second step: \( \text{CO} + \text{O} \rightarrow \text{CO}_2 \)). Reaction barriers of both steps are low (<0.5 eV), indicating high catalytic activity of the supported graphene. In the first step of the reaction, the barrier of 0.47 eV originates from the breaking of the O-O bond and also one C-O bond between the \( \text{O}_2 \) and graphene. To get some rough ideas for how much these two bond breaking processes contribute to the barrier, we performed c-NEB calculations for the breaking of O-O bond on graphene (without the presence of the CO molecule). The results of calculations are shown in Figure S3 in the supporting information, from which we can see that around 0.31 eV of energy is needed to break the O-O bond on graphene surface. After one O atom of \( \text{O}_2 \) is taken away in the first step of reaction, the binding between the left-over O atom and the graphene becomes stronger, which can be seen from the fact that in this case the C-O bond is shorten from 1.50 Å to 1.35 Å. The shorter C-O bond length means more energy is needed to break the bond. The reaction barrier of the second step of the reaction is estimated to be 0.40 eV (Fig. 5b). The barrier is mainly from the breaking of the left-over C-O bond. The barriers of two steps are similar, suggesting that two steps are equally important and the reaction rate of the catalyzed CO oxidation is determined by both steps.

**Conclusion**

In summary, via *ab initio* calculations, we present a novel method of activating graphene for catalyzing chemical reaction of CO oxidation. When supported on impurity free Fe/Ni (111) surface, graphene is chemically inert. We show that after doping an impurity Zn atom in the underlying Fe layer, graphene is activated and become an excellent catalyst for CO oxidation. The reaction barriers of the two-step chemical reaction of CO oxidation catalyzed by graphene supported on Zn doped Fe/Ni (111) are estimated to be low (<0.5 eV). Unlike the previously proposed graphene-supported or covered catalysts, the catalysis center here is the graphene itself. Since CO does not bind to graphene, the CO poisoning is not an issue for the proposed catalyst here. Also, considering the cost-effective materials (Ni, Fe, Zn, C) we used here, the proposed method of activating graphene may have important applications in the future design of graphene-based catalysis. We believe that the observed substrate-impurity induced high
catalytic activity of graphene is not limited to the specific system discussed in this paper. Our studies pave a way for a new class of graphene-based solid-state catalysts.

**Computational Methods.** The first-principles calculations are performed with spin-polarized density functional theory (DFT) by utilizing Vienna ab-initio Simulation Package (VASP)²³,²⁴. Reaction barriers of catalyzed CO oxidations are determined by the Climbing Image Nudged Elastic Band (c-NEB) method²⁵. The generalized gradient approximation (GGA) in the Perdew-Burke-Ernzerh (PBE) format²⁶ and the projector-augmented wave (PAW) method²⁷ are employed in all calculations. A plane wave basis

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**Figure 5.** The energy profile of two-step CO oxidation catalyzed by graphene supported on Zn@Fe/Ni (111). The calculations are done by the c-NEB method. (a) the first step of the reaction: \( \text{CO} + \text{O}_2 \rightarrow \text{CO}_2 + \text{O}\). The initial state (IS): \(d(O-O)=1.48\ \text{Å}, d(C-O)=1.14\ \text{Å}, \text{and } d(\text{CO-O}_2)=3.10\ \text{Å}\); The transition state (TS): \(d(O-O)=2.12\ \text{Å}, d(C-O)=1.16\ \text{Å}, \text{and } d(\text{CO-O}_2)=1.85\ \text{Å}\), in the final state (FS), \(\text{CO}_2\) forms. (b) the second step of the reaction: \(\text{CO} + \text{O} \rightarrow \text{CO}_2\). The initial state (IS): \(d(C-O)=1.14\ \text{Å}, \text{and } d(\text{CO-O}_2)=3.30\ \text{Å}\); The transition state (TS): \(d(C-O)=1.17\ \text{Å}, \text{and } d(C-O)=1.76\ \text{Å}\). \(\text{CO}_2\) does not bind to the graphene surface in both cases. Note that \(d(O-O)\) denotes the O-O bond length of \(\text{O}_2\), \(d(C-O)\) the C-O bond length of \(\text{CO}\), and \(d(\text{CO-O}_2)\) the distance between the C atom in \(\text{CO}\) and the nearest O in \(\text{O}_2\).
with the cut-off energy of 450 eV and $3 \times 3 \times 3$ sampling of the Brillouin zone are used. The convergence criterion for structural relaxations is set to 0.01 eV/Å.

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**Author Contributions**

C.Z. designed and supervised the project. N.G. and Y.X. performed calculations. All contributed to the data analysis and writing of the paper.

**Additional Information**

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