First steps toward the development of SONATE, a Compact Accelerator driven Neutron Source

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Abstract.
Facilities providing bright thermal neutron beams are of primary importance for various research topics such as condensed matter experiments, neutron-imaging or medical applications. Currently these are mainly spallation sources and nuclear reactors. However, these later facilities are ageing and the political context does not favor the building of new ones. This is the case in CEA-Saclay (France), where the Orphée reactor is planned to shutdown in 2019. Therefore, another local facility, affordable by one country, able to provide high brilliance neutron beams has to be built. At CEA-Saclay, a compact accelerator driven neutron source, SONATE, is investigated in taking advantage of the IPHI accelerator able to deliver a 3 MeV proton beam with an intensity up to 100 mA. In the future, SONATE is foreseen to operate with 20 MeV protons to increase the neutron brightness. In addition to the difficulties to operate such high intensity accelerators, the other challenges regard the target-moderator-reflector (TMR) design which is crucial to maximize the neutron flux at the detector location. At CEA-Saclay, several experiments were performed between 2016 and 2019 with the IPHI accelerator. Geant4 simulations were also developed. They demonstrate the feasibility of such concept and enable to find the best TMR configuration for the future SONATE facility. These developments are reported in this article.

1 Introduction

In Europe, after the ILL nuclear reactor shutdown foreseen in 2030-2040 (optimistic scenario from [1]), the neutron beam time will drop roughly by 40%. The neutron beam time in Europe will then depend mainly on four big facilities: ESS (Sweden, start foreseen in 2024), ISIS (UK), SINQ (Switzerland) and FRM-II (Germany). The beam time will then be expensive and dedicated to high impact experiments. In this context, the development of new small facilities affordable by one country are necessary to support the large neutron community. It is of primary importance to develop new experimental techniques and to prepare scientists to use the big facilities. Fission, spallation and fusion reaction based facilities could be a solution but are not suited as local facilities because of either political choices, their high costs or their low neutron yields. Compact accelerator driven neutron sources (CANS), based on (p/d/n) reaction with light nucleus targets, can now be competitive with nuclear reactors in term of neutron brightness because of the recent developments of high intensity accelerators, and at low cost (less than 200 Meuros).

In France at CEA-Saclay, since the Orphée reactor will shutdown in autumn 2019, a CANS is expected to be developed taking advantage of the IPHI proton accelerator. This installation is able to deliver protons with an energy (E\textsubscript{p}) of 3 MeV and a high intensity up to 100 mA, in continuous or pulsed mode. For this purpose a target-moderator-reflector (TMR) assembly providing the highest well collimated thermal (E\textsubscript{n}<100 mev) neutron flux is searched for.

This paper starts by presenting the criteria leading to select the best suited target to generate the primary neutrons. Then the validation of the Geant4 simulation (version 10.04.p02) [2] is presented for different TMR configurations. Finally, the first optimization steps to find the best TMR configuration, for 3 MeV proton beam of few mA, are detailed focusing on the moderator and reflector materials and geometries.

2 Target design criteria

The target has to maximize the neutron yield (Y\textsubscript{n}) for 3 MeV protons and to handle the heat load deposited by the beam. At E\textsubscript{p}=3 MeV, to maximize Y\textsubscript{n} the target nucleus should have a low Coulomb barrier (a low nuclear charge) and a low reaction energy threshold (a low nucleon binding energy). To that respect the stable nuclei having the higher (p,n) reaction cross-section at 3 MeV are \textsuperscript{7}Li and \textsuperscript{9}Be [4]. Table 1 shows that for \textsuperscript{7}Li and \textsuperscript{9}Be thick target (thick enough to stop the protons inside the target), Y\textsubscript{n} are respectively 2.38x10\textsuperscript{-4} n/p and 5.23x10\textsuperscript{-5} n/p [3]. In increasing E\textsubscript{p} to 5 MeV, Y\textsubscript{n} will increase by a factor 4.2 and 9.5 respectively for \textsuperscript{7}Li and \textsuperscript{9}Be, which
Table 1: Neutronic and thermal properties for different material foreseen as a target. The neutron yields are from [3].

| Nucleus | $\rho$ [g/cm$^3$] | $E_{\text{th}}$ [MeV] | $Y_n$ [n/p] @3MeV | $Y_n$ [n/p] @5MeV | $T_f$ [°C] | $\sigma_{th}$ [W/m/K] |
|--------|------------------|----------------------|--------------------|--------------------|------------|---------------------|
| $^7$Li  | 0.534            | 1.88                 | 2.38x10$^{-4}$     | 1.01x10$^{-3}$     | 180        | 85                  |
| $^9$Be  | 1.848            | 1.85                 | 5.23x10$^{-5}$     | 4.95x10$^{-4}$     | 1280       | 200                 |

is a significant gain. In addition, to maximize $Y_n$, the target should have a high density ($\rho$) and a high purity since impurities will only lead to proton energy losses and so, to a lower reaction rate. To have a thermal neutron flux ($\phi_{th}$) which can compete with a reactor flux, equal to $10^5-10^6$ n/cm$^2$/s at the detector position, the proton beam intensity has to be around 20-100 mA. For $E_p=3$ MeV, this leads to a deposited power inside the target between 60 kW and 300 kW. Since the lithium fusion temperature ($T_f$) and thermal conductivity ($\sigma_{th}$) are very low (Table 1), operating a solid lithium target is difficult. It will only be possible with a large target, not suited to maximize neutron flux, or with a rotating target, with complex mechanics. Recently a liquid lithium target (LiLiT) has been successfully developed at the SARAF facility [5] and could be the key to operate a lithium target with high intensity beams. In the SONATE project, a solid beryllium target sustaining 50 kW is studied in priority. It has as advantages good thermal properties although $Y_n$ is four times less than with lithium (Table 1). To optimize the TMR configuration Monte-Carlo simulations are performed. Considering the low neutron yields, to reduce the computation time $^9$Be(p,n) reactions are not simulated. A neutron source term is built by hands based on experimental data (see Section 3). One drawback of using experimental data is that the emitted neutron energy and angle are not kinematically correlated. This adds to the simulations an uncertainty.

3 Experimental results and Geant4 simulation validation

To develop SONATE, different steps are being followed. In 2016, experiments where performed with a 10 W beam. These have firstly validated the experimental neutron angular distribution from Howard [6] and discarded Marion’s one [7] as shown in [8]. The former is thus used as input in the Geant4 simulation. The neutrons are preferentially emitted at 83 degrees w.r.t the beam axis. The neutron energy distribution is built from Howard data [6] and has an average energy of 640 keV. More details are given in [9]. With these ingredients, in the 2016 TMR configuration presented in Figure 1 (left), the simulation agrees with the experimental data with an accuracy better than 20 % as shown in Figure 2 (IPHI-2016 legends). The Geant4 simulation has also been successfully benchmarked against MCNP6 [10]. This validates the neutron transportation process in Geant4. In 2019, an experiment to test the thermal resistance and lifetime of a solid beryllium target with a 3 kW beam has been performed. The inherent difficulty in operating such a target with a high beam intensity (here 1 mA in average) is the proton blistersing effect. A high number of protons is implanted in the target and since the hydrogen solubility in Be and its diffusion coefficient are low, hydrogen atoms start to accumulate. Blisters appear and damage the target. With the beryllium target used for the experiment no blistering effect has been observed so far. The main difference regarding the 2016 and 2019 TMR experimental configurations are sketched in Figure 1 and summarized in Table 2. The Geant4 simulation also agrees with the data taken in the 2019 configuration inside their uncertainties as presented in Figure 2 (IPHI-2019 legends).
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Table 1: Neutronic and thermal properties for di-

| Power (W) | Tube diameter (mm) | Beam diameter (mm) | Material | Backing | Dimensions | Axial distance (cm) |
|-----------|--------------------|--------------------|----------|---------|------------|-------------------|
| 2016      | 10                  | 60                 | 16       | 0.5 mm  | 30x30x40 cm³ | 2.5               |
| 2019      | 3.0 kW              | 110                | 50       | 20 mm   | 50x50x60 cm³ | 4                 |

Between the 2016 and 2019 TMR configurations (Figure 2), at 2 meters, the simulated \( \phi_{th} \) dropped by a factor 1.7. This is mainly because the target thickness increases from 0.5 mm to 20 mm (Table 2). Indeed, in beryllium, a 640 keV (average initial neutron energy) neutron has a mean free path around 17 mm and the neutron mean free path averaged over the initial neutron energy spectrum is around 22 mm. These are similar to the 20 mm target thickness. Therefore, in average neutrons interact once before reaching the moderator. This spatially dilutes the neutrons leading to a less efficient neutron extraction. This target thickness was initially chosen for mechanical and thermal constraint reasons. This underlines the need to find the best compromise between mechanical, thermal, and neutronic constraints to maximize the neutron flux.

4 Moderator and reflector design

To provide the highest well collimated thermal flux outside the TMR assembly, the moderator and reflector materials and geometries, along with their coupling, have to be carefully investigated. During this optimization process, it has been systematically checked that the moderator/reflector material descriptions take into account the thermal scattering law (TSL) necessary to accurately transport thermal neutrons. In Geant4 (version 10.04.p02), TSL are taken from the ENDF/B-VII.0 data library [11].

Good moderator and reflector materials should both have a low capture cross-section and a high elastic cross-section. In addition, the moderator material should have a low nuclear mass to maximize the energy transfer between the neutron and the recoil nucleus. This is to minimize the number of collisions to thermalize the neutrons and so to minimize the neutron spatial dilution. Regarding the cross-sections in JEFF-3.3 data library, the best moderator materials are found to be polyethylene, light water, heavy water and graphite. These were expected based on previous reactor physics studies as summarized in [12, 13].

In the following, the TMR configuration is the 2016 one with the 2019 accelerator/target part (Figure 1). Starting from this configuration only the moderator material is changed. As expected, Figure 3 shows that for a given moderator volume and neutron extractor geometry, more thermal neutrons are extracted with hydrogenous materials. For heavy water and graphite larger dimensions are needed to thermalize the neutrons. This dilutes the spatial neutron distributions, which is not ideal to maximize \( \phi_{th} \) for collimated beam when only one extraction channel is used. The thermal fluxes obtained with polyethylene (CH₂) and light water (H₂O) are similar (Figure 3). However the thermal beam quality, defined by \( \phi_{th}/\phi_{tot.} \), is higher by a factor 1.5 with CH₂ compared to H₂O because ¹⁶O has a higher mass than ¹²C. To conclude, polyethylene is the best moderator to get a well collimated beam along with a high thermal neutron flux for one extraction channel and a low beam power. For higher beam power H₂O will be considered because of the CH₂ low thermal conductivity. For multiple beam extractions, D₂O could be also a good candidate.

Figure 3: Thermal neutron flux as a function of the distance (see Figure 1), for different moderator materials. The TMR configuration is the 2016 one with the 2019 accelerator/target part.

Figure 4: Thermal neutron flux as a function of the distance (see Figure 1), for different polyethylene moderator dimensions.
The impact of the polyethylene moderator dimensions on $\phi_a$ has then been studied. Figure 4 shows that the bigger the moderator size is, the higher $\phi_a$ is (neutrons are more thermalized). It also shows that starting from a 30x30x40 cm$^3$ moderator, $\phi_a$ starts to converge to an asymptotic value as the moderator size increases. It is then not necessary to increase the moderator volume above this volume.

To increase the thermal neutron flux outside the TMR assembly, a reflector can be placed around the moderator. Its role is to send back the neutrons to the moderator volume. The quality of a reflector is given by its albedo parameter ($\beta$) (details are given in [12, 13]). For a reflector size around three times the neutron diffusion length, $\beta$ converges toward an asymptotic value:

$$\beta_{\text{as}} = \frac{1 - 2 \sqrt{\Sigma_t/3 \Sigma_s}}{1 + 2 \sqrt{\Sigma_t/3 \Sigma_s}}$$  \hspace{1cm} \text{(1)}$$

with $\Sigma_t = \Sigma_{\text{tot}}, \Sigma_{\text{tot}}$, $\Sigma_s$, and $\Sigma_{\text{tot}}$ the transport, total, scattering and absorption macroscopic cross-sections. $\bar{\mu} = 2/3A$ is the average cosine of the scattering neutron angle in the laboratory frame. $A$ is the nuclear mass. From Equation (1), a good reflector is characterized by a low $\Sigma_s$ and a high $\Sigma_{\text{tot}}$. Heavy water, beryllium, graphite and lead (which is also a gamma shielding) can be used as a reflector. Heavy water is not considered anymore because its thermal neutron diffusion length is more than 1 meter, which will result in having a very large reflector. Taking into account these considerations, a 15x15x20 cm$^3$ polyethylene moderator is wrapped in a ~10 cm thick reflector. Figure 5 presents the impact of the reflector material on the thermal flux. At 1 meter, $\phi_a$ is higher by a factor 1.5 with beryllium and graphite compared with polyethylene and lead. Polyethylene is a poor reflector because of its high capture cross-section. Knowing that $\beta_{\text{as}}$ is reached for a material thickness equal to three times the thermal neutron diffusion length in the material, calculations showed that for beryllium and graphite, $\beta_{\text{as}}$ are respectively reached for a 60 cm and 165 cm reflector thickness. In increasing the beryllium reflector thickness from 10 cm to 60 cm $\phi_a$ is expected to increase approximately by a factor 2. This has to be confirmed by experiments.

5 Conclusion

Since the Orphée reactor will shutdown in 2019, an alternative solution to provide neutron beams is being investigated in CEA-Saclay. Taking advantage of the PHI high intensity proton accelerator, a compact accelerator driven neutron source is expected to be developed. The first experiments performed from 2016 to 2019 allowed to gain experience in operating solid beryllium targets with high intensity proton beams and to validate the Geant4 simulation software for two TMR configurations. The first optimization steps have shown that a TMR configuration with a small polyethylene moderator coupled to a big beryllium (or graphite) reflector will maximize the thermal neutron flux outside the TMR assembly for one neutron extractor. In 2020, the average beam intensity will be increased from 1 mA to 17 mA. The deposited power in the target will be around 50 kW. As the power density limit is set to 0.5 kW/cm$^2$, the target has to be redesigned. New studies are ongoing to find the best TMR configuration for this new target. In a near future, cold moderator and neutron extractor geometries will be investigated to increase the neutron flux and to develop new TMR configurations dedicated to other neutronic instruments. In a longer future, the accelerator will be upgraded to provide 20 MeV protons.

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