Abstract. We have investigated the nature of the magnetic white dwarf LP 790-29 = LHS 2293 by polarimetric monitoring, searching for short-term variability. No periodicity was found and we can exclude rotation periods between 4 sec and 1.5 hour with a high confidence. Maximum amplitudes of sinusoidal variations are $\Delta R < 0.009\,\text{mag}$ and $\Delta V_R < 0.7\%$ for a mean value of the $R$–band circular polarization of $V_R = +9.1 \pm 0.3\%$. Combined with earlier results by other authors, our observation suggests that LP 790-29 is, in fact, an extremely slowly rotating single white dwarf and not an unrecognized fast rotator and/or disguised cataclysmic variable.

Key words: stars: white dwarfs – stars: rotation – stars: individual: LP790-29
The high-field magnetic white dwarf LP790-29: not a fast rotator *

K. Beuermann and K. Reinsch

Universit¨ ats-Sternwarte, Geismarlandstr. 11, D-37083 G¨ ottingen, Germany

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1. Introduction

The large majority of white dwarfs are slow rotators with equatorial velocities $v_{\text{rot}} \sin i < 15 \text{ km s}^{-1}$ and rotational periods $P_{\text{rot}} \gtrsim 1 \text{ hr}$ (Heber et al. 1997, Koester et al. 1998). Much slower rotation is not detectable in non-magnetic white dwarfs, but easily measurable in magnetic ones by polarimetric monitoring (Schmidt & Norsworthy 1991, Berdyugin & Piïrola 1999).

Among the magnetic white dwarfs, there is a surprising dichotomy in the distribution of $P_{\text{rot}}$ for magnetic white dwarfs with all stars having either $P_{\text{rot}} < 20 \text{ d}$ or $P_{\text{rot}} > 100 \text{ yrs}$ (Schmidt & Norsworthy 1991). Some magnetic white dwarfs rotate surprisingly fast while others are apparently extremely slow. Among the fast ones are the DA white dwarf RE0317-853 with $B \simeq 500 \text{ MG}$, $T_{\text{eff}} \simeq 40000 \text{ K}$, and $P_{\text{rot}} = 12 \text{ min}$ (Barstow et al. 1995, Ferrario et al. 1997, Burleigh et al. 1999); the DA star PG1015+014 with $B \simeq 120 \text{ MG}$, $T_{\text{eff}} \simeq 14000 \text{ K}$, and $P_{\text{rot}} = 99 \text{ min}$ (Wickramasinghe & Cropper 1988, Schmidt & Norsworthy 1991), and the DAB white dwarf Feige 7 with $B \simeq 35 \text{ MG}$, $T_{\text{eff}} \simeq 20000 \text{ K}$, $P_{\text{rot}} = 2.2 \text{ hr}$ (Liebert et al. 1977, Achilles et al. 1992). Five systems seem to be very slow rotators with $P_{\text{rot}} > 100 \text{ yr}$ (Schmidt & Norsworthy 1991), among them the proven systems, GD229, G240-72 (Berdyugin & Piïrola 1999), and Grw+70°8247 (Friedrich & Jordan 2001), as well as a suspected one, LP790-29 (Liebert et al. 1978).

Slow rotation may be caused by coupling of angular momentum into the giant envelope of the progenitor star or the interstellar medium during later stages (Schmidt & Norsworthy 1991). Fast rotation may be achieved in a double degenerate low-mass donor (AM CVn binaries) transfer helium or carbon and typically end as CO white dwarfs, possibly with a substellar companion (Iben & Tutukov 1991). The hot white dwarf RE0317-853 has been suggested to be the result of a merger (Barstow et al. 1995, Ferrario et al. 1997) or a mass-transfer binary (Meyer & Meyer-Hofmeister 1999). The former appears more likely because the primary in RE0317-853 is hot with $T_{\text{eff}} \simeq 40000 \text{ K}$ and all white dwarfs in short-period cataclysmic variables are cool with $T_{\text{eff}} \simeq 9000 – 15000 \text{ K}$ (Gänsicke 2000). Although definite conclusions in any individual case may be problematic, the detection of rapid rotation of magnetic white dwarfs would clearly help to elucidate their evolutionary history. It also helps to understand the physical processes by which angular momentum is coupled into the environment of the star (Schmidt & Norsworthy 1991) and may help to define sources of gravitational wave radiation (Heyl 2000).

LP 790-29 = LHS 2293 was discovered by Liebert et al. (1978) and found to be a highly circularly polarized cool white dwarf which shows the Zeeman shifted C Swan bands. It has Stokes $V \simeq +8\%$ and $+9\%$ at wavelengths shortward of $4300 \text{ Å}$ and longward of $5500 \text{ Å}$, respectively, (Liebert et al. 1978, West 1989) which decreases to nil in between and to $2\%$ in the J-band. Bues (1999) refined its temperature to $T_{\text{eff}} \simeq 7800 \text{ K}$. The field strength was originally quoted as $B \sim 200 \text{ MG}$ (Liebert et al. 1978, Schmidt & Smith 1995), while Bues used $50 \text{ MG}$ in her spectral fitting, and Wickramasinghe & Ferrario (2000) quoted an uncertain $100 \text{ MG}$. LP 790-29 is also linearly polarized at the level of $\sim 1\%$ (West 1989). Polarimetric observations (Liebert et al. 1978, Robert & Moffat 1989, West 1989) have shown that the level of circular polarization has stayed constant over 10 years, excluding rotation periods in the range of $\sim 20 \text{ min} \lesssim P_{\text{rot}} \lesssim 100 \text{ yrs}$ (Schmidt & Norsworthy 1991). In a non-axisymmetric field geometry, the circular polarization will depend on rotation phase and a short rotation period should be readily detectable in time series of the circular polarization, provided the magnetic axis is inclined for more than a few degrees against the rotational axis and the latter does not point directly at the observer. In this communication, we report results of a

Send offprint requests to: beuermann@uni-sw.gwdg.de

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search for rapid rotation in LP 790-29, using photometric and polarimetric data taken in the Bessell R-band. Given the observed spectral dependence of Stokes V (Liebert et al. 1978, Schmidt et al. 1995), the R-band provides the best polarimetric signal of the standard photometric bands.

2. Observations and Results

We observed LP 790-29 on February 4, 2000, with the ESO 3.6m telescope equipped with the focal-reducer spectrograph and camera, EFOSC2, and a user-supplied superachromatic quarter-wave plate which was produced by Halle/Berlin and is of the same type as used in the ESO VLT FORS1 spectrograph. Flux and circular polarization were measured in the photometric Bessell R-band with a time resolution of \(\sim 30\) – \(40\) sec for a total of 1.5 h. Pairs of images with the retarder-plate position angle alternating between \(-45^\circ\) and \(+45^\circ\) were taken with exposure times of 2 s, 3 s, 5 s, and 10 s chosen in random order. The dead time between two exposures was 28–32 sec due to readout and instrument-setup times. Reading out single exposures guaranteed the best possible S/N ratio for Stokes \(V_R\). The \(R\)-band was chosen because of the higher level of circular polarization compared with that at the shorter wavelengths (Liebert et al. 1978). The measured values of \(V_R\) have been corrected for instrumental biases and linear-polarization cross-talk, \(-1.3\pm0.3\%\), which was determined to the first order from each pair of observations. The error in \(V_R\) is dominated by this systematic uncertainty, its statistical error is \(<0.1\%\).

Figure 1 shows the resulting time series. The \(R\)-band magnitude was measured against two comparison stars with similar brightness, USNO 0675 11099946 and USNO 0675 11100282, located 14 arcsec SW and 34 arcsec NE of LP 790-29. In order to search for periodicities, we computed the Analysis of Variance statistics implemented in the European Southern Observatory Munich Image Data Analysis System (MIDAS) software package (Schwarzenberg-Czerny 1989) for periods between the Nyquist limit of \(P_{rot} = 4\) sec and a maximum period \(P_{rot} = 1.5\) h (Figure 2). The expected value of this statistic for pure noise is \(\Theta_{AoV} = 1\). The critical value of the Fisher-Snedecor distribution for a 3-\(\sigma\) detection of a periodic signal in the data is \(\Theta_{crit} = 2.96\) for 9 and 140 degrees of freedom in the numerator and the denominator, respectively. The lack of a slope or curvature in the data excludes periods up to \(\sim 3\) h.

In order to estimate the upper limit on the amplitude of any periodicity in the range of 4 sec to 1.5 hour, we added artificial sinusoidal signals with various periods to the data and repeated the periodogram analysis. Amplitudes in excess of \(\Delta R = 0.006\)–0.009 mag and \(\Delta V_R = 0.5\)–0.7% can be excluded, with the lowest sensitivity corresponding to periods close to the local maxima of the spectral window function. Hence, we find that any coherent periodicity present must have a fractional modulation of Stokes \(V_R < 0.7\%\).

No significant photometric or polarimetric variability was detected. The periodogram in Fig. 2 is divided into two sections of which the left-hand one covers the more relevant longer periods that might be expected for a merger.
The right-hand panel demonstrates that no periodicity is found down to $P_{\text{rot}} = 4$ sec, the break-up period of a $1.2M_\odot$ white dwarf. For 6.950 trial periods each, in the photometric and in the polarimetric data, only a few frequency bins reach a significance above $3\sigma$, entirely consistent with expectation for the detection of spurious lines in a wide frequency band. Folding of the data on the frequencies with $\sim 3\sigma$ significance uncorrected for bandwidth demonstrates that none yields a sinusoidal modulation with an amplitude exceeding either 0.004 mag in $I_R$ or 0.4% in $V_R$. Since the addition of an artificial sinusoidal signal to the data produces a peak in the periodogram with FWHM $\sim 3$ pixels, the number of independent trial periods within the total bandwidth is reduced to $\sim 2.300$. Corrected for the bandwidth, the Fisher-Snedecor critical value for a $3\sigma$ detection of a periodic signal in the data then is $\Theta_{cr} = 4.39$. We conclude that there is no evidence for periodic variation in both the photometric and the polarimetric variation.

The mean of our $R$–band circular polarization values is $V_R = 9.1 \pm 0.3\%$. We summarize this and other published polarization measurements in the red part of the spectrum (quasi R–bands) in Tab. 1. The spectropolarimetric observations of Liebert et al. (1978) and Schmidt et al. (1995) have been averaged over pass bands as close as possible to the Kron-Cousins $R$-band. Since the circular polarization falls off from a maximum at 5750 Å towards longer wavelengths (Schmidt et al. 1995, their Fig. 2), the values quoted in Table 1 are not strictly comparable. We have tried to account for this uncertainty in assigning the errors. The quoted circular polarizations are consistent with each other, except for the low value measured from the polarization spectrum of Schmidt et al. (1995). Taken at face value, the data in Table 1 suggest that there is a variation...
over the time span of 23 years with a period of this order. The sparcity of the data calls, however, for more extensive polarimetric monitoring with a single instrument, either broad band polarimetry or preferably spectropolarimetry as described by Jordan & Friedrich (2001).

### 3. Discussion and Conclusion

We have performed a sensitive search for short periodicities in the \( R \)-band flux and circular polarization of the highly magnetic white dwarf LP790-29 which was previously thought to have a rotational period \( P_{\text{rot}} \sim 1 \text{ min} \). At an effective temperature of about 8000 K (Liebert et al. 1978, Bues 1999) the cooling age of LP 790-29 is \( t_{\text{cool}} \sim 2 \times 10^3 \text{ yrs} \) (Anselowitz et al. 1999), but may be shorter if it originated from the merger of a cool white dwarf with its companion (Segretain et al. 1997). Hence, even a cool white dwarf might still be a fast rotator, although over the time the original rotational velocity may have been reduced by magnetic braking.

Rapid rotation would be the signature of a white dwarf spun up in a cataclysmic variable of the AM Herculis type (Meyer & Meyer-Hofmeister 1999), of the AM Canes Venaticorum type (Iben & Tutukov 1991), or in a double degenerate merger (e.g. Segretain et al. 1997, Ferrario et al. 1997). A merger may rotate at the disruption limit. Segretain et al. argued, however, that the merger loses 90\% of the initial angular momentum by a strong wind, yielding an initial rotational period \( P_{\text{rot}} \sim 1 \text{ min} \). At an effective temperature of about 8000 K (Liebert et al. 1978, Bues 1999) the cooling age of LP790-29 is \( t_{\text{cool}} \sim 2 \times 10^3 \text{ yrs} \) (Anselowitz et al. 1999), but may be shorter if it originated from the merger of a cool white dwarf with its companion (Segretain et al. 1997). Hence, even a cool white dwarf might still be a fast rotator, although over the time the original rotational velocity may have been reduced by magnetic braking.

Our principal result is the absence of variability in LP790-29 with periods between 4 sec and about 3 hours and amplitudes \( \Delta R > 0.009 \text{ mag} \) and \( \Delta V_R > 0.7 \% \). This includes the absence of photometric variability of the type one might expect in a short period binary. Hence, even a cool white dwarf may have a rotational period longer than \( P_{\text{rot}} \sim 1 \text{ h} \), although the limit \( P_{\text{rot}} > 100 \text{ yrs} \) set by West (1989) and Schmidt & Norsworthy (1991) may be premature in view of the low level of the 1994 circular polarization by Schmidt et al. (1995). Nevertheless, these results suggest that LP 790-29 is, in fact, an exceedingly slow rotator. It seems worthwhile to follow up the possibility of a period of about a quarter of a century (see also the paper by Jordan & Friedrich 2001) by monitoring the level of circular polarization.

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