Ionospheric irregularities and scintillations: a direct comparison of in situ density observations with ground-based L-band receivers

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Abstract

Ionospheric irregularities can affect satellite communication and navigation by causing scintillations of radio signals. The scintillations are routinely measured using ground-based networks of receivers. This study presents observations of ionospheric irregularities by Langmuir probes on the Swarm satellites. They are compared with amplitude scintillation events recorded by the Global Positioning System-Scintillation Network and Decision Aid (GPS-SCINDA) receiver installed in Mbarara (Lat: 0.6°S, Lon: 30.8°E, Mag. lat: 10.2°S). The study covers the years from 2014 to 2018 when both data sets were available. It was found that the ground-based amplitude scintillations were enhanced when Swarm registered ionospheric irregularities for a large number of passes. The number of matching observations was greater for Swarm A and C which orbited at lower altitudes compared to Swarm B. However, some counterexamples, i.e., cases when in situ electron density fluctuations were not associated with any observed L-band amplitude scintillation and vice versa, were also found. Therefore, mismatches between observed irregularity structures and scintillations can occur just over a few minutes and within distances of a few tens of kilometers. The amplitude scintillation strength, characterized by the $S_4$ index was estimated from the electron density data using the well-known phase screen model for weak scattering. The derived amplitude scintillation was on average lower for Swarm B than for A and C and less in accordance with the observed range. Irregularities at an altitude of about 450 km contribute strongly to scintillations in the L-band, while irregularities at about 510-km altitude contribute significantly less. We infer that in situ density fluctuations observed on passes over or near Mbarara may be used to indicate the risk that ionospheric radio wave scintillations occur at that site.

Keywords: Equatorial ionosphere, Ionospheric irregularities, Amplitude scintillation

Introduction

Ionospheric irregularities occur near the geomagnetic equator after sunset and multiple studies have observed the extension of irregularity occurrence after midnight (e.g., Dao et al. 2011; Huang et al. 2014). They may be seen as field-aligned structures characterized by plasma density depletions, observed by satellites and radar backscatter in the topside ionosphere (Woodman and La Hoz 1976; McClure et al. 1977; Portillo et al. 2008). The spatial scales of equatorial plasma irregularities extend over at least four orders of magnitude (Lühr et al. 2014; Xiong et al. 2016; Rino et al. 2016). An assumed scenario is that larger-scale bubbles are created by a Rayleigh–Taylor instability (RTI) which starts from the bottom side of the ionosphere (Abdu 2005; Kelley 2009; Schunk and Nagy 2009; Portillo et al. 2008). The bubbles then rise and develop spatial gradients triggering further instabilities,
which result in a cascade from larger to small-scale structures typically for turbulent phenomena.

Ionospheric irregularities in the low latitudes may disturb the amplitude and phase of propagating radio signals, an effect called scintillation (Yeh and Liu 1982; Aarons 1982; Kintner et al. 2007). Scintillations depend on many factors, like the irregularity strength and layer thickness, signal frequency, local time, season, solar and magnetic activity, the satellite zenith angle and the angle between the ray path and the Earth's magnetic field (Aarons 1982; Wernik et al. 2007). Scintillation normally occurs when the Fresnel dimension of the propagating radio wave is of the order of irregularity scales in the ionosphere (Briggs and Parkin 1963). They are expected to originate mainly from medium-scale irregularities between the largest and smallest scales of a typical disturbed equatorial ionosphere, nearer to small scales. It is not clear to what extent the Fresnel theory which considers wave diffraction at hard boundaries gives robust estimates for the turbulent ionosphere. State-of-the-art numerical simulations have recently been able to resolve L1 Fresnel scale irregularities (Yokoyama 2017; Rino et al. 2018). Yokoyama (2017) recommended the need to directly resolve Fresnel scale irregularities so as to quantitatively evaluate scintillations in simulation studies.

Scintillations may affect the performance of satellite communication and navigation systems, for example, the Global Navigation Satellite System (GNSS). The intensity of the amplitude scintillation can be so strong that a receiver loses lock on an acquired signal. When several paths are affected, a navigation system may fail to give position (Wernik et al. 2007; Kintner et al. 2007; Paznukhov et al. 2012). The ability to evaluate the risk of occurrence of scintillations is useful for accurate navigation applications. Therefore, the study of scintillations and ionospheric irregularities is an interesting scientific topic for plasma turbulence, radio wave propagation and also important when developing applications and systems.

Previous studies have shown that scintillations recorded on the ground are associated with ionospheric irregularities encountered in situ (Roddy et al. 2010; Nishioka et al. 2011; Akala et al. 2017; Andima et al. 2018). Roddy et al. (2010) and Nishioka et al. (2011) compared ionospheric irregularity observations made by Communication Navigation Outage Forecasting System (C/NOFS) and 250-MHz scintillations over Christmas Island (3°N, 157°W). Akala et al. (2017) compared SCINDA GNSS data at Addis Ababa (Lat 9.03°N, Lon 38.77°E, Mag. lat 0.18°N) with C/NOFS data. Andima et al. (2018) characterized the low latitude ion density irregularities measured by C/NOFS from 2011 to 2013. This was supported by ground-based data from the SCINDA receivers at Makerere (Lon 32.6°E, Lat 0.3°N, and Dip. lat −9.3°N) and Nairobi (Lon 36.8°E, Lat −1.3°N, and Dip. lat −10.8 °N). However, most of these studies used ionospheric irregularity observations made by C/NOFS. The European Space Agency’s (ESA) Swarm satellite mission has provided an opportunity to compare scintillation occurrences and ionosphere irregularities from a meridional point of view. Olwendo et al. (2019) specifically compared irregularity structures recorded by Swarm at 2 Hz frequency with amplitude scintillation over Pwani (Lon: 39.78°E, Lat: 3.24°S). The Swarm satellites can measure electron density using a faceplate at an even greater frequency of 16 Hz which enables smaller-scale structures to be identified in electron density (Nishioka et al. 2011). The high-resolution density estimates sampled at 16 Hz correspond to a spatial scale of about 500 m, which is of the order of the Fresnel scale relevant for scintillations in the L-band at GNSS frequencies.

Furthermore, to overcome the limited spatial coverage of ground-based receivers, efforts have been made to derive scintillation for prediction/monitoring purposes from in situ electron density measurements as an alternative. One approach involves the derivation of scintillation from in situ electron density measurements using the weak phase screen model (Basu et al. 1976; Rino 1979; Crane 1977; Wernik et al. 2007; Patel et al. 2011; Liu et al. 2012; Alfonsi et al. 2017). To test the ionospheric scintillation model, coincident observations of electron density in the ionosphere and scintillation signals on the ground is necessary (Liu et al. 2012). The Swarm high-resolution faceplate electron density measurements were recommended by Alfonsi et al. (2017) as input to the Rino’s power law phase screen model to derive scintillation of amplitude.

Limitations in spacecraft orbital parameters and orientation have meant that C/NOFS, which orbited at a low inclination angle of about 13°, sampled ionospheric irregularities mainly in the zonal direction. As far as we know, no attempt has yet been made in calculating expected scintillation amplitudes from Swarm 16-Hz density data. Therefore, in the present study, we compared measurements of near Fresnel scale resolved irregularities recorded in situ by Swarm with observed scintillations at L1 on the ground to demonstrate whether observed 500-m scale irregularities correspond to detected scintillations. In addition, we also modeled and predicted the expected level of scintillations from the observed density variation using the power–law phase screen model. The 16-Hz Swarm electron density and GPS-SCINDA scintillation measurements were used for the period from 2014 to 2018 when both data sets were available.
Data and methods

Data
The Swarm mission is made up of three identical satellites, i.e., Swarm A, B, and C, each with payloads designed to measure the Earth’s magnetic field and space environment (Friis-Christensen et al. 2006). After commissioning and orbit consolidation was completed in mid-March 2014, a polar orbit (inclination ≈ 87.5°), with Swarm A and C orbiting side by side at a longitudinal separation of about 1.4° was reached. By the end of 2018, Swarm A and C altitude, near the equator and that of Swarm B was about 512 km. In this study, the Swarm faceplate electron density (Ne) measurements at a frequency of 16 Hz were used to identify irregularity structures in the topside ionosphere and these are readily available at http://earth.esa.int/swarm. The Swarm pass segments with magnetic latitudes between −30° and +30° were considered. Only faceplate Ne data acquired during orbits that passed within ±5° longitude of the ground site were considered for detailed comparison, similar to Roddy et al. (2010) and Sarkar and Gwal (2014). The occurrence of amplitude scintillation was determined using the GPS-SCINDA receiver installed at Mbarara, roughly near the location of the EIA southern crest. The GPS-SCINDA samples raw data at a frequency of 50 Hz. It records the carrier-to-noise ratio (C/No) intensity and phase of GPS L1 and L2 signals, as well as the azimuth and elevation of the visible satellites (Akala et al. 2012). It is a real-time GPS data acquisition and ionospheric analysis system which computes amplitude scintillation (S4) and Total Electron Content (TEC) using the full temporal resolution of the receiver (Groves et al. 1997).

Data processing

Ionospheric irregularity identification from swarm
To identify ionospheric irregularities, a similar method used by Huang et al. (2014) and Aol et al. (2020) was applied. The 16-Hz Ne measurements were passed through a 2-s running mean filter corresponding to a wavelength of about 15 km along the satellite track. The running mean Ne determined over a 2-s interval was subtracted from the original observations to obtain the residuals dNe = Ne − Nε. The standard deviation of the residuals std (dNe) was then used to represent the density perturbation. We mainly focused on small-scale equatorial plasma structures. A threshold value of std (dNe) = 1 × 10^{10} m^{-3}, similar to that by Huang et al. (2014) was adopted.

Amplitude scintillation from GPS-SCINDA
The GPS-SCINDA determines S4 index using the equation: S4² = (⟨I²⟩ − ⟨I⟩²)/⟨I⟩², where I represents the signal intensity. The S4 index is defined as the normalized ratio of the standard deviation of signal intensity fluctuations to the mean signal intensity (Groves et al. 1997; Kil et al. 2002; Nishioka et al. 2011). The S4 index is computed at 60 s intervals. Processing of the GPS-SCINDA data was done using the GPS-TEC application software (Seemala and Valladares 2011). In this study, amplitude scintillations with S4 > 0.3 were considered as relevant scintillation events, similar to Akala et al. (2011), Alfonsi et al. (2013), and Wang et al. (2014). Only S4 data corresponding to ray path elevations above 30° was considered to avoid the multipath effects. For comparison with Swarm data, the IPPs of the GPS satellites were extrapolated to Swarm A and C heights to judge as accurately as possible how close in space and time to the GPS signal path a Swarm satellite was observing the plasma structures. Scintillation data at times when Swarm was within ±5° longitudinal range from the longitude of the ground site were considered. Scintillation data that were close to Swarm passes in terms of latitude and longitude were selected.

Amplitude scintillation from the phase screen model
To derive S4 from the 16-Hz Ne data, using Rino’s phase screen model, the data were first split into 8-s long segments similar to Wernik et al. (2007) and Liu et al. (2012). For each segment, a detrended Ne represented by dNe was determined by subtracting a linear fit to the data segment. A Hanning window was applied to each segment to prevent energy leakage (Liu et al. 2012). Spectral analysis was then carried out for each segment using the Fast Fourier Transform algorithm to obtain the Power Spectral Density (PSD). The power spectra of 8-s long 16-Hz data segments were calculated, corresponding to scale sizes from about 300 m to 1 km along the satellite track. The one-dimensional spectral index p was calculated using log–log least-squares fit to the PSD. The modeled amplitude scintillation was obtained from the equation (Rino 1979):

\[ S4^2 = (r_e k)^2 L \sec \theta C Z^{\nu - 0.5} \left\{ \frac{\Gamma[(2.5 - \nu)/2]}{2\pi^{0.5} \Gamma[(\nu + 0.5)/2](\nu - 0.5)} \right\} F, \]
where \( C_s = 8\pi^{3/2} (\Delta N_e^2) q_0^2 \nu^{-2} \Gamma(v + 0.5) / \Gamma(v - 0.5) \) is the turbulence strength (Rumsey 1975), the Fresnel zone parameter \( Z = \lambda Z_R = \pi \theta / 4, \lambda Z_R = z_s / z + z_s, \) \( z \) is the distance to the phase screen, \( z_s \) is the distance to the GPS satellite, \( r_e = 2.8 \times 10^{-15} \) m is the classical electron radius, \( \theta \) is the zenith angle of the signal path, and \( L \) is the irregular layer thickness. We have assumed \( L = 200 \) km (Basu et al. 1976; Patel et al. 2011). \( \lambda \) is the wavelength at GPS L1 frequency, \( q_0 = 2\pi / L_o \) is the outer-scale cut-off number, \( L_o \) is the outer scale of the irregularity, \( \langle \Delta N_e^2 \rangle \) is the variance of \( N_e \) at the satellite location and the angle brackets denote ensemble averaging. Rino (1979) noted that, the in situ spectral index \( p \) is one less than the three-dimensional phase spectral index \( \nu \) used in Eq. 1. We, therefore, used the formula \( p = 2\nu - 1 \) to obtain \( \nu \) from the spectra of the in situ \( N_e \) fluctuations (Rino 1979). 

\( F \) in Equation 1 is the geometry-dependent Fresnel filter factor defined in Equation (34) of Rino (1979). For isotropic irregularities \( F = 1 \) (Wernik et al. 2007). A detailed derivation of the modeled \( S_4 \) in Eq. 1 can be found in Patel et al. (2011) and references therein.

**Results and discussion**

**Correspondence between occurrence of ionospheric irregularities and scintillations**

This section presents the occurrence of ionospheric irregularities observed by Swarm A, B, and C in comparison with simultaneous measurements of \( S_4 \) index values on the ground. Figure 1 shows examples of \( N_e \) variations for Swarm A, B, and C passes close to Mbarara longitude (\( \approx 30.8^\circ \)). The results include variations in \( N_e \) in logarithmic scale and standard deviation of \( dN_e \) (std\((dN_e)) \) in panels (a) and (b), respectively. Multi-peak variations in \( N_e \) were observed in panel (a) of Fig. 1 due to the sufficient temporal resolution of the 16-Hz measurements. From panel (b) of Fig. 1, it is observed that std\((dN_e)) \) clearly captures the fluctuations in \( N_e \) between about \( \pm 15^\circ \) quasi-dipole latitude (QLat). The depletions are certainly associated with nighttime ionospheric irregularities (Kelley 2009; Pi et al. 1997; Kintner et al. 2007). Radio waves at HF, VHF and higher frequencies are particularly affected by the plasma turbulence generated through the RTI (Yeh and Liu 1982).

An example of coincident observations for Swarm A and C and GPS-SCINDA on 2015-03-09 is presented in Fig. 2. Panel (i)(a) of Fig. 2 presents the latitudinal variation of \( N_e \) and \( S_4 \) index values, while panel (i)(b) presents the longitudinal variation. The longitude plot serves to identify whether a Swarm satellite pass was close to a signal path where \( S_4 \) index values could be determined. The \( S_4 \) for all the GPS IPPs in view are shown with black dots, while the colored IPPs (see legend) correspond to the period when Swarm was in the view near Mbarara. From panel (i), multiple bubble-like \( N_e \) bite-outs are seen along especially Swarm A orbit track. The large-scale longitudinal bubble structure is sometimes observed with the two Swarm satellites (Xiong et al. 2016), but for small-scale ionospheric irregularities the separation is too large for a significant correlation between A and C as seen in Fig. 2. Panel (ii) of Fig. 2 shows the Swarm A and C ground track on 2015-03-09 and the ground track of the pseudo-random noise (PRN) number in view assuming

![Fig. 1](image-url)
a piercing height of 460 km corresponding to about the Swarm A and C orbit. Most of the GPS satellites in view showed $S_4 > 0.3$ around the same time when density fluctuations were obtained from Swarm A and C. The IPPs of the GPS signal of PRN 29 were the closest to the irregularity structure observed by Swarm A and C, and the signal showed $S_4$ values $> 0.3$. In this event, particularly ionospheric irregularities around a high background density seem to be associated with scintillations observed at the ground in Mbarara. The spatial distribution in panel (ii) of Fig. 2 shows that the ionospheric irregularities encountered by Swarm A and C and the GPS satellites are clustered at a specific volume around the ground GPS-SCINDA. Therefore, coincidental detection of the high $S_4$ and std($dN_e$) values is a common volume of ionospheric perturbations is an indication that the std($dN_e$) may be used as indicator of potential scintillation occurrence.

An Example for Swarm B, which orbits at a higher altitude than the lower pair, is presented in Fig. 3: Swarm B encountered ionospheric irregularities, seen as $N_e$ depletions in panel (i), along its track as it orbited near Mbarara. The relative fluctuations along track B on 2015-04-02 follow the typical double-peaked background density across the equator. The obtained std($dN_e$) values were up to $0.5 \times 10^{11}$ m$^{-3}$ (shown in Fig. 1) which is well above our threshold. $S_4 > 0.3$, up to 0.45 was recorded by the GPS satellite of PRN 24 in view as seen from panel(i) of Fig. 3. Also, this Swarm B satellite pass on 2015-04-02 demonstrates a good agreement between the increase in the level of $S_4$ and the strength of $N_e$ perturbations for GPS satellite of PRN 24. To examine the spatial distribution of Swarm B track and GPS observations in relation to the location of the ground-based GPS-SCINDA, the ground tracks of Swarm B and PRNs are presented in Fig. 3 panel (ii). From the spatial distribution observed in panel (ii) of Fig. 3, Swarm B and the GPS satellite of PRN 24 in view seem to be sampling overlapping regions of large and perturbed areas, similar to that presented in Figure 2 for Swarm A and C. Therefore, the coincidental observation of relatively high $S_4$ and std($dN_e$) at Swarm B altitudes are causally related in the event presented in Fig. 2.

In a substantial fraction of the Swarm satellite passes, there was no one-to-one correspondence between the occurrence of ionospheric irregularities
and scintillation. An example is presented in Fig. 4 for Swarm A and C on 2014-09-07. Figure 4 shows a case where ionospheric irregularities were observed on Swarm A and C pass with only a mild effect on $S_4$ index for GPS satellites of PRN 4, 7, 8, 28, and 30 at the time of Swarm A and C passes. At least for this event, the ionospheric irregularities encountered by Swarm A and C seem to be intermittent and patchy such that the GPS signals became only moderately perturbed. Also, for the Swarm passes presented in Fig. 4, the background density was lower compared to the events shown in Figs. 2 and 3. Nishioka et al. (2011) also observed that intense in situ ionospheric irregularities recorded using C/NOFS PLP were accompanied by only mild VHF
scintillation on 2008-10-10 Ancon near Jicamarca. Nishioka et al. (2011) concluded that the observed contradiction was due to the dynamics of the EIA, whose development is affected by the duration of the Pre-Reversal Enhancement (PRE). Our study continues by comparing the statistical occurrence of ionospheric irregularities and amplitude scintillation and the results are presented in the following section.

Statistics of occurrence of ionospheric irregularities and amplitude scintillation
Summary plots such as those presented in Figs. 2, 3, and 4 were generated for Swarm nighttime orbits for which $S_4$ data were available over Mbarara. The outputs of the summary plots could be categorized into four cases, similar to those presented by Roddy et al. (2010), considering the presence (or not) of ionospheric irregularities and scintillation. These cases are presented in Fig. 5. In about 41.9% of the cases for Swarm A, 45.3% for Swarm C, and 39.0% for Swarm B, no irregularity structures were detected by the Swarm satellites and no L-band scintillations were recorded by the GPS-SCINDA receiver. In about 30.3% of the cases for Swarm A, 27.3% for Swarm C, and 17.3% for Swarm B, irregularity structures were detected and GPS L-band scintillations were also seen. Roddy et al. (2010) and Sarkar and Gwal (2014) also observed a high percentage of agreement considering the presence (absence) of situ ionospheric irregularities and ground-based scintillation. These two categories (i.e., irregularities and scintillation and no irregularities and no scintillation) are consistent with a hypothesis, that irregularities seen by Swarm are associated with ground-based L-band scintillations when the signal path is close to the irregularity detection and scintillations are absent or small when no irregularities are seen. However, there were also “negative” detections (irregularities but no scintillation and vice versa). An explanation may be that some scintillations were produced by irregularity structures at lower altitudes which may not have extended to Swarm altitudes (Burke et al. 2003; Roddy et al. 2010). Also, a GPS signal may evade passing ionospheric irregularities due to high intermittency and spatial inhomogeneity of the ionospheric turbulence.

In terms of a possible risk assessment, with Swarm A data as input (confirm Fig. 5), in 42.6% of the passes, irregularities were detected. At 71.1% of these detections the threshold for scintillations was also exceeded at the receiver in Mbarara, while 28.9% of detections were “false alarms”, i.e., not accompanied by scintillations. 57.4% of the passes had no clear indications of irregularities in the Swarm data. In 72.9% of these passes with no irregularities also no scintillations occurred, but 27.1% did have scintillations at Mbarara. For Swarm C, in 37.2% of the passes, irregularities were detected. At 73.4% of these detections, scintillations was also exceeded at the receiver in Mbarara, while 26.6% of the detections were “false alarms”, i.e., not accompanied by scintillations. 62.7% of the passes had no clear indications of irregularities in the Swarm data. In 72.2% of these passes with no irregularities also no scintillations occurred, but 27.8% did have scintillations at Mbarara. For Swarm B, in 35.8% of the passes irregularities were detected. At 51.4% of these detections the scintillations was also exceeded at the receiver in Mbarara, while 48.6% of detections were “false alarms”, i.e., not accompanied by scintillations. 64.3% of the passes had no clear indications of irregularities in the Swarm data. In 60.7% of these passes with no irregularities also no scintillations occurred, but 39.3% did have scintillations at Mbarara. Therefore, the 16-Hz Swarm data may indicate an unbiased statistical chance/risk that radio wave scintillations occur at the ground.

Several parameters, including local time, latitude, longitude and season, affect the development of ionosphere disturbances. Therefore, it was interesting to also compare seasonal and local time dependence of GPS amplitude scintillation to that of ionospheric irregularities observed over Mbarara by Swarm. Figure 6 presents the percentage occurrence of ionospheric irregularities and $S_4$ as a function of local time grouped into bins of 1 h. Even when the events included in Fig. 5 were filtered according to spatial and temporal proximity, the alignment still allowed for differences by a few tens of km and minutes between the Swarm orbit and GPS path. The percentage occurrence of ionospheric irregularities observed by the three Swarm satellites increased
Figure 7 presents the seasonal dependence of occurrence of ionospheric irregularities observed by Swarm in panels (a–c) and $S_4$ recorded by the GPS-SCINDA in panel (d). To generate Fig. 7, for each month (y-axis), the std($dN_e$) and $S_4$ were averaged over 0.1 hr local time. The highest values of std($dN_e$) and $S_4$ were observed during the equinoxes and December solstice. Overall, during June solstice, small values of std($dN_e$) and $S_4$ were observed. Similar seasonal distribution of ionospheric irregularities and amplitude scintillation was reported by Akala et al. (2011) and Paznukhov et al. (2012). The seasonal dependence of ionospheric irregularities as well as amplitude scintillation is adequately explained by geomagnetic arguments, i.e., the alignment of the solar terminator and local magnetic meridian (Tsunoda 1985). Ionospheric irregularities and scintillations are most commonly observed during seasons when the solar terminator is best aligned with the geomagnetic meridian (Tsunoda 1985). Figure 7 also offers further proof for the connection of the GPS-SCINDA recorded amplitude scintillation occurrences in our research region with Swarm’s in situ observed ionospheric irregularities.

Furthermore, we analyzed the relationship between $S_4$ and $N_e$ perturbations associated with plasma bubbles. The results are presented in Fig. 8, combining the data for Swarm A, C, and B (see legend). The results show a linear dependence between $S_4$ and std($dN_e$) with moderate positive correlation coefficients (see legend) for the years from 2014 to 2018, although very few data points were available for this analysis. Due to scarcity of Swarm orbital tracks which coincided with the location of the ground GPS-SCINDA during the period of study, we were limited in the number of events where we could directly associate in situ ionospheric irregularity events with the ground-based $S_4$ measurements. Therefore, the moderate positive correlation coefficients between $S_4$ and std($dN_e$) can be attributed to this relatively low number of passes and a typically high variability of irregularity strength, indicated by std($dN_e$), along satellite tracks. A future larger number of suitable passes will allow to empirically study the relation between ionospheric irregularities and scintillation strengths in more detail. Nevertheless, we modeled the amplitude scintillations from the available Swarm 16-Hz $N_e$ data. The results are presented in the following section.

Amplitude scintillation from the weak phase screen model

In this section, we present model results of amplitude scintillation from Swarm $N_e$ measurements. Conducting spectral analysis of the high-resolution $N_e$ data from Swarm was the first step to estimate the resulting amplitude scintillation. Figure 9 presents arbitrary example
Fig. 7 Month and local time (LT) variations of std(dNe) (panels (a–c) and $S_4$ index (panel (d)) for the years from 2014 to 2018.

Fig. 8 Relation between the GPS-SCINDA scintillation intensity and Swarm $N_e$ perturbations for the years from 2014 to 2018.
outputs obtained after the spectral analysis procedure. Panels (i) and (ii) of Fig. 9 show the 8-s-long segments of $N_e$ and their corresponding PSDs, respectively, for Swarm A and C on 2015-07-03 and B on 2015-07-25. The PSD is plotted as a function of the wave number $k$ ($k = \omega/v$, $v = 7.5$ km s$^{-1}$ is the speed of the Swarm satellites, $\omega$ is the angular frequency) in panel (ii) of Fig. 9. A key parameter to be considered when determining the scintillation level is the spectral index $p$ (Wernik et al. 2007). Therefore, the value of the spectral slope or spectral index $p$ estimated by linear least square fitting are indicated and the uncertainty in the slopes is provided in parenthesis for each slope in panel (ii) of Fig. 9. It is important to recall that the spectrum obtained from in situ $N_e$ data is a one-dimensional scan of the power spectrum of $dN_e$. From panel (ii) of Fig. 9, if we ignore the somewhat flat portion of the PSD at small $k$ values, the regions in the spectra tend to follow a power–law, i.e., $P(k) \propto k^{-p}$. The power–law relation has also been observed for data from other satellite missions (Basu et al. 1980; Wernik et al. 2007; Su et al. 2001; Rodrigues et al. 2009; Lühr et al. 2014).

The spectral indices $p$ were determined as the slopes of straight lines fitted to log–log power spectra of the $N_e$ irregularities observed for the years from 2014 to 2018. These are presented in Fig. 10 which reveals that most values of the spectral index, $p$, is between about 0.5 and 2.5. A range of $p = 2.0 - 2.2$ was obtained by Patel et al. (2011) for the Stretched Rohini Satellite Series (SROSS-C2) and $p = 2.0$ had the highest occurrence rate. Lühr et al. (2014) obtained $p$ from the high-resolution Challenging Minisatellite Payload (CHAMP) magnetic field measurements over the frequency range 1–25 Hz and they observed spectral indices between 1.4 and 2.6 with a peak occurrence rate around 1.9, while Wernik et al. (2003) obtained $p$ for the plasma density data recorded by the Dynamic Explorer-2 (DE-2) satellite within the frequency range of 1–20 Hz and a peak in percentage
occurrence of $p$ was within the range of 1.0 and 2.5. In general, the highest percentage occurrence of $p$ obtained in Fig. 10 is within the range observed in previous studies. Rino (1979) and Singh et al. (1997) used spectral indices 2.5, 2.6, and 2.66 as inputs to the phase screen model. Crane (1977) reported $p = 3$, which seems to be slightly higher than the measured values. Alfonsi et al. (2017) took a value of spectral index $= 1.3$. From the results presented by Rino (1979), the 1-GHz amplitude scintillation data fit the model computations to a much better extent using $\nu = 1.25$. Therefore, in this study we adopted a value of $p = 1.5$ as input to the model which leads to a value of $\nu = 1.25$ (e.g., Rino 1979; Singh et al. 1997).

Vats et al. (1981) considered the outer scale length, $L_o$ to be 50 km for F-region irregularities, whereas Rino (1979) took $L_o$ to be 200 km. Patel et al. (2011) varied $L_o$ between 50 km and 200 km and they presented results using $L_o = 175$ km as the most feasible value. In the present modeling, we also adopted a value of $L_o = 175$ km. Examples of modeled outputs for Swarm A, B, and C along the satellite tracks are presented in Fig. 11. The figure presents the $N_e$ fluctuations along the satellite track,
corresponding \( \text{std}(dN_e) \), the turbulence strength \( C_s \), and the modeled \( S_4 \), in that order. The \( S_4 \) index frequently varied from about 0 to about 1.0 for Swarm A satellite pass on 2015-03-09 and from about 0 to 0.4 for Swarm C satellite pass on 2015-03-09. The calculated \( S_4 \) for Swarm A and C is in a reasonable agreement with the observed increase of \( S_4 \) for GPS PRN 29 (see Fig. 2). The weak scattering model seems to produce the amplitude scintillations at about the observed level when using the Swarm A and C \( N_e \) measurements. The calculated \( S_4 \) for Swarm B was between 0 and 1.0 which is in agreement with the observed \( S_4 \) range (see Fig. 3), particularly for GPS PRN 29.

We analyzed the relationship between the \( S_4 \) index values measured by the GPS-SCINDA and modeled \( S_4 \) index values for the years from 2014 to 2018 and the results are presented in Fig. 12. Swarm satellite passes within a longitude range of \( \pm 5^\circ \) of the ground site were considered and the \( S_4 \) index values recorded by the GPS-SCINDA were determined for each PRN which was in view during a Swarm satellite pass within a longitude range of \( \pm 1^\circ \) of the Swarm longitude. To distinguish scintillation phenomena from other noise effects, and to remain within the weak scatter assumption, only observed \( S_4 \) index values >0.3 were used. Generally, the measured and modeled \( S_4 \) show moderate positive correlation for Swarm A and C. The results presented in Fig. 12 show that the number of data points available was limited and the moderate positive correlation obtained may be due to low counting statistics. In general, solar cycle 24 was characterized by very low solar activity compared to cycles that preceded it (Basu 2013). Therefore, we managed to identify very few events where Swarm coincided with GPS satellites over Mbarara. The relatively larger scatter of modeled versus observed \( S_4 \) is consistent with discrepancies that were found in the analysis of the events above and explained by a high degree of intermittency and spatial localization of small-scale irregularities within larger-scale structures such as equatorial plasma bubbles (EPBs). Small-scale irregularities conglomerate in such an EPB environment, but do not fill it. On the other hand, Swarm B shows no significant or even a negative correlation between measured and modeled \( S_4 \).

Nishioka et al. (2011) described how VHF scintillations seemed to originate from bottom-side irregularities that remained undetected by the C/NOFS satellite at an orbit of as low as 400 km. Here, such bottom-side irregularities seem to have a negligible statistical effect for the case of Swarm A and C (where both too high and too low modeled \( S_4 \) occur). But for Swarm B, the effects of ionospheric irregularities confined to below 500 km seem important. Also, the Swarm B data cover the post-sunset sector at increasingly later LT than A and C (Zakharenkova et al. 2016) which might contribute to the statistical non-agreement between modeled and observed \( S_4 \). Comparable modeling for the high latitudes had been undertaken using 64-Hz Retarding Potential Analyzer (RPA) data from the DE-2 by Wernik et al. (2007). In the study by Wernik et al. (2007), the modeled \( S_4 \) was roughly one order of magnitude too low (their Figures 5, 13-16). Furthermore, the DE-2 orbit was on average higher than Swarm A and C and this might have contributed to the low modeled L-band \( S_4 \) by Wernik et al. (2007). Patel et al. (2011) derived
appropriate relations for $S_4$, and phase variance from the weak scatter theory using Stretched Rohini Satellite Series (SROSS-C2) satellite at VHF. Their modeled results reproduced successfully the observed VHF scintillation intensity monthly and seasonal variations. Alfonsi et al. (2017) extended the phase screen model to strong scattering and modeled $S_4$ at L-band frequency around the equatorial latitudes using ion density measurements collected by the DE-2 RPA. The observed density variations were scaled to a model peak density using the IRI model by Alfonsi et al. (2017).

In addition to the spatial distribution, other factors related to the modeling may have caused discrepancies between modeled and measured $S_4$ in the results presented in Fig. 12. One factor is the assumption that the ionospheric irregularities are isotropic which leads to an overestimate of turbulence strength (Wernik et al. 2007). Furthermore, in the Rino’s phase screen model, the outer scale of turbulence is assumed to be very high, meaning that the outer-scale wave number is far less than the wave disturbances adding to the amplitude scintillation (Alfonsi et al. 2017). Liu et al. (2012) estimated the outer scale and found that it varies from 1 km to several hundreds of kilometers, with the distribution having maximum at scales less than 25 km for dip latitude $< 20^\circ$ and at scales 25–50 km for dip latitudes $> 20^\circ$. Thus, in most cases, the outer scale is much larger than the Fresnel’s scale for frequencies of interest and any effects of the outer scale on the intensity scintillation can be neglected.

It is well accepted that the scale sizes of the field-aligned irregularities that are primarily responsible for L-band scintillation are about 400 m. According to the Nyquist theorem, to correctly represent these scale sizes, the medium would have to be sampled at every 200 m or less. Considering the Swarm satellite velocity of 7.5 km s$^{-1}$, this would correspond to the sampling frequency of 37.5 Hz. A comparison of calculated $S_4$ based on the simple phase screen model and the observed $S_4$ shows a large scatter, but for the lower Swarm A and C no significant systematic bias. A scaling of the model $S_4$ as it was done for DE-2 seems not needed, probably owing to the lower altitude.

**Conclusions**

In this study, measurements of $N_e$ variations by the Swarm satellites were compared with L-band scintillation data at the ground over Mbarara. The findings of this study are summarized below:

1. For selected events, the observed ionospheric irregularities along Swarm tracks could be identified with moderate increases in $S_4$ at the ground. In some cases, ionospheric irregularities were not associated with significant scintillation or amplitude scintillations were seen without Swarm recording sufficient density variations. Similar disagreements between ground-based scintillation observation and in situ ionospheric irregularities were reported by Roddy et al. (2010) and Sarkar and Gwal (2014). This may be explained by small-scale irregularities arising and disappearing within a few minutes, and having spatial extents of only tens of kilometers (Basu et al. 1978; Saito et al. 2008). Some observed disagreements may also be due to scintillation causing irregularities confined to lower altitudes (Roddy et al. 2010; Burke et al. 2003) and they are therefore not encountered by Swarm at 445 km altitude and above.

2. Statistically, during satellite passes there is a reasonable correspondence between in situ-measured irregularities and L-band scintillations. The distributions of ionospheric irregularities and scintillations over both LT and seasons show similar patterns. This is in agreement with the results obtained in earlier studies (Burke et al. 2003; Andima et al. 2018). Therefore, the similar LT and seasonal trend obtained provide statistical evidence for the association of amplitude scintillation events recorded by the GPS-SCINDA in our study area to ionospheric irregularities recorded in situ by Swarm faceplate.

3. Using the weak phase screen model, the turbulence strength $C_s$, 1-D spectral index $p$ and $S_4$ were derived from the Swarm 16-Hz $N_e$ data. In general, there were discrepancies between the derived and observed $S_4$ for Swarm. The discrepancies may have risen because of poor spatial distribution and other factors related to the modeling. Despite the limitations associated with the phase screen model in Rino’s formulation, the results presented here act as an initial step to describe weak scintillation at Swarm satellite altitudes and therefore encouraging further studies. Also, while the weak phase screen model gives nearly the correct range of $S_4$ values when using Swarm A and C faceplate $N_e$ data, a more refined two-component model could be tried in future works, when a larger data set is available, and the relative directions with respect to the magnetic field is taken into account.

While there are no concrete plans to obtain Swarm data in near real-time, this might be a possibility for other missions. In the mean time, the statistics of irregularities seen by Swarm have been analyzed (Wan et al. 2018; Aol et al. 2020), which could be a basis of empirical models for irregularities. With limitations, which are investigated and pointed out in this work a global risk for L-band scintillations can be inferred from the Swarm data,
complementing existing scintillation models which are based on ground-based observations. The Swarm data set has good global coverage and therefore, the potential to study irregularities and scintillations over areas where no or very few ground scintillation data are available, over the oceans and parts of central Africa.

Abbreviations
C/No: Carrier-to-noise ratio; C/NOFS: Communication Navigation Outage Forecasting System; CHAMP: Challenging Minisatellite Payload; DE-2: Dynamic Explorer-2; EPBs: Equatorial plasma bubbles; EIA: Equatorial ionization anomaly; ESA: European Space Agency; GNSS: Global Navigation Satellite System; GPS-SCINDA: Global Positioning System-Scintillation Network and Decision Aid; IPPS: Ionospheric pierce points; PRE: Pre-reversal enhancement; PRN: Pseudo-random noise; PSD: Power spectral density; RPA: Retarding potential analyzer; RTI: Rayleigh–Taylor instability; SROSS-C2: Stretched Rohini Satellite Series; TEC: Total electron content; QLat: Quasi-dipole latitude.

Acknowledgements
This study was financially supported by the International Science Program (ISP) of the Uppsala University, Sweden. The official Swarm website is http://earth.esa.int/swarm, and the ftp server for Swarm data distribution is ftp://swarm-earth.esa.int/. The authors acknowledge the ESA Swarm team for the Swarm mission. We are also thankful to the Air Force Research Laboratory (AFRL), USA, for supplying the GPS-SCINDA receiver used in this study. The authors also acknowledge the anonymous reviewers for their careful reading of our manuscript and many insightful comments and suggestions which improved the quality of this paper. The authors acknowledge the ESA European Space Agency and the Technical University of Denmark for financial support (grant 4000109587/13/INB).

Authors’ contributions
SA, SB and EJ designed and conducted the current research and prepared the manuscript. SA and EJ analyzed the GPS-SCINDA scintillation data. SA analyzed the Swarm election density data, interpreted the results and wrote the manuscript. SB participated in the Swarm data analysis, interpretation of the results, and editing of the draft. EJ was involved in the interpretation and draft editing as well as in applying for grants for the project. All authors read and approved the final manuscript.

Funding
This study was financially supported by the International Science Program (ISP) of the Uppsala University, Sweden (grant no. UGA 02).

Availability of data and materials
The official website of Swarm is http://earth.esa.int/swarm and the ftp server for Swarm data distribution is ftp://swarm-earth.esa.int. The authors acknowledge the ESA Swarm team for the Swarm mission. We are also thankful to the Air Force Research Laboratory (AFRL), USA, for supplying the GPS-SCINDA receiver used in this study. The authors also acknowledge the anonymous reviewers for their careful reading of our manuscript and many insightful comments and suggestions which improved the quality of this paper. The authors acknowledge the ESA European Space Agency and the Technical University of Denmark for financial support (grant 4000109587/13/INB).

Competing interests
The authors declare that they have no competing interests.

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Received: 10 February 2020 Accepted: 13 October 2020
Published online: 30 October 2020

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