The first deep X-ray and optical observations of the closest isolated radio pulsar

A. Tiengo*, R. Mignani†, A. De Luca*,**, P. Esposito‡, S. Mereghetti* and A. Pellizzoni‡

*INAF - Istituto di Astrofisica Spaziale e Fisica Cosmica - Milano, via E. Bassini 15, I-20133
Milano, Italy
†University College London, Mullard Space Science Laboratory, Holmbury St. Mary, Dorking,
Surrey RH5 6NT, UK
**IUSS - Istituto Universitario di Studi Superiori, viale Lungo Ticino Sforza 56, I-27100 Pavia,
Italy
‡INAF - Osservatorio Astronomico di Cagliari, località Poggio dei Pini, strada 54, I-09012
Capoterra, Italy

Abstract. With a distance of 170 pc, PSR J2144–3933 is the closest isolated radio pulsar currently
known. It is also the slowest and least energetic radio pulsar; indeed, its radio emission is difficult to
account for with standard pulsar models, since its position in the $P - \dot{P}$ diagram is far beyond typical
“death lines”. Here we present the first deep X-ray and optical observations of PSR J2144–3933,
performed in 2009 with XMM-Newton and the VLT, from which we can set one of the most robust
upper limits on the surface temperature of a neutron star. We have also explored the possibility
of measuring the neutron star mass from the gravitational lensing effect on a background optical
source.

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INTRODUCTION

The timing parameters of PSR J2144–3933 are unique among the ~2000 pulsars currently
known: it is the radio pulsar with the longest spin period ($P=8.51$ s, [1]), with the
lowest rotational energy loss ($\dot{E}_{\text{rot}} = 4\pi^2 I \dot{P}/P^3 \approx 2.6 \times 10^{28}$ erg s$^{-1}$) and the farthest
below the radio pulsar death line (see Figure 1, left panel), where radio emission should
be inhibited ([2]). The interest in this source further increased when accurate distance
($172^{+20}_{-15}$ pc, corrected for the Lutz–Kelker bias [3]) and proper motion ($\mu = 166 \pm 1$ mas
yr$^{-1}$) were measured with VLBI observations at the Australian Long Baseline Array [4].
Although PSR J2144–3933 is the closest isolated radio pulsar known to date, it has never
been detected outside of the radio waveband. The pulsar sky region was serendipitously
observed for 5 ks with ROSAT/HRI [5] in 1997, for 20 ks in a EUVE deep survey observation [6] and for 200 s by GALEX [7] during the all-sky imaging survey. We report the
results of the first deep observations of PSR J2144–3933 in the X-ray and optical bands
with XMM-Newton and the VLT (results are presented in more detail in [8]).
OBSERVATIONS AND RESULTS

X-ray observation

PSR J2144–3933 was observed by XMM-Newton for about 40 ks on 2009 October 24. After filtering out the time intervals affected by a high level of particle background, the net exposure times were 22 ks for the EPIC PN [9] and 26 ks for the two EPIC MOS cameras [10]. After standard event selection, X-ray images were produced in several energy bands and analyzed with different source detection algorithms, but no X-ray source was found at the position of PSR J2144–3933 (see Figure 1, middle panel). The 3σ upper limit on the source net count rate in a 10′′ radius circle in the 0.2–10 keV energy range is 1.5×10⁻³ and 8.3×10⁻⁴ counts/s for the PN and the sum of the two MOS, respectively. The count rate upper limit for the PN in the 0.2-1 keV energy band, where our observation is most sensitive to the presumably soft X-ray emission of PSR J2144–3933, is 4.9×10⁻⁴ counts/s.

Optical observation

We observed PSR J2144–3933 with the VLT at the ESO Paranal observatory on 2009 August 21 with the FORS2 blue-sensitive CCD detector at the Antu UT1 telescope. The integration times were 8850 s with the U and B filters, and 2950 s with the V filter. As can be seen in Figure 1 (right panel), a relatively bright (U = 23.39 ± 0.05, B = 23.82 ± 0.03, V = 23.76 ± 0.05) object is detected 1.2′′ from the proper-motion-corrected position of PSR J2144–3933. Considering that the 1σ uncertainty of our astrometry is 0.24′′ and that the source is slightly extended, this object is very likely unrelated to the pulsar and it is probably a background galaxy. After correcting for atmospheric extinction and taking into account the contamination from the nearby source, the 3σ upper limits to the emission of PSR J2144–3933 are U > 25.3, B > 26.6, and V > 25.5.

DISCUSSION AND CONCLUSIONS

Assuming a column density of NH = 10²⁰ cm⁻² and an interstellar reddening E(B−V) = 0.02, the non detection of PSR J2144–3933 in our deep X-ray and optical observations corresponds to the following 3σ upper limits on the surface temperature (measured at infinity, assuming a blackbody spectrum): 2.3×10⁵ K for a 13 km radius neutron star, 4.4×10⁵ K for a 500 m radius hot spot, and 1.9×10⁶ K for a 10 m radius polar cap. The upper limits on the luminosity of the pulsar non-thermal emission (assuming a power-law spectrum with photon index Γ = 2) are instead 7×10²⁷ erg s⁻¹ in the 0.5–2 keV energy band and 5.6×10²⁶ erg s⁻¹ in the B optical band.

Given the small and well-determined distance of PSR J2144–3933 and the high quality of our observations, these limits are very robust and their values are among the lowest ever obtained for a neutron star. However, the large characteristic age (τc = P/(2P) ≃ 3.4×10⁸ years) and extremely low rotational energy loss of PSR J2144–3933 make them not particularly constraining both for the study of neutron star cooling evolution
FIGURE 1. Left panel: $P - \dot{P}$ diagram for radio pulsars (dots), radio-quiet $\gamma$-ray pulsars (diamonds), RRATs (crosses), XDINSs (triangles), AXPs/SGRs (squares), and PSR J2144–3933 (large star). Data from the ATNF Pulsar Catalogue ([11], http://www.atnf.csiro.au/research/pulsar/psrcat) and from the McGill SGR/AXP Online Catalog (http://www.physics.mcgill.ca/pulsar/magnetar/main.html). The arrow indicates the $\dot{P}$ upper limit for SGR 0418+5729 [12]. Dotted lines represent dipole magnetic fields from $10^9$ to $10^{15}$ G and characteristic ages from $10^3$ to $10^{10}$ years; the dashed line is a typical death line ($B/P^2 = 1.7 \times 10^{14}$ G s$^{-2}$, [13]). Middle panel: XMM-Newton PN (slightly smoothed) image of the $\sim 1.5' \times 3'$ field around PSR J2144–3933 in the 0.2–1 keV energy range. The $10'' \times 20''$ box indicates the sky region shown in the right panel. Right panel: ESO/VLT B-filter image of the $\sim 10'' \times 20''$ field around PSR J2144–3933. The pulsar proper-motion-corrected position is marked by a circle with a radius corresponding to the image astrometric uncertainty (0.24$''$).

and pulsar non-thermal emission processes: in fact, the surface temperature is expected to drop well below $10^5$ K after $\sim 10^7$ years and typical efficiencies of radio pulsars in converting rotational energy into the X-ray and optical radiation are much lower than our upper limits of 30% and 2%, respectively. Similarly, polar cap emission (for polar caps with radii $> 10$ m) could have been detected by our X-ray observation only for an unprecedented polar cap emission efficiency $> 40\%$.

On the other hand, the peculiar position of PSR J2144–3933 in the $P - \dot{P}$ diagram (see Figure 1, left panel), in the period range of AXPs/SGRs [14] and XDINSs [15], but with a weaker dipolar magnetic field ($B = (3c^2I/8\pi R^6 P\dot{P})^{1/2} \approx 1.9 \times 10^{12}$ G), might indicate that PSR J2144–3933 was born as a magnetar, but its magnetic field has now decayed to an intensity typical of radio pulsars. In such a case, it would be significantly younger than indicated by its characteristic age and would have also been heated by the decay of its field [16]. The recent discovery that the transient SGR 0418+5729 has a small period derivative [12] has shown that a pulsar with timing parameters similar to those of PSR J2144–3933 (see the left panel of Figure 1) can be a bright X-ray source.

The background galaxy that worsened our sensitivity to detect PSR J2144–3933 in the VLT images, might potentially be a tool to measure the neutron star mass through gravitational lensing (see, e.g., [17]). However, the displacement of a background source
Δφ = 1.2" away from a M = 1.4 M⊙ neutron star at a distance D = 170 pc is expected to be only δφ = \frac{4GM}{c^2DΔφ} ≃ 0.06 mas, which is beyond the capability of existing optical instruments. The situation is not going to improve in the next years, since the pulsar proper motion is directed away from the nearby source. A second gravitational lensing effect that would be in principle observable in this case, is the formation of a secondary ghost image of the galaxy close to the position of the pulsar, with a flux smaller than the one of the original source by a factor A− = 0.5 – \frac{u^2+2}{2u(u^2+4)^{1/2}}, where u = Δφ \left( \frac{4GM}{c^2D} \right)^{-1/2}.

For a canonical neutron star mass, the B = 23.8 magnitude background galaxy would have a secondary image of magnitude B = 45, while only a lens of at least 2 × 10^4 M⊙ would produce a detectable image in our VLT data! Not a particularly constraining upper limit for a neutron star mass, but the first one reported to date with this method...

A detectable lensing effect, that would lead to a measure of the neutron star mass, would instead be possible in case deeper observations of this field will be able to find dimmer sources along the sky trajectory of PSR J2144–3933. This pulsar is an ideal target for this challenging measurement thanks to its small distance, its very well determined position and proper motion and its vanishing flux outside the radio band.

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