SEARCH FOR CP AND CPT VIOLATION IN
POSITRONIUM DECAY

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Positronium, the electron-positron bound state, is described to a good approximation by pure QED. The states of positronium have definite values of C and P. Consequently, positronium is an attractive system for the investigation of possible violations of the discrete symmetries in the leptonic sector. We discuss signals for CP and CPT violation in the decay of spin-polarized orthopositronium and show where such correlations might arise in the context of the Standard-Model Extension.

1. Introduction

Positronium has long been used for precision tests of QED and in searches for violations of the discrete symmetries C, P, and T, singly and in various combinations. The physics of positronium is governed almost completely by QED, so that any sizable discrepancies between predictions and measured quantities would be signs of new physics. The bound states of positronium have simple properties under P and C, which makes positronium a particularly attractive system for tests of the discrete symmetries.

Tests of the discrete symmetries in positronium have taken three basic forms:

1. Search for decays to the wrong number of photons. (See Ref. 1 for a recent review.) Since positronium states are eigenstates of C, with eigenvalues \( C = (-1)^{l+s} \), ground state spin-singlet positronium—parapositronium—has even parity under C, and decays to an even number of photons, while ground state spin-triplet positronium—orthopositronium (o-Ps)—has odd C parity, and decays to an odd number of photons.
2. Tests involving transition rates.\(^2\) These can involve searches for forbid-
den transitions or for asymmetries caused by symmetry violation.

Tests involving correlations of the momenta or polarizations of the final state photons for decays from spin-polarized orthopositronium (see Fig. 1). Some measurable correlations are allowed, such as $\langle \hat{k}_i \cdot \hat{k}_j \rangle$ or $\langle \vec{S} \cdot \vec{h}_1 \rangle$ where $\hat{k}_i$ is a unit vector in the direction of the momentum of the $i^{th}$ final state photon, $\vec{h}_i$ describes the helicity of the $i^{th}$ photon, and $\vec{S}$ describes the orthopositronium spin. (The photons are numbered in order of decreasing energy.) QED predictions for the allowed correlations can be obtained using standard methods and provide both an opportunity for precision tests of bound state QED and known quantities for testing techniques used to search for the forbidden correlations. Correlations such as $\langle (\vec{S} \cdot \hat{k}_1)(\vec{S} \cdot \hat{k}_1 \times \hat{k}_2) \rangle$ and $\langle \vec{S} \cdot \hat{k}_1 \times \hat{k}_2 \rangle$ are forbidden and a non-vanishing value (apart from the effects of final state interactions) would be evidence for symmetry violation.

Here we review the most recent (null) measurements of the symmetry violating correlations and discuss their dependence on the various Standard-Model Extension (SME) coefficients.

2. Symmetry violating correlations

Two symmetry violating correlations have been extensively studied:

2.1. $\langle \vec{S} \cdot \hat{k}_1 \times \hat{k}_2 \rangle$

This correlation is CP symmetric but violates CPT. A non-vanishing correlation implies a difference between the number of decays with the normal to the decay plane parallel (+) to the spin direction and the number with
the normal antiparallel (−) to the spin direction. The measured quantity is the asymmetry

\[ A = \frac{N_+ - N_-}{N_+ + N_-}. \]  

(1)

The CPT violating coefficient was obtained from \( A \) by including a factor representing the “analyzing power” of the experiment. The asymmetry was measured by several groups starting with Arbic et al. in 1988.\(^7\) They found the CPT violating angular coefficient

\[ C_n = 0.020 \pm 0.023. \]  

(2)

Andrukhovich, Antovich, and Berestov reported a measurement of the asymmetry in 2000: \( A = 0.0008 \pm 0.00091 \).\(^8\) This was quoted by Vetter and Freedman\(^9\) as giving a coefficient of

\[ C_n = 0.0140 \pm 0.0190. \]  

(3)

Finally, the 2003 result of Vetter and Freedman,\(^9\) measured using the Gammasphere, was

\[ C_n = 0.0071 \pm 0.0062. \]  

(4)

2.2. \( \langle (\hat{S} \cdot \hat{k}_1)(\hat{S} \cdot \hat{k}_1 \times \hat{k}_2) \rangle \)

This correlation violates P, T, and CP, but not CPT. It was first measured by Skalsey and Van House in 1991.\(^10\) They found a CP violating coefficient of

\[ C_{\text{CP}} = -0.0056 \pm 0.0154. \]  

(5)

The most recent measurement was by Yamazaki et al. in 2010.\(^11\) Their result was

\[ C_{\text{CP}} = 0.0013 \pm 0.0022, \]  

(6)

giving a factor of seven improvement. Both results are consistent with zero.

3. Analysis of the symmetry-violating correlations

The Standard Model (SM) gives small but non-vanishing contributions for the correlations of Sec. 2. The operator \( \hat{S} \cdot \hat{k}_1 \times \hat{k}_2 \) is odd under time reversal, but the o-Ps states are unstable and are not eigenstates of T. The correlation \( \langle \hat{S} \cdot \hat{k}_1 \times \hat{k}_2 \rangle \) has a contribution at the \( O(\alpha^2) \) level arising from final-state interactions involving standard QED light-by-light scattering.\(^5\)
The CP-violating correlation \( \langle (\hat{S} \cdot \hat{k}_1)(\hat{S} \cdot \hat{k}_1 \times \hat{k}_2) \rangle \) has small non-vanishing weak interaction contributions.5

The Standard Model Extension provides a context for thinking about CPT and CP violation beyond the SM. In distinction to SM effects, SME contributions would be expected to depend on sidereal time. The discrete symmetry properties of the correlation operators of Sec. 2 are shown in Table 1. The SME operators that share these symmetry properties have coefficients \( b_j, g_{jk0}, (k_{AF})_j \) for \( \hat{S} \cdot \hat{k}_1 \times \hat{k}_2 \) and \( c_{0j}, c_{j0}, (k_{F})_{0jk} \) for \( (\hat{S} \cdot \hat{k}_1)(\hat{S} \cdot \hat{k}_1 \times \hat{k}_2) \).12 Explicit results for the symmetry violating correlations in the SME would be expected to involve these coefficients. The corresponding calculations are underway.

Table 1. Discrete symmetry properties of the correlation operators.

| Operator                      | C | P | T | CP | CPT |
|--------------------------------|---|---|---|----|-----|
| \( \hat{S} \cdot \hat{k}_1 \times \hat{k}_2 \) | + | + | - | +  | -   |
| \( (\hat{S} \cdot \hat{k}_1)(\hat{S} \cdot \hat{k}_1 \times \hat{k}_2) \) | + | - | - | -  | +   |

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